


CMR INSTITUTE OF TECHNOLOGY		USN <input type="text"/>									
Internal Assessment Test - I								CMR			
Sub:	ANTENNAS AND PROPAGATION						Code:	10EC64			
Date:	28/03/2017	Duration:	90 mins	Max Marks:	50	Sem:	VI	Branch:	ECE(C,D)		
Answer Any FIVE FULL Questions											
								Marks		OBE	
										CO	RBT
1.	<p>Explain the following terms as related to antenna system:</p> <p>(i) Directivity, (ii) Effective Height, (iii) Beam Area, (iv) Effective Aperture</p> <p>(i)</p> <p>The directivity of an antenna is equal to the ratio of the maximum power density $P(\theta, \phi)_{max}$ (watts/m²) to its average value over a sphere as observed in the far field of an antenna.</p> $D = \frac{P(\theta, \phi)_{max}}{P(\theta, \phi)_{av}} \rightarrow \text{Directivity from pattern.}$ <p>The directivity is dimensionless ratio ≥ 1.</p> <p>The average power density over a sphere,</p> $P(\theta, \phi)_{av} = \frac{1}{4\pi} \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} P(\theta, \phi) \sin\theta d\theta d\phi$ $= \frac{1}{4\pi} \iint_{4\pi} P(\theta, \phi) d\Omega \quad (\text{watts}^{-1})$ $D = \frac{P(\theta, \phi)_{max}}{\left(\frac{1}{4\pi}\right) \iint_{4\pi} P(\theta, \phi) d\Omega} = \frac{1}{\frac{1}{4\pi} \iint_{4\pi} \left[\frac{P(\theta, \phi)}{P(\theta, \phi)_{max}}\right] d\Omega}$ $= \frac{4\pi}{\iint_{4\pi} P_n(\theta, \phi) d\Omega} = \frac{4\pi}{\Omega_A} \quad \text{Directivity from beam area } \Omega_A$ <p>where, $P_n(\theta, \phi) d\Omega = \frac{P(\theta, \phi)}{P(\theta, \phi)_{max}} = \text{normalized power pattern}$</p>						[10]	CO1	L1		

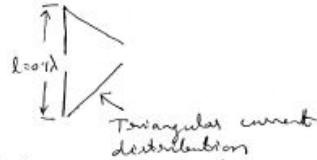
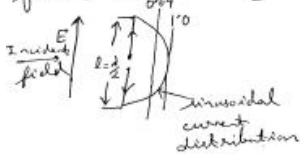
(ii)

Effective Height

Multiplying the effective height by the incident field E of the same polarization gives the voltage V induced.

$$\therefore V = hE$$

Accordingly, the effective height may be defined as the ratio of the induced voltage to the incident field, or $h = \frac{V}{E}$.



consider, a vertical dipole of length $l = \frac{\lambda}{2}$ immersed in an incident field E .

→ If the current distribution was uniform, effective height would be l .

→ But actual current distribution is nearly sinusoidal with an average value $\frac{2}{\pi} = 0.64$ (of the maximum)

\therefore The effective height $\approx h = 0.64l$.

Assumed antenna is oriented for maximum response.

Now say $l = 0.1 \lambda$

The current tapers almost linearly from the central fed point to zero at the ends in a triangular distribution.

The average current is $\frac{1}{2}$ of the maximum.

\therefore effective height = $0.5l$.

∴ Effective height can also be defined, by considering the transmitting case of the antenna and equating the effective height to the physical height (or length l) and multiplying by the ~~the~~ normalized average current.

$$h_e = \frac{1}{I_0} \int_0^{l_p} I(z) dz = \frac{I_{av}}{I_0} h_p \text{ (m)}$$

where, h_e = effective height (m)

h_p = physical ~ (m)

I_{av} = Average current (A)

Effective height can also be expressed as a vector quantity, For linear polarization,

$$V = \vec{h}_e \cdot \vec{E} = h_e E \cos \theta$$

h_e → effective height and ~~polarization angle of antenna~~

\vec{E} → Field Intensity ~ ~ ~ of incid. wave

θ → angle b/w polarization angles of antenna and wave in degree.

(iii)

Beam area or beam solid angle Ω_A of an antenna is given by the integral of the normalized power pattern over a sphere (4π sr)

$$\Omega_A = \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} P_n(\theta, \phi) \sin \theta d\theta d\phi$$

$$\text{and } \Omega_A = \int \int_{4\pi} P_n(\theta, \phi) d\Omega \text{ (sr)}$$

where, $d\Omega = \sin \theta d\theta d\phi$, sr.

The beam area Ω_A is the solid angle through which all of the power radiated by the antenna would stream if $P(\theta, \phi)$ maintained its maximum value over Ω_A and was zero elsewhere.

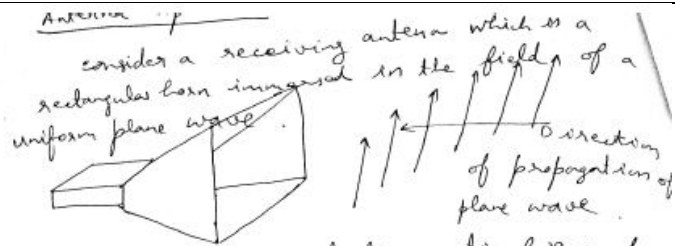
∴ The power radiated = $P(\theta, \phi) \Omega_A$ watts.

The beam area can be given approximately in terms of the angles subtended by the half-power points of the main lobe in the two principal planes.

$$\therefore \text{Beam area} \cong \Omega_A \cong \theta_{HP} \phi_{HP}$$

where, θ_{HP} and ϕ_{HP} are the HPBW in the two principal planes, minor lobes being neglected.

(iv)



plane wave incident on electromagnetic horn of physical aperture A_p .

Let the Poynting vector, or power density, of the plane wave be $S \text{ watts/m}^2$
 Let A the area, or physical aperture of the horn be $A_p \text{ m}^2$.

If the horn A extracts all the power from the wave over its entire physical aperture. Then the total power P absorbed from the wave is,

$$P = \frac{E^2}{Z} \cdot A_p = S A_p W \quad \text{①}$$

Thus, the horn considered to have an aperture, (the ~~the~~ total power it extracts from passing wave being proportional to the aperture or area of its mouth.

— But the field response of the horn is not uniform across the aperture A .

→ As \vec{E} at the sidewalls must be zero.
 \therefore The effective aperture A_e of the horn is less than the physical aperture A_p .

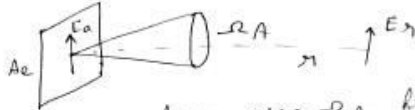
$$E_{AP} = \frac{A_e}{A_p} \text{ (dimensionless) Aperture efficiency}$$

For horn and parabolic reflector antenna, E_{AP} is in the range of 50 to 80%.

$$E_{AP}$$

$$(0.5 \leq E_{AP} \leq 0.8)$$

- Large dipole or patch arrays with uniform field to the edge of physical aperture may attain higher aperture efficiencies approaching 100%.



Radiation over beam area A_r from aperture A_e consider as antenna with effective aperture A_e . It radiates all its power in a conical pattern: beam area A_r .

Assuming uniform field E_a over the aperture, the power radiated is,

$$P = \frac{E_a^2}{Z_0} A_e \text{ (W)} \quad \dots (1)$$

where, Z_0 = Intrinsic impedance of medium (377 Ω for free space)

Assuming a uniform field E_r in the far field at a distance r , the power radiated is,

$$P = \frac{E_r^2}{Z_0} r^2 A_r \text{ (W)} \quad \dots (2)$$

$$A_e \therefore \frac{E_a^2}{Z_0} A_e = \frac{E_r^2}{Z_0} r^2 A_r$$

$$\text{Now } E_r = \frac{E_a A_e}{r \lambda}$$

$$\Rightarrow \frac{E_r^2}{E_a^2} = \frac{A_e^2}{r^2 \lambda^2}$$

$$\Rightarrow A_e = \frac{A_e^2}{\lambda^2} \cdot r^2 A_r$$

$$\Rightarrow \text{A}$$

$$A_e = \frac{E_r^2}{E_a^2} r^2 A_r$$

$$= \frac{A_e^2}{A_e^2} \cdot \frac{E_a^2}{\lambda^2} \cdot r^2 A_r$$

$$= \frac{A_e A_r}{\lambda^2}$$

$$\Rightarrow A_e = \frac{\lambda^2}{4\pi A_r}$$

$$\Rightarrow \lambda^2 = A_e A_r$$

Aperture beam area relation.

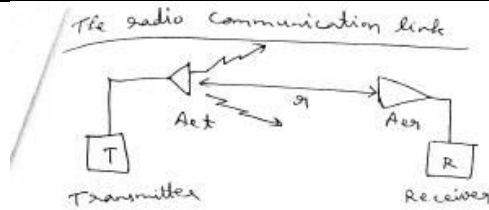
$$\text{Now, } D = \frac{4\pi}{A_r} = \frac{4\pi}{\frac{\lambda^2}{A_e}} = 4\pi \cdot \frac{A_e}{\lambda^2}$$

2.(a) State and prove Friis transmission formula.

[04]

CO1

L1



Free Transmission formula

Assume lossless, matched antennas.
 Let, Tx feed a power P_t to a transmitting antenna of effective aperture A_{et} .

At a dist. r a receiving antenna is placed.
 A_{er} → effective aperture of receiving antenna.
 $\therefore A_{er}$ intercepts some power radiated by the transmitting antenna of effective aperture A_{et} and delivers to the receiver R.

Assume transmitting antenna is isotropic.
 \therefore Power/area available at the receiving antenna is,

$$S_r = \frac{P_t}{4\pi r^2} \text{ W/m}^2 \quad \text{--- (1)}$$

Now any antenna has gain G_t .
 The power available/unit area,

$$S_r = \frac{P_t G_t}{4\pi r^2} \text{ W/m}^2 \quad \text{--- (2)}$$

Now the power collected by the lossless, matched receiving antenna of effective aperture A_{er}

$$P_r = S_r A_{er} = \frac{P_t G_t}{4\pi r^2} A_{er} \quad \text{--- (3)}$$

Now, gain of the transmitting antenna can be expressed as,
 $G_t = \frac{4\pi A_{et}}{\lambda^2}$ [if antenna is lossless, $e_{ca} = 1 \Rightarrow G_t = e_{ca} D = D$]

$$\therefore P_r = \frac{P_t A_{er}}{4\pi r^2} \cdot \frac{4\pi A_{et}}{\lambda^2}$$

$$\Rightarrow \frac{P_r}{P_t} = \frac{A_{er} A_{et}}{r^2 \lambda^2} \quad \text{(dimensionless)}$$

Free Transmission formula.

where, P_r → received power W
 P_t → transmitted " W
 A_{et} → effective aperture of transmitting antenna m^2
 A_{er} → effective aperture of receiving antenna m^2
 r = distance b/w antennas
 λ = wavelength.

2.(b) Calculate the directivity for the following sources with power patterns:

- (i) $U = U_m \sin^2 \theta \sin^2 \varphi$; for $0 \leq \theta \leq \pi$ and $0 \leq \varphi \leq \pi$
 $= 0$; elsewhere

[06] CO1 L3

(ii) $U = U_m \sin^2 \theta \sin^3 \phi$; for $0 \leq \theta \leq \pi$ and $0 \leq \phi \leq \pi$
 $= 0$; elsewhere

2.(b) (i) $v = U_m \sin^2 \theta \sin^3 \phi$

$$\therefore \mathcal{R}_A = \iint p_n(\theta, \phi) \sin \theta \, d\theta \, d\phi$$

$$\therefore \mathcal{D} = \frac{4\pi U_m}{\mathcal{R}_A}$$

$$P_{rad} = \int_{\theta=0}^{\pi} \int_{\phi=0}^{\pi} U_m \sin^2 \theta \sin^3 \phi \cdot \sin \theta \, d\theta \, d\phi$$

$$= U_m \int_{\theta=0}^{\pi} \sin^3 \theta \, d\theta \int_{\phi=0}^{\pi} \sin^3 \phi \, d\phi$$

$$= U_m \cdot \frac{4}{3} \cdot \frac{\pi}{2} = \frac{4\pi}{6} U_m$$

$$\therefore \mathcal{D} = \frac{4\pi U_m}{\frac{4\pi}{6} U_m} = 6$$

Now, $\int_{\theta=0}^{\pi} \sin^3 \theta \, d\theta$

$$= \int_{\theta=0}^{\pi} \frac{3}{4} \sin \theta \, d\theta - \frac{1}{4} \int_{\theta=0}^{\pi} \sin^3 \theta \, d\theta$$

$$= \frac{3}{4} [-\cos \theta]_0^{\pi} + \frac{1}{4} \left[\frac{\cos^3 \theta}{3} \right]_0^{\pi}$$

$$= \frac{3}{4} [\cos \theta]_0^{\pi} + \frac{1}{12} [\cos^3 \pi - \cos^3 0]$$

$$= \frac{3}{4} [1 + 1] + \frac{1}{12} [-1 - 1]$$

$$= \frac{6}{4} - \frac{2}{12} = \frac{18-2}{12} = \frac{16}{12} = \frac{4}{3}$$

Also, $\int_{\phi=0}^{\pi} \sin^3 \phi \, d\phi = \frac{1}{2} \int_{\phi=0}^{\pi} (1 - \cos^2 \phi) \, d\phi$

$$= \frac{1}{2} [\pi] - \frac{1}{2} \left[\frac{\sin^2 \phi}{2} \right]_0^{\pi} = \frac{\pi}{2}$$

$$[\sin^3 \theta = 3 \sin \theta - 4 \sin^3 \theta$$

$$\Rightarrow 4 \sin^3 \theta = 3 \sin \theta - \sin^3 \theta$$

$$\Rightarrow \sin^3 \theta = \frac{3}{4} \sin \theta - \frac{1}{4} \sin^3 \theta$$

$$[\cos^2 \theta = 1 - 2 \sin^2 \theta$$

$$\Rightarrow 2 \sin^2 \theta = 1 - \cos^2 \theta$$

$$\Rightarrow \sin^2 \theta = \frac{1}{2} (1 - \cos^2 \theta)$$

2.(b)(ii) $U = U_m \sin^2 \theta \sin^3 \phi$

$$\mathcal{D} = \frac{4\pi}{\mathcal{R}_A} = \frac{4\pi}{\int_{\theta=0}^{\pi} \int_{\phi=0}^{\pi} p_n(\theta, \phi) \sin \theta \, d\theta \, d\phi}$$

$$= \frac{4\pi}{\int_{\theta=0}^{\pi} \int_{\phi=0}^{\pi} \sin^2 \theta \sin^3 \phi \sin \theta \, d\theta \, d\phi}$$

$$= \frac{4\pi}{\int_{\theta=0}^{\pi} \sin^3 \theta \, d\theta \int_{\phi=0}^{\pi} \sin^3 \phi \, d\phi} = \frac{4\pi}{\frac{4}{3} \times \frac{4}{3}} = \frac{9\pi}{4} = 7.06$$

3.(a) Show that maximum effective aperture of short dipole is $0.119\lambda^2$.

[06]

CO3

L3

Soln. The maximum effective aperture of an antenna
 λ^2 ,
 $A_{em} = \frac{V^2}{4SR_n}$

where, $V = \epsilon l$

where, $\epsilon \rightarrow$ effective electric field intensity
 $l \rightarrow$ length of the dipole,

The radiation resistance R_r of a short dipole of length l with uniform current is,

$$R_r = \frac{80\pi^2 l^2}{\lambda^2} \left(\frac{I_{av}}{I_0}\right)^2$$

$$= 790 \left(\frac{I_{av}}{I_0}\right)^2 \left(\frac{l}{\lambda}\right)^2$$

where $\lambda \rightarrow$ wavelength

$I_{av} \rightarrow$ average current

$I_0 \rightarrow$ terminal current

The power density, or Poynting vector, of the incident wave at the dipole is related to the field intensity by,

$$S = \frac{E^2}{Z} = \frac{E^2}{120\pi}$$

where, $Z \rightarrow$ intrinsic impedance of the medium

Here medium is free space,

$$\therefore Z = 120\pi \Omega$$

$$\therefore A_{em} = \frac{V^2}{4SR_n} = \frac{V^2}{4} \cdot \frac{120\pi}{E^2} \cdot \frac{\lambda^2}{80\pi^2 l^2} \cdot \left(\frac{I_0}{I_{av}}\right)^2$$

The maximum effective aperture of a short dipole (for $I_{av} = I_0$),

$$A_{em} = \frac{V^2 \cdot 120\pi}{4 E^2} \cdot \frac{\lambda^2}{80\pi^2 l^2}$$

$$= \frac{120\pi \lambda^2}{320\pi^2} = \frac{120}{32\pi} \lambda^2 = 0.119 \lambda^2$$

$$(10) D = \frac{4\pi A_e}{\lambda^2} = \frac{4\pi \times 0.119 \lambda^2}{\lambda^2} = 1.5$$

3.(b) Show that the directivity for unidirectional operation is $2(n+1)$ for an intensity variation of $U = U_m \cos^n \theta$.

[04]

CO1

L3

$$U = U_m \cos^n \theta, \quad D = \frac{4\pi}{\int \int P_n(\theta, \phi) d\Omega}$$

$$= \frac{4\pi}{\frac{1}{2} \int_0^{2\pi} \int_0^{\pi} \cos^n \theta \sin \theta d\theta d\phi}$$

$$\theta=0 \quad \phi=0$$

$$= \frac{4\pi \cdot 2}{2\pi \int_0^{\pi/2} \cos^n \theta \sin \theta d\theta}$$

$$= \frac{2}{\frac{\cos(\theta) \cdot \frac{1}{n+1}}{n+1} \Big|_0^{\pi/2}} = 2(n+1)$$

$$\int_0^{\pi/2} \cos^n \theta \sin \theta d\theta$$
 set, $\cos \theta = z$
 $\Rightarrow -\sin \theta d\theta = dz$
 $\theta=0 \quad z=1$
 $\theta=\pi/2 \quad z=0$

$$\therefore \int_0^{\pi/2} \cos^n \theta \sin \theta d\theta = - \int_1^0 z^n dz = \left[\frac{z^{n+1}}{n+1} \right]_1^0 = \frac{1}{n+1}$$

4. Derive the expression for total field for two isotropic point sources with same amplitude and opposite phase. Plot the field pattern when the two isotropic sources are spaced $\lambda/2$ apart.

$$E = E_0 e^{j\psi} - E_0 e^{-j\psi}$$

$$= 2E_0 \cdot \frac{1}{2} (\cos \psi/2 + j \sin \psi/2 - \cos \psi/2 + j \sin \psi/2)$$

$$= 2E_0 \cdot \frac{1}{2} 2j \sin \psi/2 = 2E_0 j \sin \psi/2$$

$$= 2j E_0 \sin \left(\frac{d}{a} \cos \phi \right) \quad \text{--- (1)}$$

The presence of j in (1) indicates the phase reversal of one of the sources results in a 90° phase shift of the total field as compared with the total field for case 1.

Taking $2jE_0 = 1$ and considering the special case of $d = \lambda/2$,

$$E = \sin \left(\frac{\lambda}{2} \cos \phi \right) \quad \text{--- (2)}$$

[10] CO2 L3

The directions ϕ_m of maximum field are obtained when,

$$\frac{\pi}{2} \cos \phi_m = \pm (2k+1) \frac{\pi}{2}$$

where, $k=0, 1, 2, 3 \dots$

when, ~~$k=0, \phi_m = \pm 90^\circ$~~

$$\frac{\pi}{2} \cos \phi_m = \pm \frac{\pi}{2}$$

$$\Rightarrow \cos \phi_m = \pm 1$$

$$\therefore \phi_m = 0^\circ \text{ and } 180^\circ \dots 3(a)$$

The null directions ϕ_0 are given by,

$$\frac{\pi}{2} \cos \phi_0 = \pm k\pi$$

$$\text{for } k=0, \phi_0 = \pm 90^\circ$$

$$\Rightarrow \phi_0 = \pm 90^\circ \dots 3(b)$$

The half power directions are given by,

$$\sin\left(\frac{\pi}{2} \cos \phi\right) = \frac{1}{\sqrt{2}} \text{ considering } 2\epsilon_0 = 1$$

$$\frac{\pi}{2} \cos \phi = \pm (2k+1) \frac{\pi}{4}$$

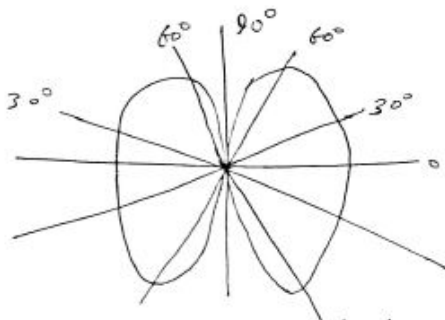
$$\text{for } k=0, \frac{\pi}{2} \cos \phi = \pm \frac{\pi}{4}$$

$$\Rightarrow \cos \phi = \pm \frac{1}{2}$$

$$\Rightarrow \phi = \pm 60^\circ, \pm 120^\circ \dots 3(c)$$

The field patterns as in (3) are shown in fig. 1.

The pattern is a relatively broad figure-of-eight with the maximum field in the same direction as the line joining the sources (x-axis).



The pattern is a relatively broad figure-of-eight ~~rotation of~~ with the maximum field in the same direction as the line joining the two sources (x-axis).

- The space pattern is a fig. of revolution about the x-axis.

- The two sources can be described as ~~the~~ a simple-type of "end-fire" array.

5. Draw the field pattern of a broadside array of 5 elements. Each element is at a distance of $\lambda/2$ from its neighbour. Assume $\delta = -\pi$.

[10] CO2 L3

B) Broadside array. $n=5$, $d = \frac{\lambda}{2}$, $\delta = -\pi$

$\psi = 0$

Condition for maxima
 $\psi = 0$
 $\Rightarrow n \cos \phi = 0$
 $\Rightarrow \phi = 0$
 $\Rightarrow \phi = \pm 90^\circ$

$d_s = \frac{\lambda}{2} d = \frac{\lambda}{2} \cdot \frac{\lambda}{2} = \frac{\lambda^2}{4} = \pi$

Condition for null:

$$\frac{n\psi}{2} = \pm 2k\pi$$

$$\frac{5}{2} (\pi \cos \phi) = \pm 2k\pi$$

$$\Rightarrow \cos \phi = \pm \frac{2k}{5}$$

$$\Rightarrow \phi = \pm \cos^{-1} \left[\pm \frac{2k}{5} \right]$$

for $k=1$, $\phi = \pm \cos^{-1} \left[\pm \frac{2}{5} \right]$
 $\Rightarrow \phi = \pm 66.42^\circ$
 and $\phi = \pm 113.57^\circ$

for $k=2$, $\phi = \pm \cos^{-1} \left[\pm \frac{4}{5} \right]$
 $\Rightarrow \phi = \pm 36.869^\circ$
 and $\phi = \pm 143.13^\circ$

for $k=3$, \rightarrow discarded

Direction of side lobes

$$\frac{n\psi}{2} = \pm (2k+1) \frac{\pi}{2}$$

$$\Rightarrow \frac{5}{2} \pi \cos \phi = \pm \frac{(2k+1)\pi}{2}$$

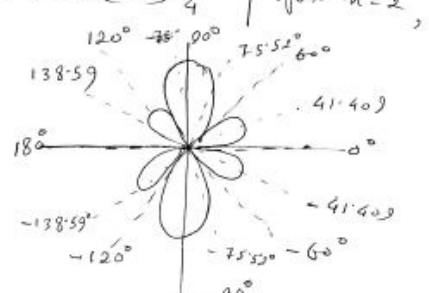
$$\Rightarrow 5 \cos \phi = \pm (2k+1)$$

$$\Rightarrow \cos \phi = \pm \frac{1}{5} (2k+1)$$

$$\Rightarrow \phi = \pm \cos^{-1} \left[\pm \frac{1}{5} (2k+1) \right]$$

for $k=1$, $\phi = \pm \cos^{-1} \left[\pm \frac{3}{5} \right]$
 $\Rightarrow \phi = \pm \cos^{-1} \left(\frac{3}{5} \right) = \pm 53.130^\circ$
 and $\phi = \pm \cos^{-1} \left(-\frac{3}{5} \right) = \pm 126.86^\circ$

for $k=2$, $\phi = \pm \cos^{-1} \left[\pm 1 \right]$
 $\Rightarrow \phi = \pm \cos^{-1} [1] = 0^\circ$
 and $\phi = \pm \cos^{-1} [-1] = 180^\circ$

6.	<p>Find peaks, nulls and beam-width of a given array of two point sources with equal amplitude and equal phase. Given $d = \lambda/2$ and $n=4$. Draw the field pattern.</p> <p>For major lobe, $\psi = 0$ Here, $d = \frac{\lambda}{2}$, $n = 4$, $\therefore \psi = \frac{2\pi}{\lambda} d \cos \phi + \delta$ $\delta = 0$ $= \frac{2\pi}{\lambda} \cdot \frac{\lambda}{2} \cos \phi + 0$ $= 2\pi \cos \phi$</p> <p>For any major lobe, $\psi = 0$ $\Rightarrow 2\pi \cos \phi = 0$ $\Rightarrow \phi = \pm 90^\circ$</p> <p>For minima, $\frac{n\psi}{2} = \pm k\pi$ $\Rightarrow \frac{2}{2} \pi \cos \phi = \pm k\pi$ $\Rightarrow \pi \cos \phi = \pm k\pi$</p> <p>For side-lobes, $\frac{n\psi}{2} = \pm (2k+1) \frac{\pi}{2}$ $\Rightarrow \cos \phi = \pm (2k+1) \frac{1}{4}$</p> <p>for $k=0$, $\phi = \pm 75.52^\circ$ for $k=1$, $\phi = \pm 41.409^\circ, \pm 138.59^\circ$ for $k=2$, discarded</p>  <p>BWFN = 60° HPBW = 30°</p>	[10]	CO2	L3
7.	<p>Derive an expression for the radiation resistance of a short dipole.</p> <p>The average Poynting vector is given by, $\vec{S} = \frac{1}{2} \text{Re} (\vec{E} \times \vec{H}^*)$ --- (1)</p> <p>The far-field components are E_θ, H_ϕ. \therefore The radial component of the Poynting vector $S_r = \frac{1}{2} \text{Re} E_\theta H_\phi^*$ --- (2)</p> <p>where, E_θ and H_ϕ^* are complex.</p> <p>Now, $\frac{E_\theta}{H_\phi} = Z$ $\Rightarrow E_\theta = H_\phi Z = H_\phi \sqrt{\frac{\mu}{\epsilon}}$</p> <p>$\therefore$ Equation (2) becomes, $S_r = \frac{1}{2} \text{Re} Z H_\phi H_\phi^*$ $= \frac{1}{2} H_\phi ^2 \text{Re} Z$ $= \frac{1}{2} H_\phi ^2 \sqrt{\frac{\mu}{\epsilon}}$ --- (3)</p>	[10]	CO3	L1

The total power P radiated is,

$$P = \iint S_n ds = \frac{1}{2} \sqrt{\frac{\mu}{\epsilon}} \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} |H_\phi|^2 r^2 \sin\theta d\theta d\phi \quad (4)$$

Now, $|H_\phi| = \frac{\omega I_0 L \sin\theta}{4\pi a^2 r}$ (5)

Substituting this into (4),

$$P = \frac{1}{2} \sqrt{\frac{\mu}{\epsilon}} \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \frac{\omega^2 I_0^2 L^2}{16\pi^2 a^4 r^2} \sin^2\theta \cdot r^2 \sin\theta d\theta d\phi$$

$$= \frac{1}{2} \sqrt{\frac{\mu}{\epsilon}} \cdot \frac{\omega^2 I_0^2 L^2}{16\pi^2 a^4} \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \sin^3\theta d\theta d\phi$$

$$= \frac{1}{2} \sqrt{\frac{\mu}{\epsilon}} \cdot \frac{\beta^2 I_0^2 L^2}{16\pi^2} \cdot 2\pi \cdot \frac{4}{3}$$

$$= \sqrt{\frac{\mu}{\epsilon}} \frac{\beta^2 I_0^2 L^2}{12\pi} \quad (6)$$

This is the average power or rate at which energy is streaming out of a sphere surrounding the dipole.

\therefore This is equal to the power radiated. Assuming no losses, it is also equal to the power delivered to the dipole.

$I_0 \rightarrow$ maximum current,

\therefore The corresponding r.m.s current = $\frac{I_0}{\sqrt{2}}$

$$\sqrt{\frac{\mu}{\epsilon}} \frac{\beta^2 I_0^2 L^2}{12\pi} = \left(\frac{I_0}{\sqrt{2}}\right)^2 R_r$$

$$\Rightarrow \sqrt{\frac{\mu}{\epsilon}} \frac{\beta^2 L^2}{6\pi} = R_r$$

$$\Rightarrow R_r = \sqrt{\frac{\mu}{\epsilon}} \frac{\beta^2 L^2}{6\pi}$$

For air or vacuum, $\sqrt{\frac{\mu}{\epsilon}} = \sqrt{\frac{\mu_0}{\epsilon_0}} = 120\pi \Omega$

\therefore Radiation resistance of a dipole with

uniform current is,

$$R_r = \frac{20}{120\pi} \cdot \frac{\beta^2 L^2}{6\pi}$$

$$= 20 \cdot \frac{(2\pi)^2}{\lambda^2} \cdot L^2 = 80\pi^2 \left(\frac{L}{\lambda}\right)^2$$

$$= 80\pi^2 L^2 \lambda^{-2} = 790 L^2 \lambda^{-2} (\Omega)$$

- It is assumed that with end-loading the dipole current is uniform.
- If no ^{end} loading dipole current is zero at the ends.
- If dipole is short, the current tapers off linearly from a max. at the center to zero at the ends.

Modifying eqn. (6) for this general case,

$$P = \sqrt{\frac{\mu}{\epsilon}} \frac{\beta^2 I_{av}^2 L^2}{12\pi} \quad (w) \quad (7)$$

The power delivered to the dipole is as before,

$$P = \frac{1}{2} I_0^2 R_A$$

where, $I_0 \rightarrow$ amp. of terminal current of center-fed dipole (peak value in time)

$$\therefore \frac{1}{2} I_0^2 R_A = \sqrt{\frac{\mu}{\epsilon}} \frac{\beta^2 I_{av}^2 L^2}{12\pi}$$

$$\begin{aligned} \therefore R_A &= \sqrt{\frac{\mu}{\epsilon}} \frac{\beta^2 L^2}{12\pi} \left(\frac{I_{av}}{I_0}\right)^2 \\ &= 790 \left(\frac{I_{av}}{I_0}\right)^2 L_A^2 \quad (\Omega) \end{aligned}$$

For a short dipole without end-loading,

$$I_{av} = \frac{I_0}{2}$$

$$\therefore R_A = 197 (L_A)^2$$

Scheme of solution
2017 IAT-1 AWP

1. (i) Directivity - 3m
 (ii) Effective height - 2m
 (iii) Beam area - 2m
 (iv) Effective Aperture - 3m

2. (a) Diagram - 1m
 Derivation & final formula - 3m

2. (b) (i) Approach - 2m
 Final answer - 1m
 (ii) Approach - 2m
 Final answer - 1m

3. (a) Diagram - 1m
 Derivation - 5m

3. (b) Approach - 2m
 Proof - 3m

4. Diagram - 2m
 Expression for total field - 3m
 Field pattern - 1m
 Maximum, null, side-lobe directions - 5m

5. Approach - 2m
 Maximum direction - 2m
 Null direction - 2m
 Field pattern - 2m
 Side-lobe direction - 2m

6. Approach - 2m
 Maximum (major-lobe direction) - 2m
 Null direction - 2m
 Side-lobe direction - 2m
 Field pattern - 2m

7. Expression for H_d - 2m
 Derivation - 6m
 Final result - 2m

Course Outcomes		PO1	PO2	PO3	PO4	PO5	PO6	PO7	PO8	PO9	PO10	PO11	PO12
CO1:	Explain the radiation mechanism of various types antennas.	3	0	1	0	0	0	0	0	0	0	0	0
CO2:	Construct array of two point sources and array of n-isotropic point sources and interpret the radiation pattern for various arrays	3	2	2	3	0	0	0	0	0	0	0	0
CO3:	Analyze dipole antennas and thin linear	3	2	2	3	0	0	0	0	0	0	0	0

	antennas and derive field components in far field antenna zone.												
CO4:	Derive the field components for loop, slot and horn antennas, assess the antennas and hence interpret their operation in broadcasting applications.	3	2	2	3	0	0	0	0	0	0	0	0
CO5:	Analyze the performances of high frequency antennas and hence rank the antennas according to the bandwidth requirement, resolution, use in remote sensing applications and broadcasting.	3	0	1	0	0	0	0	0	0	0	0	0
CO6:	Analyze ground wave, space wave, surface wave and ionospheric propagation, derive the field equations and summarize the performance in various regions of atmosphere.	3	0	1	0	0	0	0	0	0	0	0	0

Cognitive level	KEYWORDS
L1	List, define, tell, describe, identify, show, label, collect, examine, tabulate, quote, name, who, when, where, etc.
L2	summarize, describe, interpret, contrast, predict, associate, distinguish, estimate, differentiate, discuss, extend
L3	Apply, demonstrate, calculate, complete, illustrate, show, solve, examine, modify, relate, change, classify, experiment, discover.
L4	Analyze, separate, order, explain, connect, classify, arrange, divide, compare, select, explain, infer.
L5	Assess, decide, rank, grade, test, measure, recommend, convince, select, judge, explain, discriminate, support, conclude, compare, summarize.

PO1 - *Engineering knowledge*; PO2 - *Problem analysis*; PO3 - *Design/development of solutions*; PO4 - *Conduct investigations of complex problems*; PO5 - *Modern tool usage*; PO6 - *The Engineer and society*; PO7- *Environment and sustainability*; PO8 - *Ethics*; PO9 - *Individual and team work*; PO10 - *Communication*; PO11 - *Project management and finance*; PO12 - *Life-long learning*