

2nd Test Scheme of Evaluation – May 2017

Sub:	Smart Materials and Structures						Code:	14MST422	
Date:	23/05/2017	Duration:	90 mins	Max Marks:	50	Sem:	4th M.Tech	Section:	-

Question No	Description	Marks Allotted	
1	Explain in detail about the Frahm Absorber with a sketch	10	
2	Explain the characteristics of Perissogyro Vibration Absorber	10	
3	Explain how “mistuning” affects circularly symmetric structures	10	
4	Explain in details about Extrinsic Fabry-Perot Fibre Optic Sensors	10	
5	Explain in details about Brag Grating Fibre Optic Sensors	10	
6	Explain in detail about SMA controlled cantilever beam structure	10	
7	Explain in detail about Bingham Plastic model for ER/MR fluids	10	

Attempt any 5 questions.

CI

CCI

HOD

where X_{1st} is the static displacement f_0/k of the main mass. From (3a), it is clear that the vibratory amplitude of the main mass reaches zero if the natural frequency ω_2 of the absorber mass is tuned to be equal to the frequency of the forcing function at $\omega = \omega_1$ (i.e. $\omega_2 = \omega = \omega_1$). Thus, the vibratory motion of main structure is "absorbed" at its natural frequency, resulting in a null. However, the absorber mass itself vibrates with the amplitude given by eq (3b).

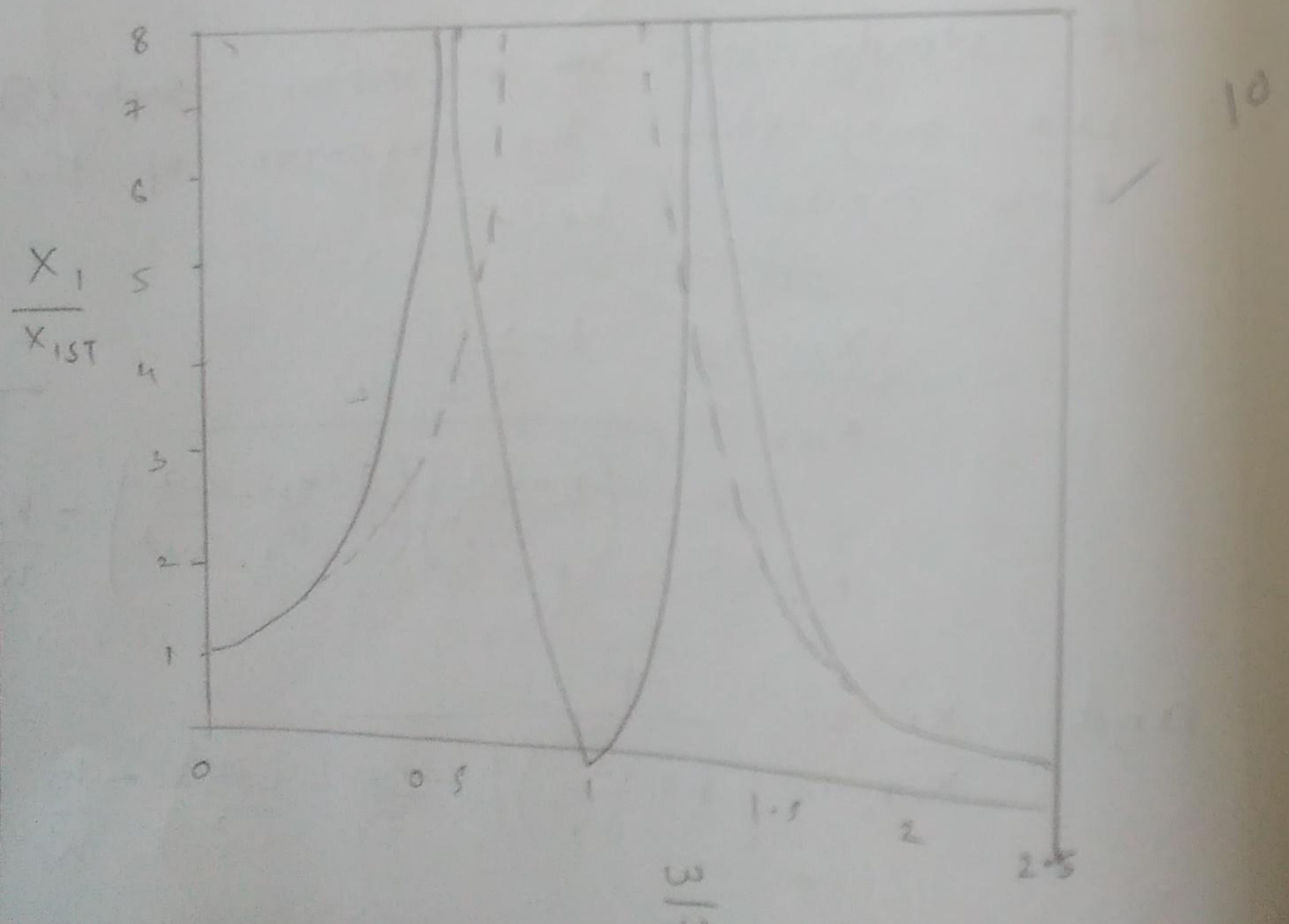


Fig-2 response of the undamped main mass with no absorber and with an undamped conventional vibration absorber. Fig 2 shows clearly shows two peaks corresponding to the two degrees of freedom and a null.

3Ans.) Circularly symmetric structures

Consider a circular plate or disk, as shown in fig 1 vibrating in a 2 nodal diameter pattern. This may represent a spinning turbine wheel around which blades may be mounted. When a forcing function acts on the plate in such a way that it matches natural mode both in space and time, then vibratory amplitudes reach a peak at the natural frequency (ω_1) corresponding to the $n=2$ mode. However, in the presence of any asymmetry, the well defined $n=2$ mode splits into two distinct but closely spaced frequencies (ω_{-2} and ω_{+2}). The spread between these frequencies depends upon the level of mistuning. The split mode

are orthogonal to each other
node line

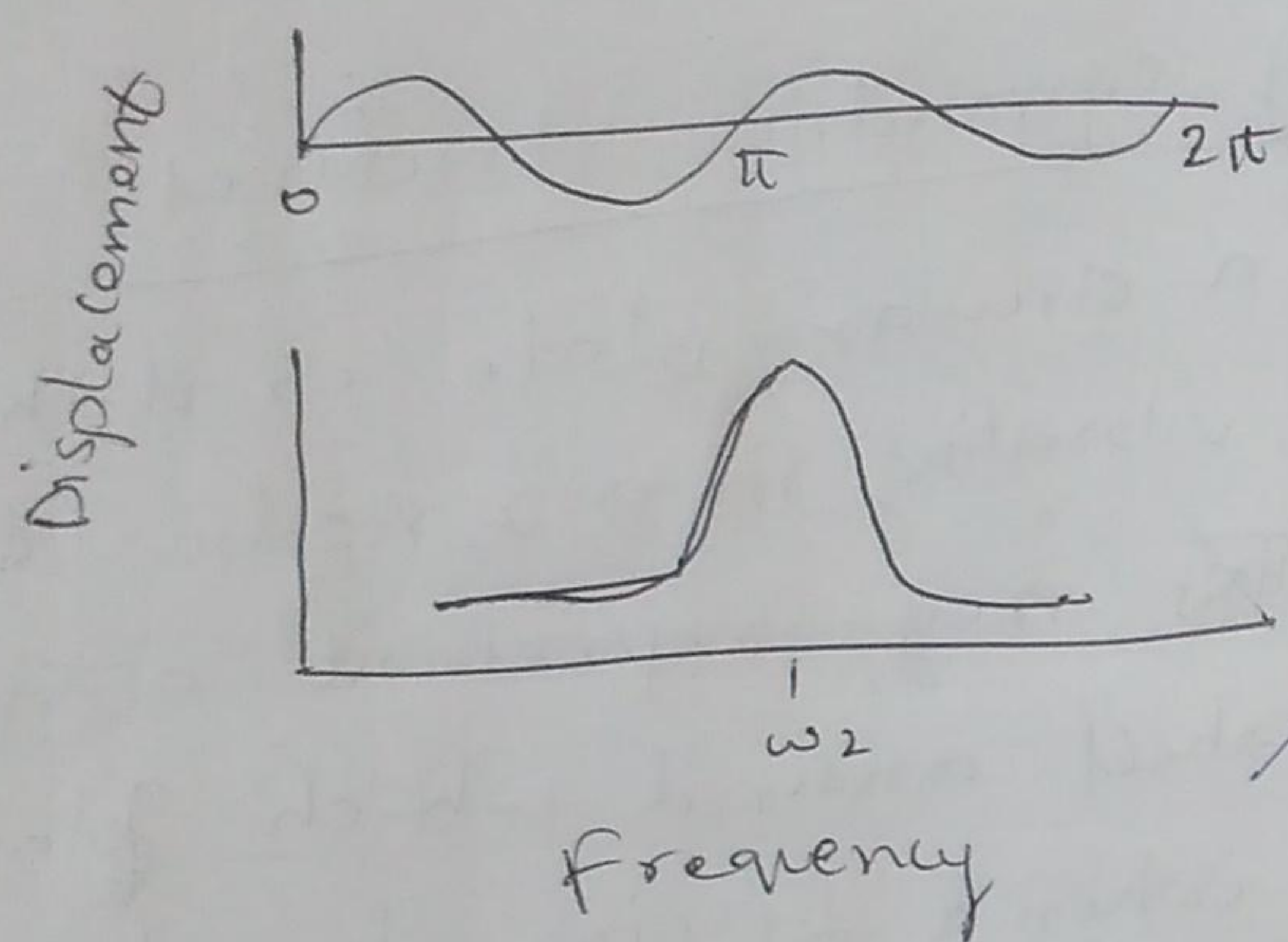
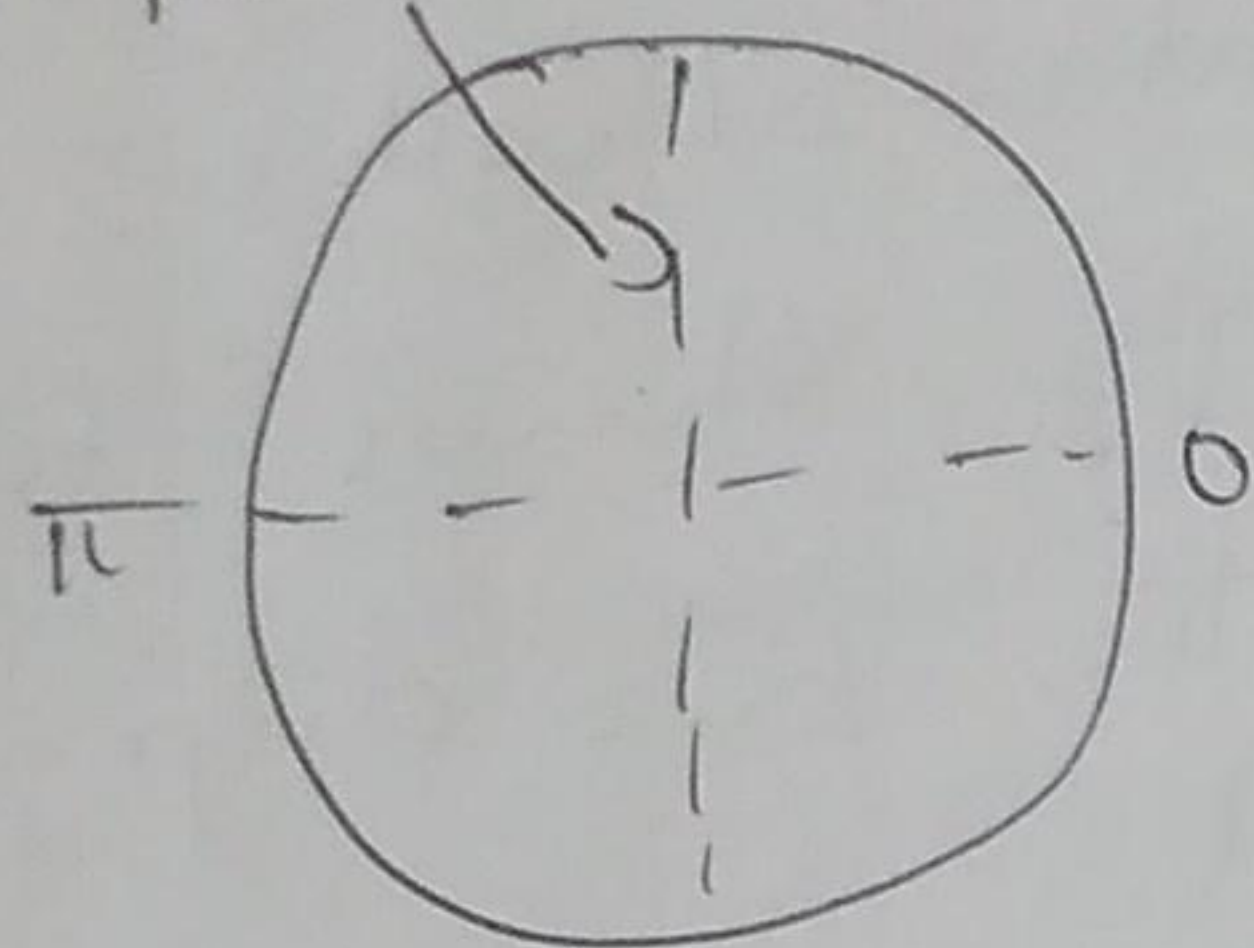


Fig ① A circular symmetric plate vibrating in a 2 nodal diameter pattern and share the same simple sinusoidal circumferential distribution but are displaced by a quarter wave. Thus every such pair represents legitimate modes of the system and contributes to the dynamic response of the structure, much as the individual nodes of the corresponding tuned system would have.

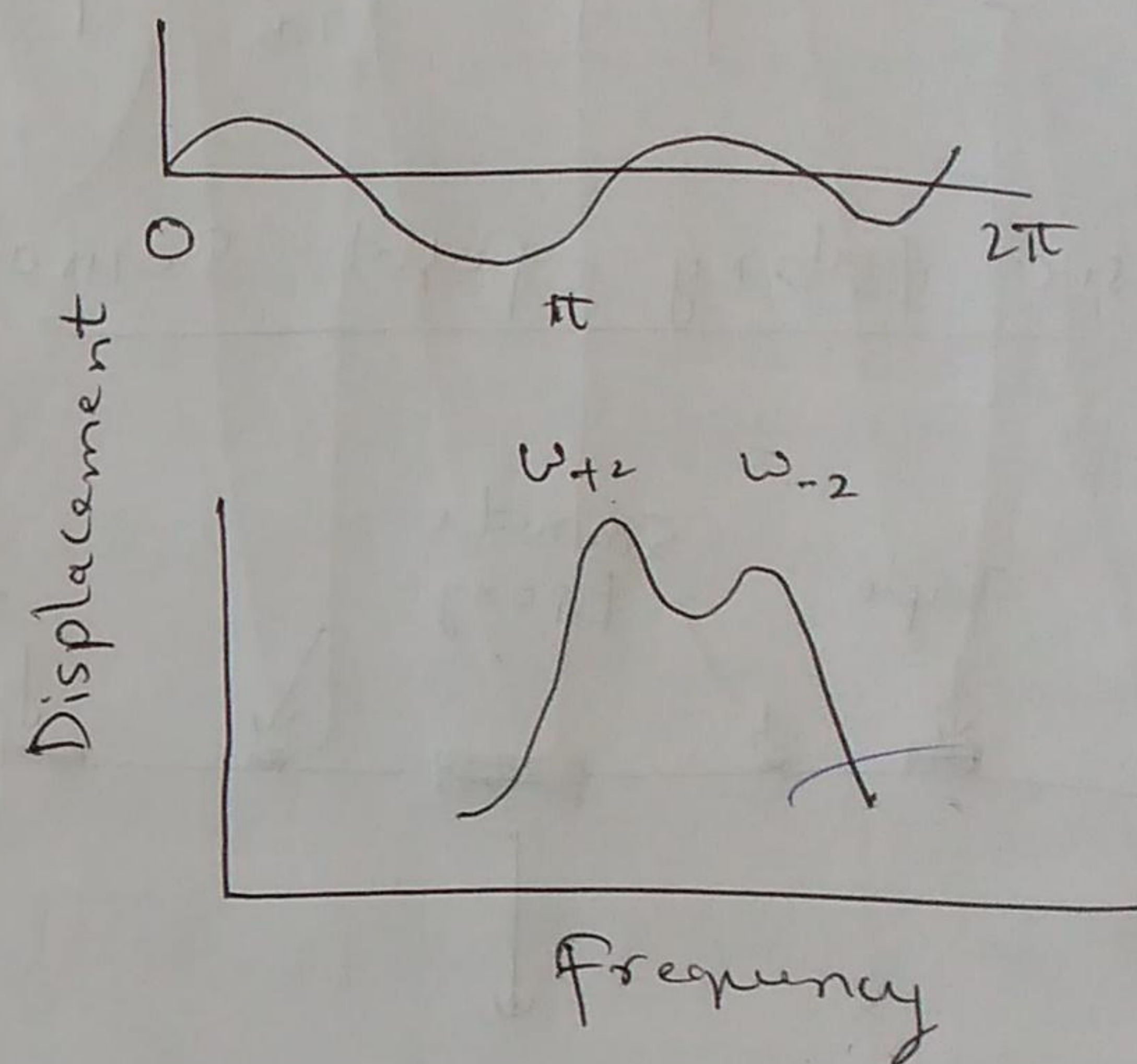
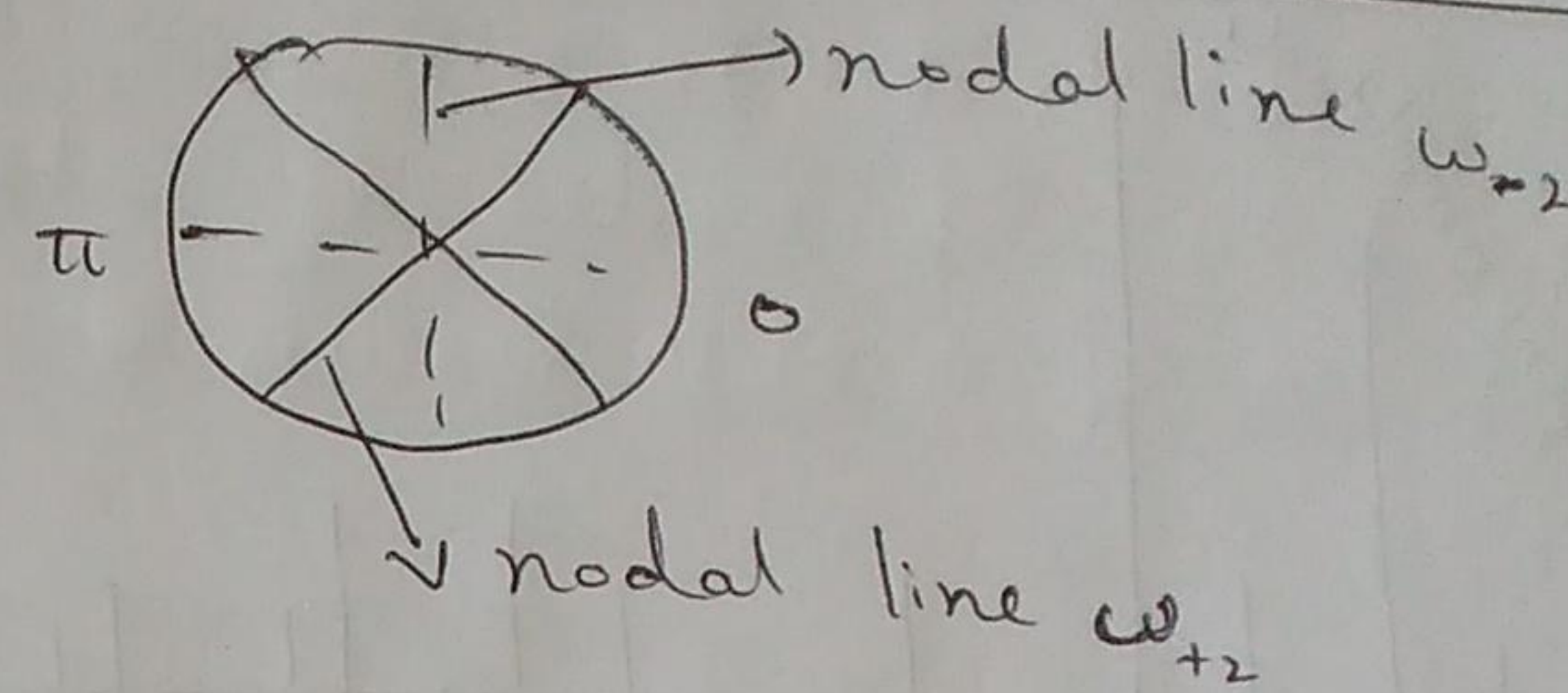


Fig ② An asymmetric circular plate vibrating in a 2 nodal diameter pattern.

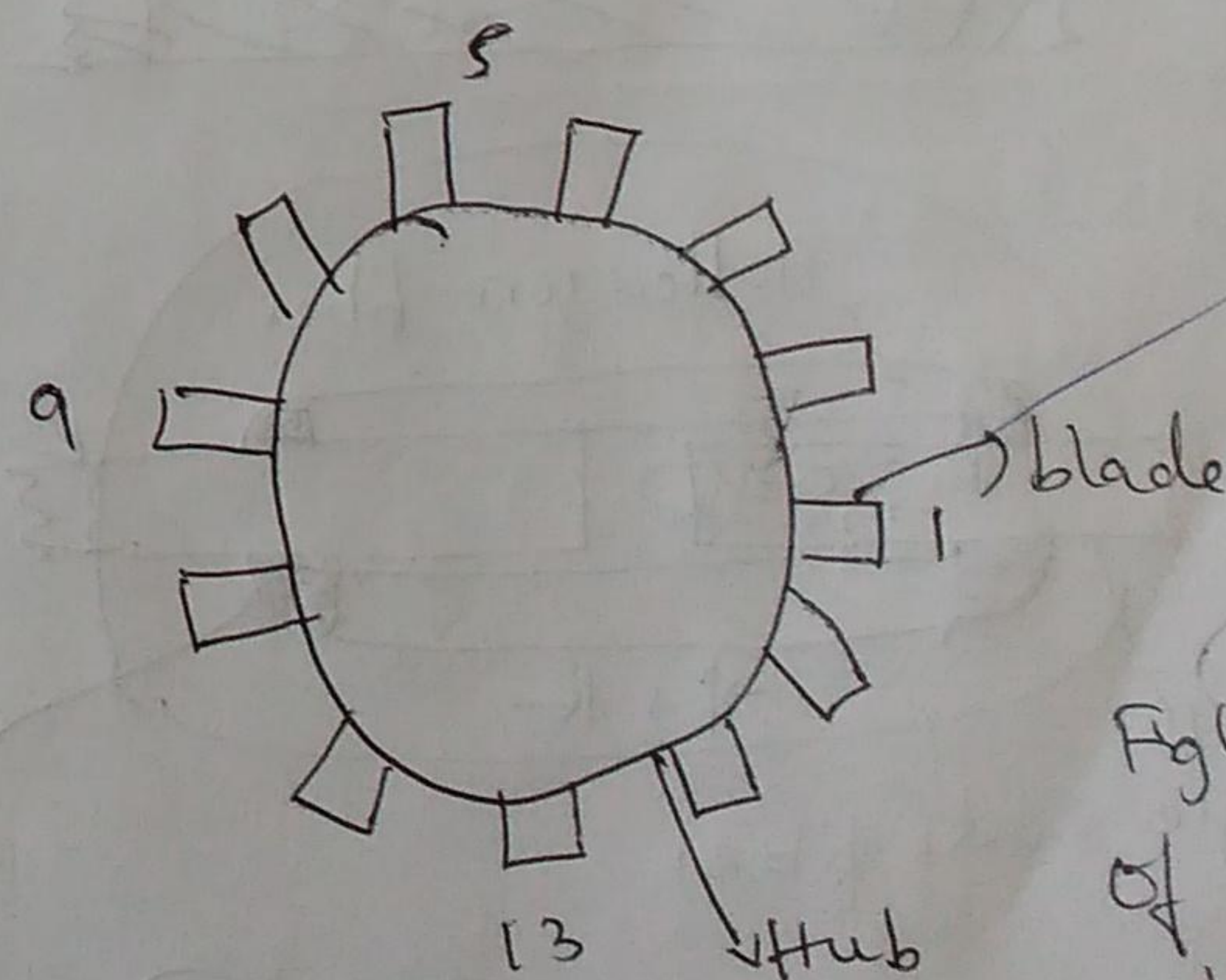
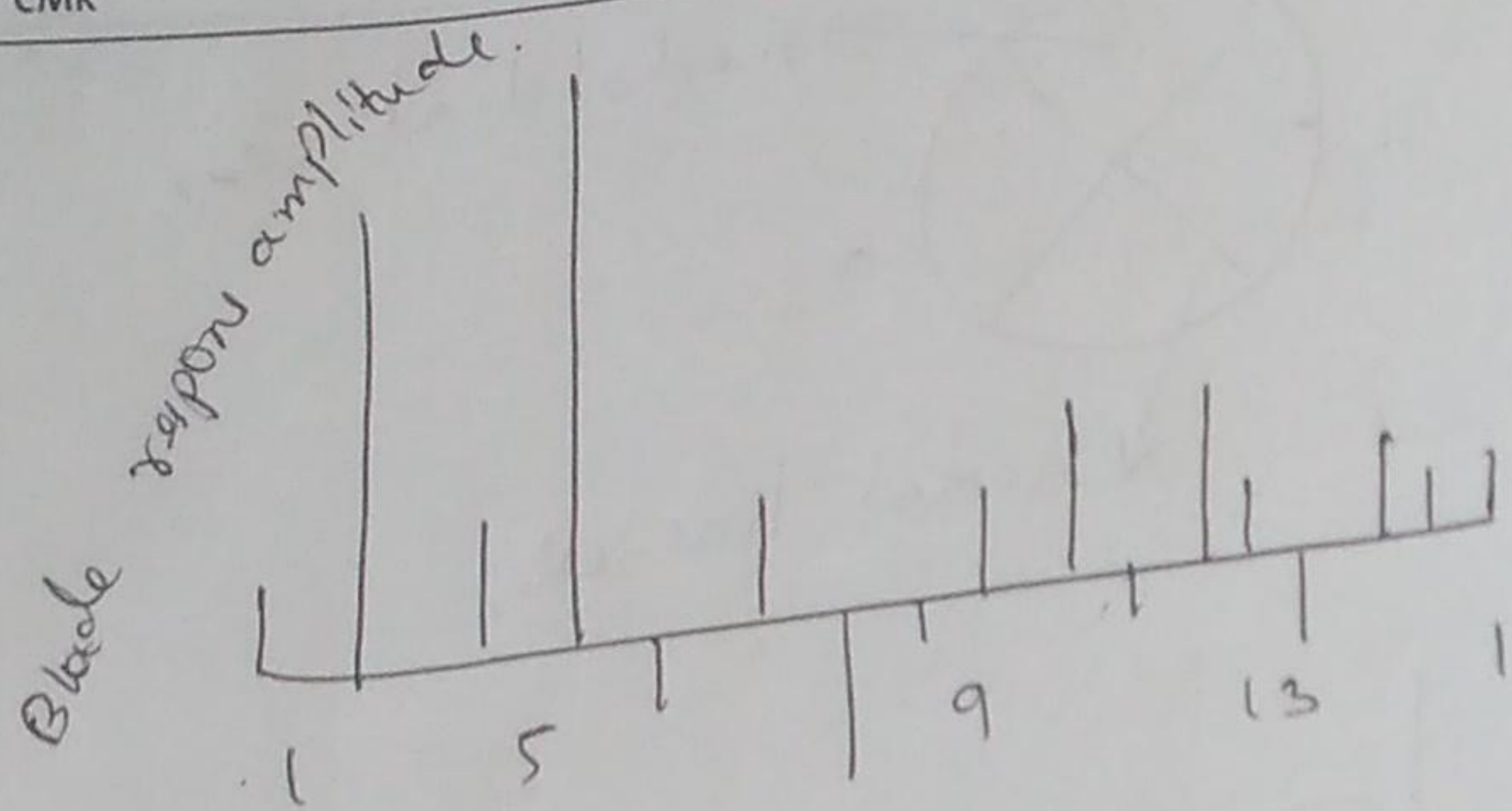


Fig ③ Amplitude of vibration of mistuned blade disk assembly



10

UAS-)

Extrinsic Fabry-Pérot Sensors

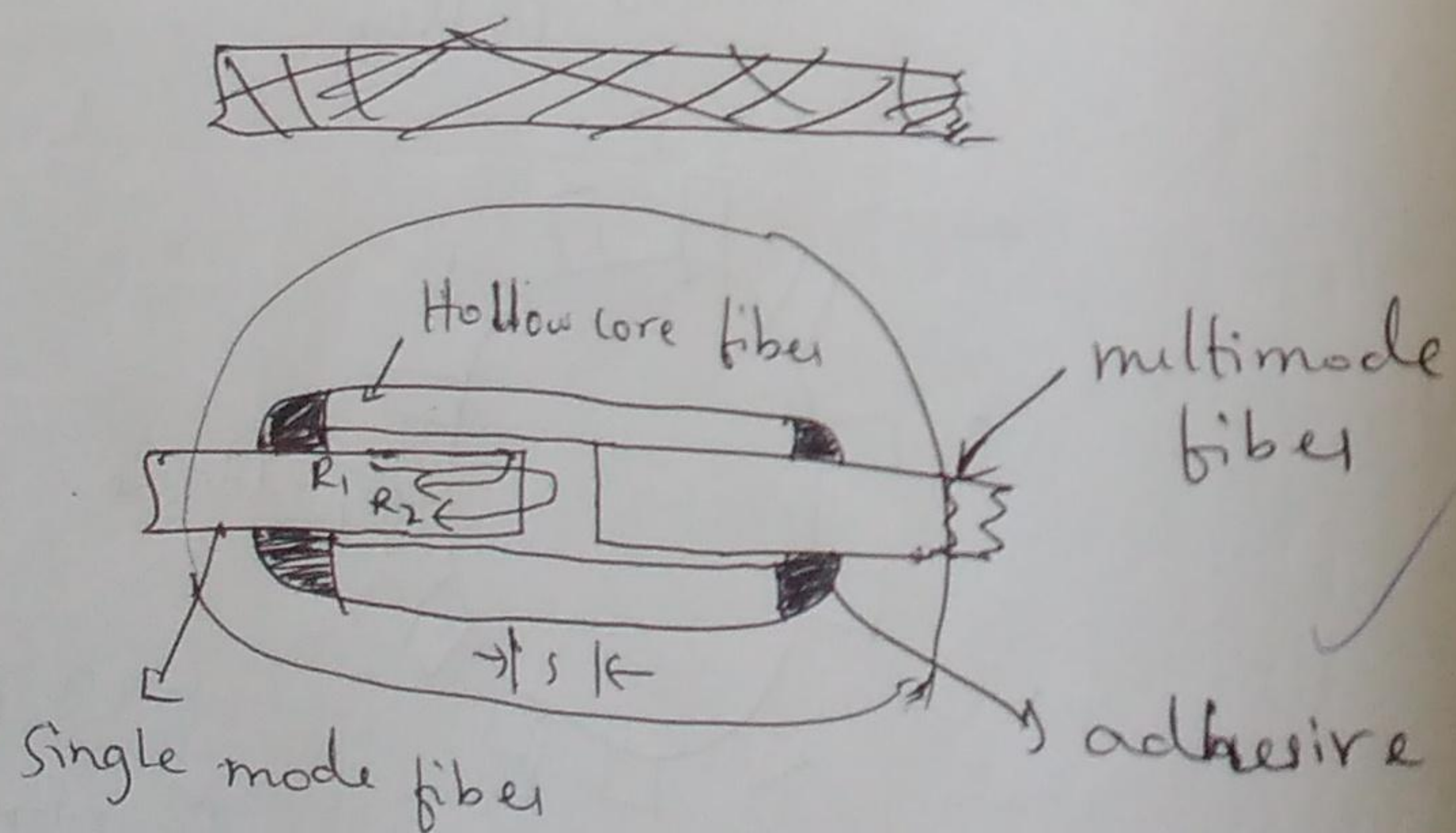
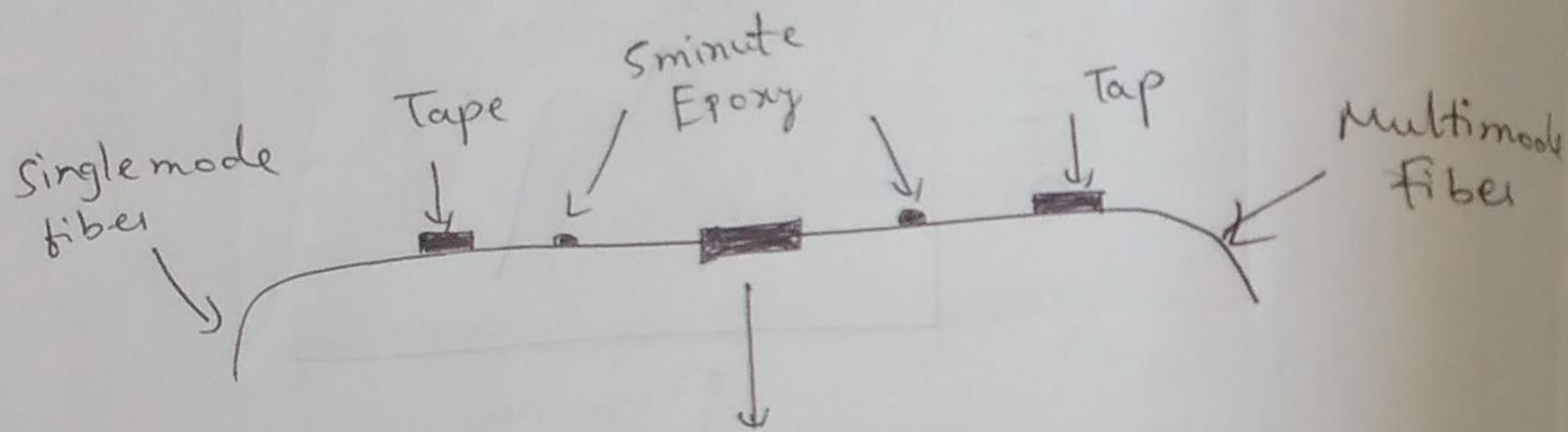
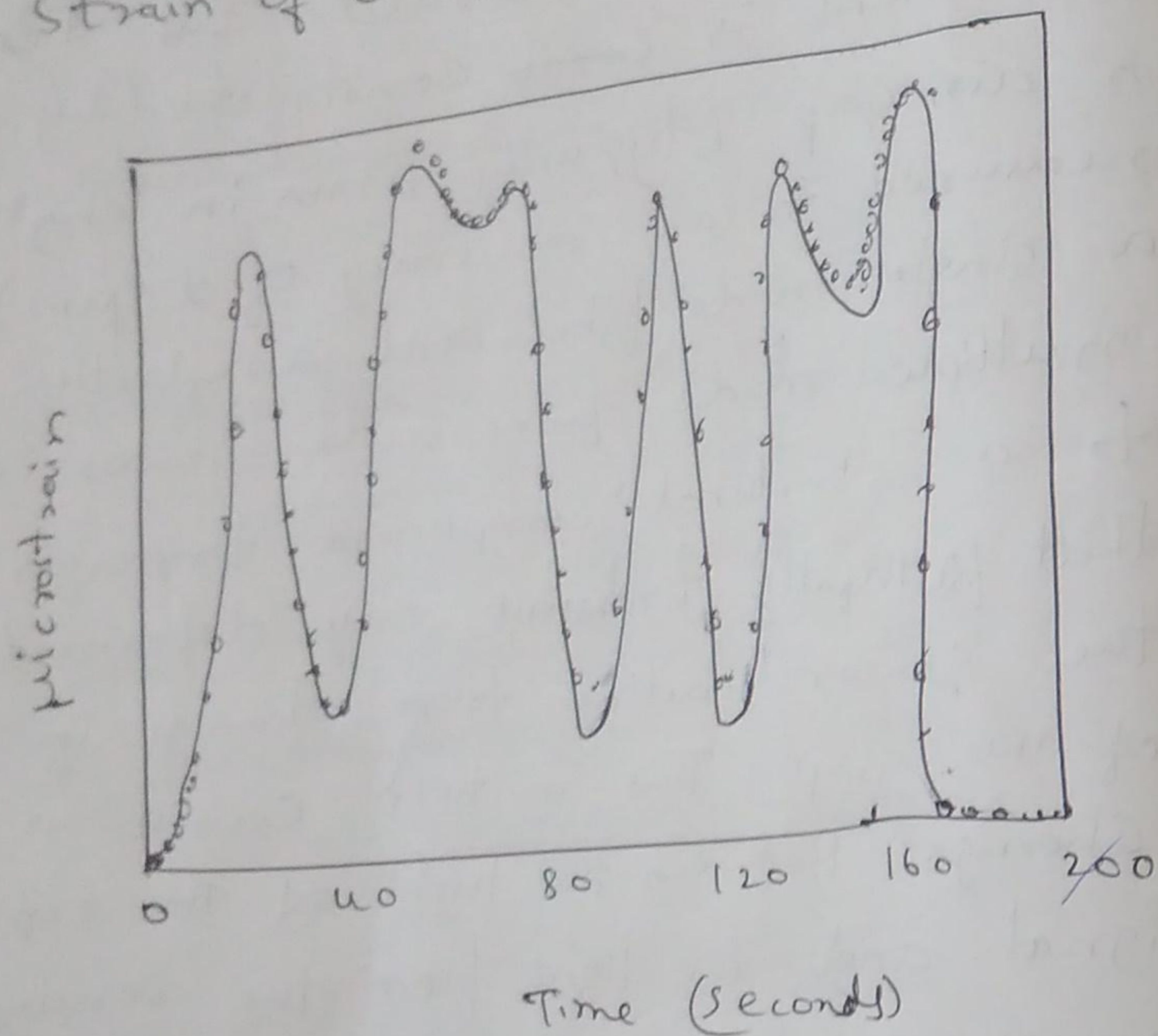


Fig ①

The basic principle that governs the operation of Fabry-Pérot sensors can be understood with reference figure ①. The instrument is constructed by providing an air gap (typically $10\mu\text{m}$ in length) or measured to an accuracy of $\pm 1\mu\text{m}$) between a single mode fiber and a reflection from a multiple mode fibre. The sensor is attached to a structural component through adhesives that faithfully transmit any deformation of the sensor leading to a change in the length of the gap. This in turn, causes a phase change between the light of the reference signal and the light from the sensing side because of interference between the two rays. This phase change is a measure of motion at the gap location and serves as the basis for an accurate measurement of strain. An excellent example of the application of this instrument to measure strains in an F-15 airframe during log load fatigue test is described by Murphy et al. (1991). The results are depicted in fig ②. Fabry-Pérot sensor

demonstrated minimum detectable motion of order of 0.1mm and a strain of 0.01 μ strain.



SAns.) Bragg grating sensor

Bragg gratings are extremely close parallel lines "written" onto a small length (1 to 20mm) of the core of a fibre so as to create a systematic perturbation of the core's refractive index. This spatially periodic variation of the index of refraction acts as a filter by

reflecting certain wavelengths while allowing others to pass through. Typical gratings have resonant wavelengths of 400 to 2000nm when broadband light traveling in the optical fiber encounters the grating, light at a wavelength proportional to the Bragg spacing is reflected back. This may be observed as a gap in the spectrum of the transmitted light, or as a peak in that of the reflected light. When the grating is strained, the Bragg spacing and thus the reflected wavelength change accordingly; the resulting change in spectrum is used as a measure of strain. This is a robust sensing scheme because light wavelength does not change as it passes through other fibers or connectors.

7 Ans)

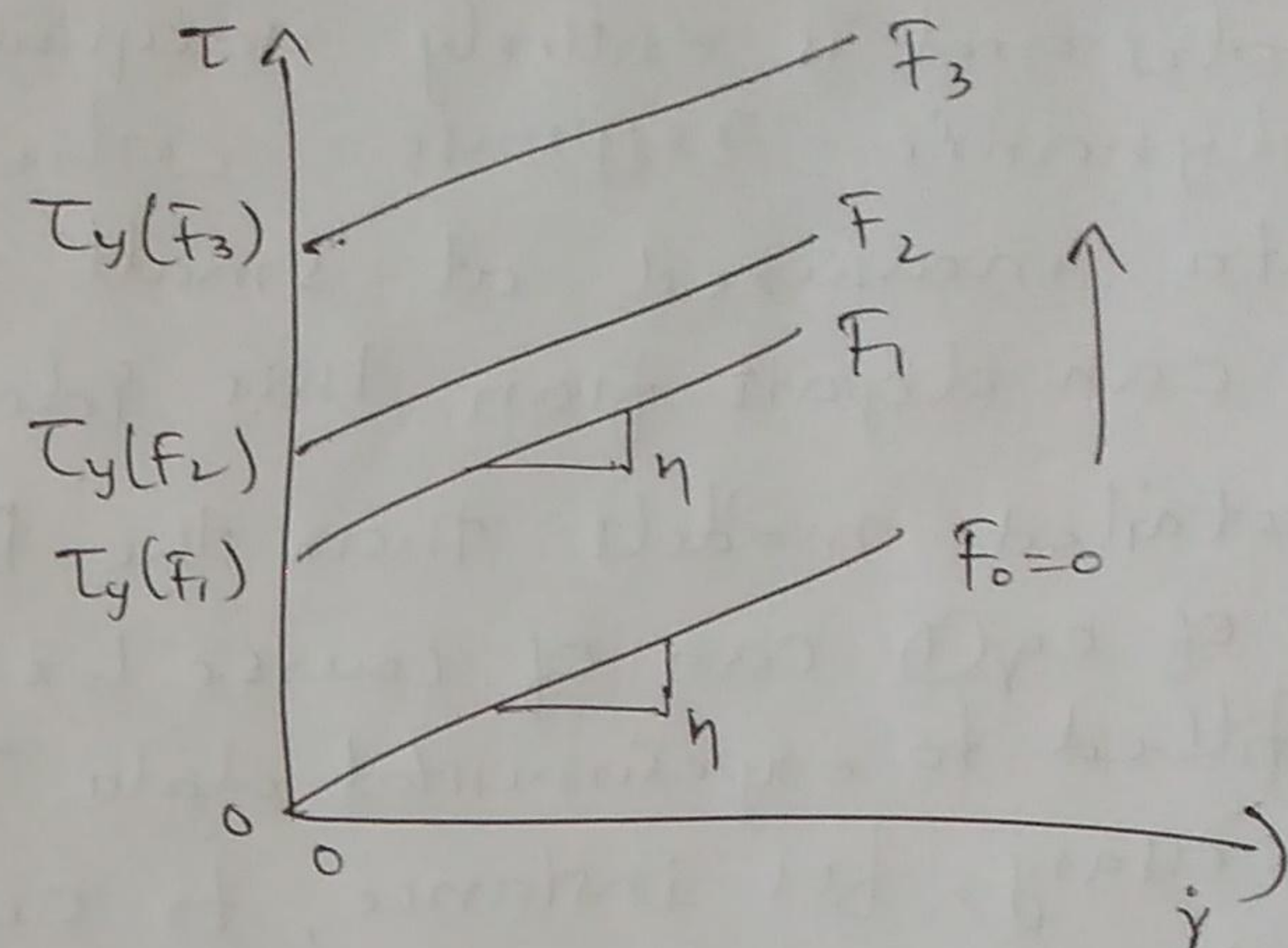
Bingham plastic model for ER/MR fluids

One might expect the ~~formation~~ formation of fibrils within an ER or MR fluid would increase the fluid's viscosity, changes little if at all. The effect of the fibrils is instead to produce a shear stress that is largely independent of the strain rate.

This is commonly referred to as the yield stress and denoted T_y . Adding this term to the Newtonian model results in the Bingham plastic model which has the stress-strain rate relation.

$$T_s = T_y(f) + \eta \dot{\gamma}$$

where in a given application f is the strength of the applied electric or magnetic field (i.e. E or H). The response predicted by this model is plotted in fig(1) which depicts the strong dependence of the yield stress on the field strength.



Fig(1) shear stress versus shear strain rate for the Bingham plastic material model.

This model or extensions of it that predict similar overall response, is by far the most popular for use in the design of devices that depend on the post yield shear resistance of an ER or MR fluid.

In practice, the dynamic viscosity is determined by a linear regression fit of a line to experimental data, and the intersection of this line with the shear stress axis is taken as the value of the yield stress T_y .

Although this is a good approximation at

strain rates and is entirely adequate for
 small dynamic response calculations,
 the data measured at small strain
 rates can depart from this idealization,
 more detailed models than the Bingham
 plastic of eq (1) can of course be devised
 and fitted to experimental data. It may
 be necessary, for instance, to capture the
 dynamics associated with the movement
 of the fluid, or to represent represent
 the finite compliance of the container.

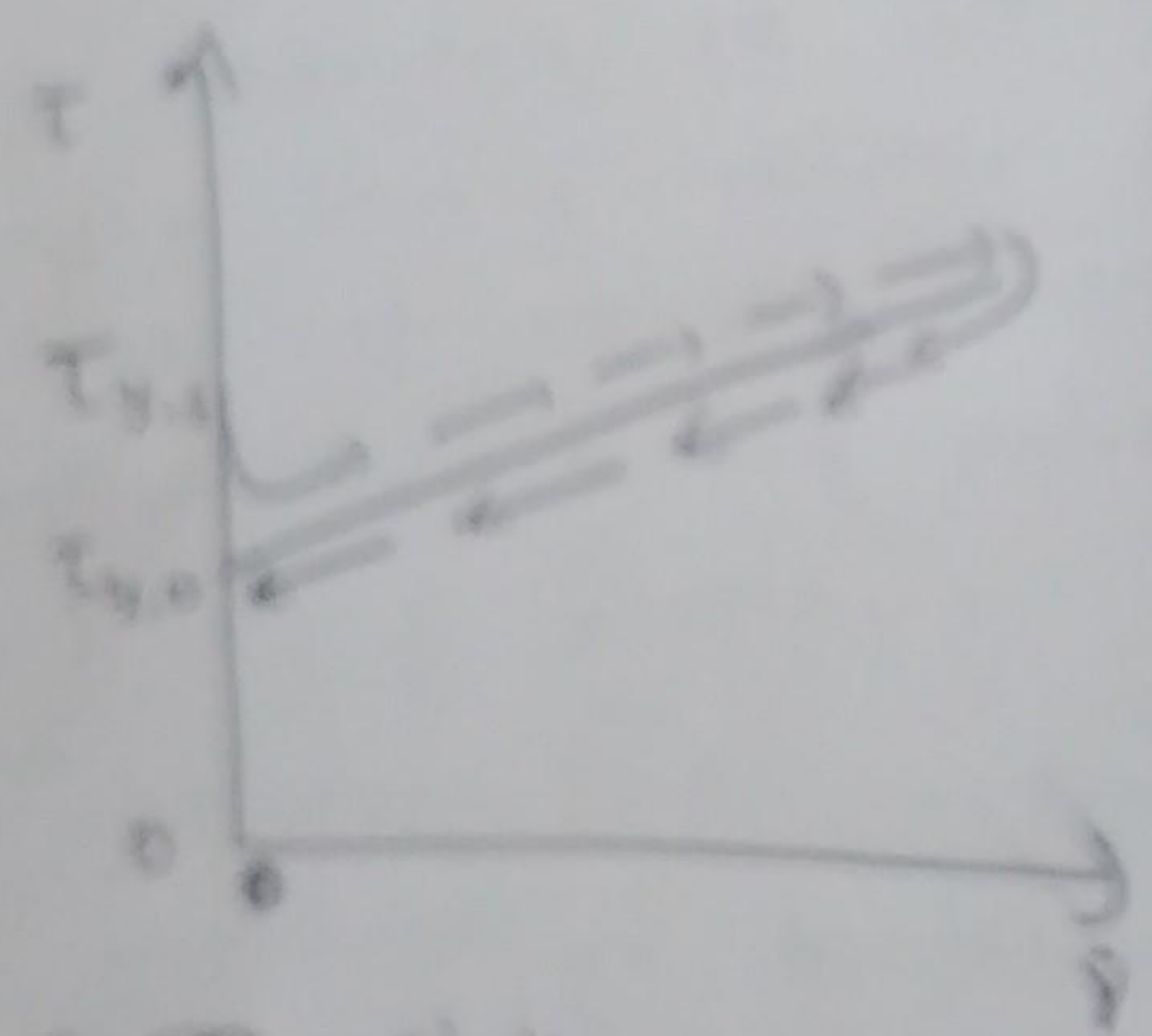


Fig. 2.1 Static and dynamic shear
 stress responses

An example of this extension of the
~~Bingham model to the case of~~
 fluid damping is presented below