

# CMR Institute of Technology, Bangalore DEPARTMENT OF ELECTRICAL & ELECTRONICS ENGINEERING II - INTERNAL ASSESSMENT

Semester: 8-CBCS 2017 Date: 19 Jun 2021
Subject: POWER SYSTEM OPERATION & CONTROL (17EE81) Time: 09:00 AM - 10:30 AM

Faculty: Ms Sanitha Max Marks: 50

Instructions to Students :  Answer any 5 questions. Each question carries 10 marks.  Answer any 5 question(s)											
						Q.No		Marks	СО	PO	BT/CL
						1	Explain the functions of various components in a steam turbine governing system with necessary diagram.	10	CO3	PO2	L2
2	Derive the transfer function and the block diagram of complete ALFC loop.	10	CO3	PO3	L3						
3	Draw the block diagram of two area system with necessary equations.	10	CO3	PO3	L3						
4	What are tie line oscillations? What determines the frequency of these oscillations?	10	CO3	PO1	L1						
5	Two control areas of capacity 500 MW and $10000  \text{MW}$ are connected through a tieline. The parameters of each area on its own capacity are R= $1  \text{Hz/pu}  \text{MW}$ and D= $0.02  \text{pu}  \text{MW/Hz}$ . There is an increase of $200  \text{MW}$ of load in area $2.0  \text{Determine}$ the steady state frequency deviation and change in tie line flow.	10	CO3	PO3	L3						
6	Two control areas are connected via a tie line with the following characteristics.  Area 1: R1=1 %, D1=0.8 pu , base MVA 1000  Area 2: R2=2 %, D2=1.0 pu , base MVA 1000  A load increase of 100 MW occurs in area 1.What is the new steady state frequency and the change in tie line flow if the nominal frequency is 50 Hz. Repeat if the load changes in area 2	10	CO3	PO3	L4						

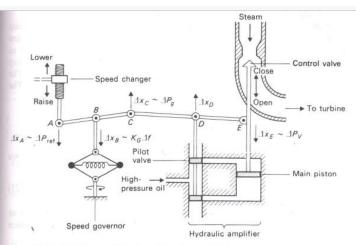


Figure 9-7 Simplified functional diagram of the primary ALFC loop.

megawatt increment  $\Delta P_V$ . This flow increase translates into a turbine power increment  $\Delta P_T$  in the turbine (not shown in the figure).

Very large mechanical forces are needed to position the main valve (or gate) against the high steam (or water) pressure, and these forces are obtained via several stages of hydraulic amplifiers. In our simplified version we show only one stage. The input to this amplifier is the position  $x_D$  of the pilot valve. The output is the position  $x_E$  of the main piston. Because the high-pressure hydraulic fluid exerts only a slight differential force on the pilot valve, the force amplification is very great.

The position of the pilot valve can be affected via the linkage system in three ways:

- 1. Directly, by the speed changer. A small downward movement of the linkage point A corresponds to an increase  $\Delta P_{\rm ref}$  in the reference power setting.
- 2. Indirectly, via feedback, due to position changes of the main piston.
- 3. *Indirectly*, via feedback, due to position changes of linkage point *B* resulting from speed changes.

It should prove a useful exercise for the reader to find, qualitatively, the workings of the mechanism. For example, give a "raise" command to the speed changer and prove that this indeed results in an increase in turbine output. Prove also that a speed drop will give the same effect.

measured in minimileters but in our analysis we shall rather express them a power increments expressed in megawatts or per-unit megawatts as the case mabe. The movements are assumed positive in the directions of the arrows. The governor output command  $\Delta P_g$  is measured by the position change  $\Delta x_C$ . The governor has two inputs:

- 1. Changes  $\Delta P_{ref}$  in the reference power setting
- 2. Changes  $\Delta f$  in the speed of frequency of the generator, as measured by  $\Delta x_B$

An increase in  $\Delta P_g$  results from an increase in  $\Delta P_{ref}$  and a decrease in  $\Delta f$ . W thus can write for small increments

$$\Delta P_g = \Delta P_{\text{ref}} - \frac{1}{R} \, \Delta f \qquad \text{MW}$$
 (9-21)

The constant R has dimension hertz per megawatt, and is referred to a regulation or droop. (For numerical values see Example 9-2 below.) Laplace transformation of Eq. (9-21) yields

$$\Delta P_g(s) = \Delta P_{\text{ref}}(s) - \frac{1}{R} \Delta f(s)$$
 (9-22)

Using well-known block diagram symbols we have represented the governo as shown in Fig. 9-8.

### 9-3-2 Hydraulic Valve Actuator

The input position  $\Delta x_D$  of the valve actuator increases as a result of an increased command  $\Delta P_g$  but decreases due to increased valve output,  $\Delta P_V$ . Equal in creases in both  $\Delta P_g$  and  $\Delta P_V$  should result in  $\Delta x_D = 0$ . We can thus write

$$\Delta x_D = \Delta P_a - \Delta P_V \qquad \text{MW} \tag{9-23}$$

For small changes  $\Delta x_D$  the oil flow into the hydraulic motor is proportional to position  $\Delta x_D$  of the pilot valve. Thus we obtain the following relationship to the position of the main piston:

$$\Delta P_V = k_H \int \Delta x_D \, dt \tag{9-24}$$

The positive constant  $k_H$  depends upon orifice and cylinder geometries and fluid pressure.

Upon Laplace transformation of the last two equations and upon elimination of  $\Delta x_D$  we obtain the actuator transfer function

$$G_H(s) = \frac{\Delta P_V}{\Delta P_a} = \frac{1}{1 + sT_H} \tag{9-25}$$

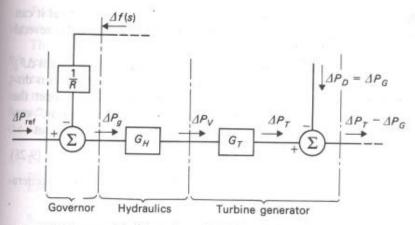


Figure 9-8 Linear model of the primary ALFC loop (minus the power system response).

where the hydraulic time constant

$$T_H = \frac{1}{k_H}$$

typically assumes values around 0.1 s.

The hydraulic valve actuator has been represented by the transfer function  $G_H(s)$  in Fig. 9-8.

## 9-3-3 Turbine-Generator Response

In normal steady state and via the mechanism described in Sec. 4-9 the turbine power  $P_T$  keeps balance with the electromechanical air-gap power  $P_G$  resulting in zero acceleration and a constant speed or frequency. Perturbations  $\Delta P_T$  and  $\Delta P_G$  in these powers will upset the above balance. If the difference power,  $\Delta P_T - \Delta P_G$ , is positive the turbine generator unit will accelerate; if negative it will decelerate.

The turbine power increment  $\Delta P_T$  depends entirely upon the valve power increment  $\Delta P_V$  and the response characteristics of the turbine. Different turbine types vary widely in this regard. It is possible to express the turbine dynamics in terms of a turbine transfer function

$$G_T = \frac{\Delta P_T}{\Delta P_V} \tag{9-26}$$

In App. D we have derived  $G_T$  for the most common turbine types. A so-called nonreheat steam turbine has the simplest transfer function, consisting of a single time constant, i.e.,

$$G_T = \frac{1}{1 + sT_T} {9-27}$$

Figure 6.18 Model of governor.

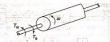


Figure 6.19 Torques acting on a generator.

Figure 6.18(b) and (c) shows how the block diagram can be reduced.  $T_G$  is the governor time considered be seen that it depends on the speed regulation R and on the gain of the hydraulic amplifier,  $K_G$ .

#### 6.6.2 Generator Model

There are two torques acting on a generator: the theft torque (due to the prime mover) and the electromagnetic torque, neglecting losses. The shaft torque tends to accelerate the generator in the positive direction of rotation and the electromagnetic torque in the negative direction, as shown in Fig. 6.19.

The total accelerating torque is given by

$$T_{\bullet} = T_{\bullet} - T_{\bullet}^{\text{equation}}$$
 of the many in consists as an  $(6.6)$ 

From Newton's laws of motion, we have for rotatory motion

$$I\alpha = T$$
 (6.7)

where I is the moment of inertia, where I is the information of inertia,  $\alpha$  is the angular acceleration, T is the net torque.

Equation (6.7) can be written as

$$I\frac{d^{2}\theta_{m}}{dt^{2}} = T_{m} - T_{e}$$
 (6.8)

where  $\theta_m$ , the rotor angle, is now converted into an angle, measured with respect to a synchronously rotating

 $\theta_m$ , the rotto angular production is a case such that  $\theta_m = \theta_m - \omega_m t \quad \text{as the product is a first production}$ where  $\omega_m$  is the synchronous speed in rad/s,  $\delta_m = \theta_m - \omega_m t$  where  $\omega_m$  is the angular displacement in rad. From Eq. (6.9) (6.9)

$$=\frac{d^2\theta_{\rm m}}{L^2} \tag{6.10}$$

Substituting into Eq. (6.8), we get

$$\frac{d^2\delta_m}{dt} = T_m - T_m N - m$$

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 $I\frac{d^2\delta_m}{dt^2}=T_m-T_e\ N-m$  Multiplying both sides by the angular velocity  $\omega_m$  , we get

$$\omega_{\rm m} I \frac{d^2 \delta_{\rm m}}{dt^2} = \omega_{\rm m} (T_{\rm m} - T_{\rm e}) N - m$$
 (6.12)

nere  $\omega_m I = M$  = angular momentum or inertia constant  $\omega_m I = M$  = mechanical power input at the shaft minus rotational losses  $\omega_m I_c = P_c$  = electrical power output minus losses

We can write Eq. (6.12) as

$$M\frac{d^2\delta_{\rm m}}{dt^2} = P_{\rm m} - P_{\rm e} = P_{\rm a} \ \mathbb{W} \tag{6.13}$$

M depends on the speed  $\omega_n$ . However, since the deviation in speed is limited, M can be assumed to be a constant. The value of M varies over a wide range depending on the rating and type of the generator. Hence, another constant H is used to specify the energy stored in the machine.

H is also called *inertia constant*. It lies in a narrow range for different machines. M and H are related as follows

$$M = \frac{2GH}{\omega_{sm}} MJ \cdot s/mech rad$$
 (6.)

where G = MVA rating of machine. In pu, M = 2H.

Now, Eq. (6.13) can be written as

$$\frac{2H}{\omega_{um}}\frac{d^2\delta_m}{dt^2} = \frac{P_m - P_e}{G} \tag{6.16}$$

We can express both  $\delta_{\mathrm{m}}$  and  $\omega_{\mathrm{m}}$  in terms of electrical radians to get  $\frac{2H}{\omega_{\rm a}}\frac{d^2\delta_{\rm m}}{dt^2} = P_{\rm m} - P_{\rm e} = P_{\rm a} \text{ pu}$ (6.17) Here,  $P_m$  = per unit mechanical power  $[P_m$  in MW/G].  $\omega = \text{synchronous speed in electrical rad/s}$   $= \frac{P}{2} \omega_m$  = specifical rad/s

$$\frac{p}{2}\omega_{\rm im}$$

Equation (6.17) is called the swing

$$P_c$$
 = acceleration power

Equation (6.17) is called the *swing equation*. We can linearize Eq. (6.17) to get

$$\frac{2H d^2 \Delta \delta}{\partial - dt^2} = \Delta P_c \Delta P_c$$
(6.18)

We express the speed deviation also in pu to get

Taking the Laplace transform we get

$$\frac{d\Delta\omega}{dt} = \frac{1}{2H} (\Delta P_{\rm m} - \Delta P_{\rm c})$$
(6.19)

 $\Delta\omega(s) = \frac{1}{2Hs} \left( \Delta P_{n}(s) - \Delta P_{e}(s) \right)$ Eq.(6.20) can be written in a block diagram form, as shown in Fig. 6.20.



Figure 6.20 Model of generator.

#### 6.6.3 Load Model

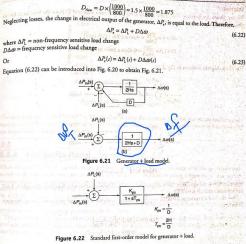
We have seen the load models in detail in Chapter 2. Some loads exhibit variation in active power drawn with respect to frequency variations. This relationship is given by

ationship is given by 
$$\Delta P_{\text{Lfeq}} = D\Delta \omega$$

or D = load damping constant

$$D = \frac{\Delta P_{\text{Lineq}}}{\Delta A_{\text{Lineq}}} \tag{6.21}$$

where  $\Delta P_{Heq}$  = frequency-dependent load.



In Fig. 6.21, 2H can be replaced by M, where both are in pu. The transfer function of Fig. 6.21 can be write ten in the form of a standard first-order transfer function  $\frac{K_{ps}}{1+T_{ps}}$ , as shown in Fig. 6.22, where

 $K_{ps}$  is the power system gain and  $T_{ps}$  is the power system time const

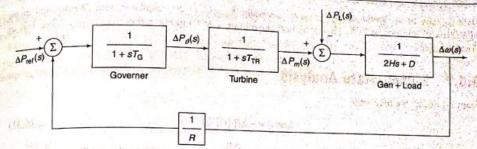


Figure 6.29 Block diagram of complete ALFC.

We are now interested in deriving the effect of change in load on the frequency without change in the reference set point. By changing the reference set point, we can set the system to give specified frequency at any load point as explained in Fig. 6.7(b). This is a secondary control to be discussed later. Here, we assume  $P_{\rm ref}$  is kept at a constant value so that  $\Delta P_{\rm ref} = 0$ . We now find the transfer function  $\frac{\Delta \omega(s)}{-\Delta P_{\rm L}(s)}$ . From the block diagram of Fig. 6.29,

$$\Delta\omega(s) = -\Delta P_{L}(s) \left[ \frac{\frac{1}{2Hs + D}}{1 + \frac{1}{R} \left( \frac{1}{2Hs + D} \right) \left( \frac{1}{1 + sT_{G}} \right) \left( \frac{1}{1 + sT_{TR}} \right)} \right]$$

$$= -\Delta P_{L}(s) \left[ \frac{(1 + sT_{G})(1 + sT_{TR})}{(2Hs + D)(1 + sT_{G})(1 + sT_{TR}) + \frac{1}{R}} \right]$$
(6.29a)

The transfer function is given by

(EE.0)

$$T(s) = \left[ \frac{(1 + sT_{\rm G})(1 + sT_{\rm TR})}{(2Hs + D)(1 + sT_{\rm G})(1 + sT_{\rm TR}) + \frac{1}{R}} \right]$$

An alternate expression for the transfer function commonly used is

$$T(s) = \begin{bmatrix} \frac{K_{ps}}{1+sT_{ps}} \\ \frac{K_{ps}}{1+\left(\frac{K_{ps}}{1+sT_{ps}}\right)\left(\frac{1}{1+sT_{T}}\right)\left(\frac{1}{1+sT_{TR}}\right)\frac{1}{R}} \end{bmatrix}$$

$$= \begin{bmatrix} \frac{K_{ps}(1+sT_{G})(1+sT_{TR})}{(1+sT_{ps})(1+sT_{G})(1+sT_{TR}) + \frac{K_{ps}}{R}} \end{bmatrix}$$
(6.30)

# 7.2 Tie-Line Control with Primary Speed Control

Let us consider a two-area system as shown in Fig. 7.1. Let us take the positive power flow to be  $P_{12}$ , to be the poon the tie-line from area 1 to area 2 is

$$\delta_{12} = \frac{E_1 E_2}{X_{12}} \sin(\delta_1 - \delta_2)$$
 (7.1)

$$X_{12} = X_1 + X_{tie} + X_2$$

Equation (7.1) can be linearized about an initial operating point  $\delta_1 = \delta_{10}$  and  $\delta_2 = \delta_{20}$  as

$$\Delta P_{12} = \frac{E_1 E_2}{X_{12}} \cos(\delta_{10} - \delta_{20}) \Delta \delta_{12}$$
 (7.2)

$$\Delta \delta_{12} = \Delta \delta_1 - \Delta \delta_2$$
 is retaining this is the stress and less (7.3)

Let 
$$T = \frac{E_1 E_2}{X_{12}} \cos(\delta_{10} - \delta_{20}) = P_{\text{max}} \cos(\delta_{10} - \delta_{20})$$
 (7.4)

where T is called the synchronizing torque coefficient (often designated as P)

Substituting Eq. (7.4) into Eq. (7.2), we get

$$\Delta P_{12} = T \left( \Delta \delta_1 - \Delta \delta_2 \right) \tag{7.5}$$

The block diagram representation of the two-area system with only primary control is shown in Fig. 7.2. A positive  $\Delta P_{12}$  means an increase in power flow from area 1 to 2. This is equivalent to a load increase in area 1 and/or decreasing load in area 2. Therefore, the feedback from  $\Delta P_{12}$  has a negative sign for area 1 and positive sign for area 2. We will now see how the system behaves for a change in the load.

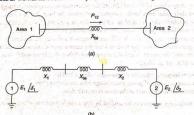


Figure 7.1 (a) Two

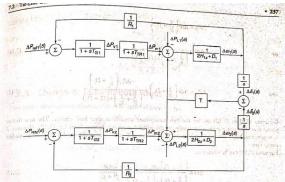


Figure 7.2 Two-area system with primary loop

#### 7.2.1 Change of Load in Area 1

Consider a load change of  $\Delta P_{\rm Li}$  in area 1. When the system has reached a steady state, both areas will have same steady-state frequency deviations. Therefore,

(or  $\Delta f - \Delta f_1 = \Delta f_2$ . Remember that in pu both  $\Delta f$  and  $\Delta w$  are the same). If we assume that the mechanical powers are constant (which means  $\Delta P_{eef}$  is constant), the tie-line and rotating masses exhibit damped oscillations called synchronizing oscillations. For area 1, we can write

called synchronizing oscillations. For area 1, we can write 
$$\Delta P_{\rm enl} - \Delta P_{11} = \Delta \omega D_{11} = \Delta \omega D_{12} = \Delta \omega D_{13} = \Delta \omega D_{14} = \Delta D_{14} = \Delta \omega D_{14} = \Delta \omega D_{14} = \Delta \omega D_{14} = \Delta D_{$$

For area 2, we have 
$$\Delta P_{ns} + \Delta P_{12} = \Delta \omega D_2 \tag{7.8}$$

$$\Delta P_{\rm ml} = \frac{c_{\rm kin}}{R_{\rm i}}$$
(7.10)

$$-\Delta P_{12} - \Delta P_{11} = \Delta \omega \left( \frac{1}{R_i} + D_i \right)$$

4

#### **Tie-Line Oscillations** 7.6

We saw in the previous section that the system state matrix is a 9 × 9 matrix for a two-area system. With some simplification, we can get a fairly good idea of the effect of the system parameters on the dynamic response. Let us make the following assumptions:

- 1. Neglect turbine and governor time constants.
- 2. Neglect damping constants D, and D,.
- 3. Both areas are identical.

With these assumptions, the two area equations reduce to

$$\Delta P_{m1}(s) = \frac{-\Delta \omega_1(s)}{D} \qquad (7.70)$$

$$\Delta P_{m2}(s) = \frac{-\Delta \omega_2(s)}{R}$$

$$\Delta \omega_{1}(s) = \frac{1}{2H_{5}} [\Delta P_{m1}(s) - \Delta P_{L1}(s) - \Delta P_{12}(s)]$$

(7.71)

$$\Delta \omega_{1}(s) = \frac{1}{2H_{s}} [\Delta P_{m1}(s) - \Delta P_{L1}(s) - \Delta P_{12}(s)]$$

$$\Delta \omega_{1}(s) = \frac{1}{1 + \frac{1}{2RH_{0}}} \left[ \frac{-\Delta P_{1,1}(s) - \Delta P_{1,2}(s)}{2H_{0}} \right]$$
(7.73)

$$\Delta \omega_2(s) = \frac{1}{1 + \frac{1}{2RHs}} \left[ \frac{-\Delta P_{12}(s) + \Delta P_{12}(s)}{2Hs} \right]_{s=0}^{100} (7.74)$$

$$\Delta P_{12}(s) = \frac{T}{s} [\Delta \omega_1(s) - \Delta \omega_2(s)]$$

$$= \frac{T}{s} \frac{1}{\left(1 + \frac{1}{2RH_s}\right)} \left[ \frac{1}{2H_s} (\Delta P_{12}(s) - \Delta P_{13}(s) - 2 \Delta P_{12}(s)) \right]$$

$$\Delta P_{12}(s) \left[ 1 + \frac{2T}{2Hs^2 + \frac{s}{R}} \right] = \frac{T}{2Hs^2 + \frac{s}{R}} [\Delta P_{12}(s) - \Delta P_{L1}(s)]$$

$$\Delta P_{12}(s) = \frac{\frac{T}{2H}}{s^2 + \frac{s}{2RH} + \frac{T}{H}} \cdot \left[ \Delta P_{12}(s) - \Delta P_{11}(s) \right]. \tag{7.75}$$

The poles of the denominator determine the oscillations in  $\Delta P_{12}$ . We compare the denominator with the standard second-order characteristic equation.

$$s^2 + 2\alpha s + \omega_n^2 = s^2 + 2\xi \omega_n s + \omega_n^2$$
 (7.76)

We can see that

$$\alpha = \frac{1}{4RH} \tag{7.77}$$

And

or

$$\omega_{\rm n} = \sqrt{\frac{T}{H}} \text{ pu or } \sqrt{\frac{2\Pi f_0 T}{H}} \text{ rad/s}$$
 (7.78)

The damping is determined by the relative values of  $\alpha$  and  $\omega_n$  and the roots of Eq. (7.75). The roots of Eq. (7.75) are

$$s_1, s_2 = -\xi \omega_n \pm j \omega_n \sqrt{1 - \xi^2}$$

$$= -\alpha \pm j \omega_n$$

$$(7.79)$$

and

$$\omega_{\rm d} = \omega_{\rm n} \sqrt{1 - \xi^2} \tag{7.80}$$

 $\alpha$  is called the damping factor or damping constant,  $\omega_n$  is the natural undamped frequency of oscillations,  $\omega_n$  is called the damped or conditional frequency and  $\xi$  is called the damping ratio. We now have the following cases:

1. When  $\xi = 1$  or  $\alpha = \omega_n$  This condition is called "critical damping."

2. When  $\xi = 0$  or  $\alpha = 0$ , the roots are purely imaginary and we get purely sinusoidal oscillations for a step change in input. In Eq. (7.74), the input is  $\Delta P_{1,2}(s) - \Delta P_{1,1}(s)$ . Therefore, for a step change in the load, we would get sustained sinusoidal oscillations of tie-line power at a frequency of  $\omega_n$ . From Eq. (7.76), we can see that  $\alpha = 0$  when  $R = \infty$ . This means that there is no governor speed control.

3. When  $\alpha < \omega_a$ ,  $\xi < 1$ , we get a pair of complex conjugate roots. The system is "under damped" and we have oscillations in tie-line power flow which have a frequency  $\omega_d$  as in Eq. (7.79). The time constant of

the system is  $1/\alpha$ .

4. When  $\alpha > \omega_a$ , we have an "over damped" system. The roots are both real.

The above analysis is only approximate, but is helpful in knowing the effect of the choice of parameters on the stability of the system. If we consider the damping constants of the load, then  $\alpha$  is modified as

$$\alpha = \frac{1}{4H} \left[ D + \frac{1}{R} \right] \tag{7.81}$$

5

S) 
$$R_1 = \frac{10000}{500} = 20 \text{ Halpun}$$
 $0_1 = 0.001 = 0.001 \text{ pumul Ha}$ 
 $R_1 = R_1 + 0_1 = 0.051 \text{ pumul Ha}$ 
 $P_2 = \frac{1}{P_2} + D_2 = 1.02 \text{ pumul Ha}$ 
 $DP_{L2} = 200 = 0.02 \text{ Pu}$ 
 $DP_{L2} = \frac{200}{10000} = 0.02 \text{ Pu}$ 
 $DF = -\frac{DP_{L2}}{B_1 + B_2} = \frac{-0.02}{0.051 + 1.02} = \frac{-0.0187 \text{ Ha}}{0.051 + 1.02}$ 

When load changes on alea 2

$$DP_{L2} = 100 \text{ mW} = \frac{100}{1000} = 0.1P4$$

$$Df = -\frac{\Delta P_{L2}}{\beta_1 + \beta_2} = \frac{-0.1}{\rho_1 + \rho_2} = \frac{-0.58 \times 10^4 \text{ P4}}{\rho_1 + \rho_2 + \rho_2 + \rho_3} = \frac{-0.032 + 3}{\rho_1 + \rho_2 + \rho_3}$$

$$\Delta P_{L12} = \frac{\Delta P_{L2} P_1}{\beta_1 + \beta_2} = \frac{0.1 \times 100.8}{\rho_1 + \rho_2} = 0.0664 \text{ P4}$$

$$\frac{P_{112} = \Delta P_{L2} P_1}{\beta_1 + \beta_2} = \frac{0.1 \times 100.8}{\rho_1 + \rho_2} = 0.0664 \text{ P4}$$