

Ques 1. Find v_a in the circuit shown

Ques 2. Find i_o in the circuit shown

Q.NO3:

Solution $Y_{11} = \frac{I_1}{V_1}\Big|_{V_2=0}$

When $b-b'$ is short circuited, $V_2 = 0$ and the network looks as shown in Fig. 16.8(a).

With port *b-b'* short circuited, $-I_2 = I_1 \times \frac{2}{4} = \frac{I_1}{2}$

$$
\therefore -I_2 = \frac{V_1}{4}
$$

$$
Y_{21} = \frac{I_2}{V_1}\Big|_{V_2=0} = -\frac{1}{4} \text{ } \text{ } U
$$

Similarly, when port $a-a'$ is short circuited, $V_1 = 0$ and the network looks as shown in Fig. 16.8 (b).

$$
Y_{22} = \frac{I_2}{V_2}\Big|_{V_1=0}
$$

$$
V_2 = I_2 Z_{eq}
$$

where Z_{eq} is the equivalent impedance as viewed from $b\hbox{-} b'.$

$$
Z_{eq} = \frac{8}{5} \Omega
$$

\n
$$
V_2 = I_2 \times \frac{8}{5}
$$

\n
$$
Y_{22} = \frac{I_2}{V_2}\Big|_{V_1=0} = \frac{5}{8} \text{ or}
$$

\n
$$
Y_{12} = \frac{I_1}{V_2}\Big|_{V_1=0}
$$

With *a-a'* short circuited, $-I_1 = \frac{2}{5}I_2$

The describing equations in terms of the admittance parameters are

$$
I_1 = 0.5 V_1 - 0.25 V_2
$$

$$
I_2 = -0.25 V_1 + 0.625 V_2
$$

Ques 4: a)

$$
\mathcal{E}_{q} \mathbf{v}_{q} \mathbf{b}_{1} \mathbf{h}_{1} \mathbf{v}_{q} \mathbf{v}_{q} = h_{11} 2_{1} + h_{12} v_{2} - 0
$$
\n
$$
I_{2} = h_{21} 2_{1} + h_{22} v_{2} - 0
$$
\n
$$
I_{3} = h_{31} 2_{1} + h_{22} v_{2} - 0
$$
\n
$$
I_{1} = \left(\frac{1}{h_{11}}\right) \mathbf{B}_{1} \mathbf{v}_{1} + \left(\frac{-h_{12}}{h_{11}}\right) v_{2} - 0
$$
\n
$$
I_{2} = \left(\frac{1}{h_{11}}\right) \mathbf{B}_{1} \mathbf{v}_{1} + \left(\frac{-h_{12}}{h_{11}}\right) v_{2} - 0
$$
\n
$$
I_{3} = h_{21} \left(\frac{1}{h_{11}} v_{1} + \frac{h_{12}}{h_{11}} v_{2}\right) + h_{22} v_{2}
$$
\n
$$
I_{2} = \left(\frac{h_{21}}{h_{11}}\right) v_{1} + \left(\frac{h_{22}}{h_{12}} - \frac{h_{12}}{h_{11}} v_{2} - 0
$$
\n
$$
I_{2} = \left(\frac{h_{21}}{h_{11}}\right) v_{1} + \left(\frac{h_{22}}{h_{11}} - \frac{h_{12}}{h_{11}} v_{2} - 0
$$
\n
$$
I_{3} = \frac{h_{11} h_{22}}{h_{11}} - \frac{h_{12} h_{31}}{h_{11}} - \frac{h_{12} h_{3
$$

Ques 4B:

Ques 5:

Antiresonance Circuit: Case 2-> Practical Circuit for parallel resonance with RL & RC

Resonance in a Two Branch RL and RC Parallel Circuit

Consider the two branch parallel circuit shown in Fig. A Let E be the voltage across each of the parallel circuit shown in the figure. The vector diagram at resonance is shown in Figure B.

$$
V = \frac{R_L}{R_L^2 + \omega_0^2 L^2} + \frac{R_C}{R_L^2 + \frac{1}{\omega_0^2 C^2}} - \frac{j\omega_0 L}{R_C^2 + \omega_0^2 L^2} + \frac{j\omega_0 C}{R_C^2 + \frac{1}{\omega_0^2 C^2}}
$$

\n
$$
Y = G_L + G_C - j(B_L - B_C)
$$

\nFor resonance, $\omega = \omega_0$
\n
$$
B_L = B_C
$$

\n
$$
\Rightarrow \frac{\omega_0 L}{R_L^2 + \omega_0^2 L^2} = \frac{\frac{1}{\omega_0 C}}{R_C^2 + \frac{1}{\omega_0^2 C^2}}
$$

\n
$$
\Rightarrow \frac{\omega_0 L}{R_L^2 + \omega_0^2 L^2} = \frac{\omega_0 C}{\omega_0^2 C^2 R_C^2 + 1}
$$

\n
$$
\Rightarrow L(\omega_0^2 C^2 R_C^2 + 1) = (R_L^2 + \omega_0^2 L^2) C
$$

\n
$$
\Rightarrow \omega_0^2 (C^2 R_C^2 L - L^2 C) = R_L^2 C - L
$$

\n
$$
\Rightarrow \omega_0^2 = \frac{R_L^2 C - L}{C^2 R_C^2 L - L^2 C}
$$

\n
$$
\Rightarrow \omega_0^2 = \frac{1}{LC} \frac{R_L^2 C - L}{R_C^2 C - L}
$$

\n
$$
\Rightarrow \omega_0^2 = \frac{1}{LC} \frac{R_L^2 - \frac{L}{C}}{(R_C^2 - \frac{L}{C})} \Rightarrow \omega_0 = \frac{1}{\sqrt{LC}} \sqrt{\frac{R_L^2 - \frac{L}{C}}{(R_C^2 - \frac{L}{C})}}
$$

This is the expression for resonant frequency. It is to be noted that

1) resonance is not possible for certain combination of circuit elements unlike in a series circuit where resonance is always possible.

2) resonance is also possible by varying of RL or Rc.

Consider the case where

$$
R_L^2 < \frac{L}{C} < R_C^2
$$

Or,

$$
R_C^2 < \frac{L}{C} < R_I^2
$$

In both the cases, the quantity under radical is negative and therefore resonance is not possible.

3) If,
\n
$$
R_L = R_C \neq \sqrt{\frac{L}{C}}
$$
\nThen,
\n
$$
\omega_0 = \frac{1}{\sqrt{LC}}
$$

as in R, L,C series circuit.

If,
$$
R_L = R_C = R = \sqrt{\frac{L}{c}}
$$

Or, $R_L^2 = R_C^2 = R^2 = \frac{L}{c} = X_L X_C$ 6

now,
$$
Y = G_L + G_c - j(B_L - B_c)
$$

\n
$$
Y = \frac{R_L}{R_L^2 + \omega_0^2 L^2} + \frac{R_C}{R_C^2 + \frac{1}{\omega_0^2 C^2}} - \frac{j\omega_0 L}{R_L^2 + \omega_0^2 L^2} + \frac{j\omega_0 C}{R_C^2 + \frac{1}{\omega_0^2 C^2}}
$$
\n
$$
Y = \frac{R_L}{R_L^2 + X_L^2} + \frac{R_C}{R_L^2 + X_C^2} - \frac{jX_L}{X_L X_C + X_L^2} + \frac{jX_C}{X_L X_C + X_C^2}
$$
\n
$$
Y = \frac{R_L}{R_L^2 + X_L^2} + \frac{R_C}{R_C^2 + X_C^2} - \frac{j}{X_C + X_L} + \frac{j}{X_C + X_L}
$$
\n
$$
Y = \frac{R}{R^2 + X_L^2} + \frac{R}{R^2 + X_C^2}
$$
\n
$$
Y = \frac{R(R^2 + X_C^2) + R(R^2 + X_L^2)}{(R^2 + X_L^2)((R^2 + X_C^2))}
$$
\n
$$
Y = \frac{1}{R}
$$
\n
$$
Y = \frac{1}{R}
$$
\n
$$
Y = \frac{1}{R}
$$
\n
$$
Z = R = \sqrt{\frac{L}{C}}
$$
\n9 from (6)

Ques 6: a

Laplace transform of
$$
u(t)
$$

Given
$$
x(t) = u(t)
$$

Then $L{x(t)} = \int_0^{\infty} x(t)e^{-st}dt = \int_0^{\infty} e^{-st}dt = \frac{1}{s}$

$$
\frac{du(t)}{dt} = \delta(t) = \begin{cases} \infty; t = 0 \\ 0; t \neq 0 \end{cases}
$$

Note: Impulse function is only defined at a point in time domain, not before and after that

Area under an impulse is unity.

$$
\int_0^\infty \!\!\delta(t) dt = 1
$$

Laplace transform of $\delta(t)$

Given
$$
x(t) = \delta(t)
$$

Then $L\{x(t)\} = \int_0^\infty x(t)e^{-st}dt = \int_0^\infty \delta(t)e^{-st}dt = \int_0^\infty \delta(t)dt = 1$

Ques 6B:

Ques 8A:

Initial-value theorem

The initial-value theorem allows us to find the initial value $x(0)$ of the function $x(t)$, directly from its Laplace transform $X(s)$

If $x(t)$, is a causal signal, then $x(0) = \lim_{s \to \infty} sX(s)$

Proof: To prove this theorem, we use time differentiation property

$$
\Rightarrow L\left\{\frac{d}{dt}x(t)\right\} = sX(s) - x(0) = \int_0^\infty \frac{d}{dt}x(t)e^{-st}dt
$$

Taking the limit $s \to \infty$ on both sides of above equation If we let $s \to \infty$ then the integral on the right side of equation vanishes due to damping factor, e^{-st}

$$
\lim_{s \to \infty} (sX(s) - x(0)) = 0
$$

\n
$$
\Rightarrow x(0) = \lim_{s \to \infty} sX(s)
$$

Ques8_B

Final-value theorem

The final-value theorem allows us to find the final value $x(\infty)$ of the function $x(t)$, directly from its Laplace transform $X(s)$

If $x(t)$, is a causal signal, then $x(\infty) = \lim_{s \to 0} sX(s)$

Proof: To prove this theorem, we use time
differentiation property

 $\Rightarrow L\left\{\frac{d}{dt}x(t)\right\}=sX(s)-x(0)=\int_0^\infty\frac{d}{dt}x(t)e^{-st}dt$ Taking the limit $s \to 0$ on both sides of above equation If we let $s \to 0$ then the integral on the right side of equation reduces to $\int_0^\infty dx(t)$

$$
\lim_{s \to 0} (sX(s) - x(0)) = \int_0^{\infty} \frac{dt}{dt} dt
$$
\n
$$
\lim_{s \to 0} S(X(s) - x(0)) = x(t) \Big|_0^{\infty}
$$
\n
$$
\lim_{s \to 0} S(X(s) - x(0)) = x(\infty) - x(0)
$$
\n
$$
\Rightarrow x(\infty) = \lim_{s \to 0} sX(s)
$$

$$
x_1 \neq 0
$$

\n $x_2 \neq 0$
\n $x_3 \neq 0$
\n $x_4 \neq 0$
\n $x_5 \neq 0$
\n $x_6 \neq 0$
\n $x_7 \neq 0$
\n $x_8 \neq 0$
\n $x_9 \neq 0$
\n $x_9 \neq 0$
\n $x_1 \neq 0$
\n $x_2 \neq 0$
\n $x_3 \neq 0$
\n $x_4 \neq 0$
\n $x_5 \neq 0$
\n $x_6 \neq 0$
\n $x_7 \neq 0$
\n $x_8 \neq 0$
\n $x_9 \neq 0$ <