

Internal Assessment Test - II

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|-------|--------------------|-----------|---------|------------|----|------|-------|----------|-------------|
| Sub: | Electronic Devices | | | | | | Code: | 18EC33 | |
| Date: | 25/01/2022 | Duration: | 90 mins | Max Marks: | 50 | Sem: | III | Branch : | ECE-A,B,C,D |

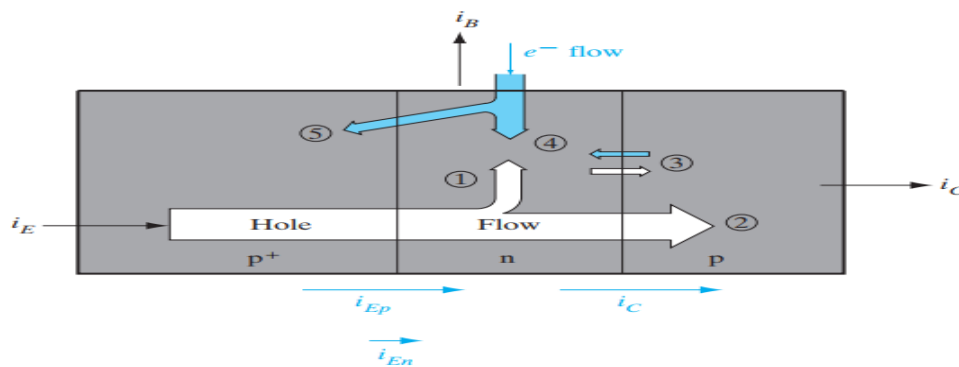
Answer Any FIVE FULL Questions

1. With a neat diagram explain the hole and electron flow in P⁺-N-P transistor with proper Biasing. What are the mechanisms (factors) which cause base current in a transistor?

Marks

[7+
3]

| | |
|-----|-----|
| OBE | |
| CO | RBT |
| CO2 | L2 |

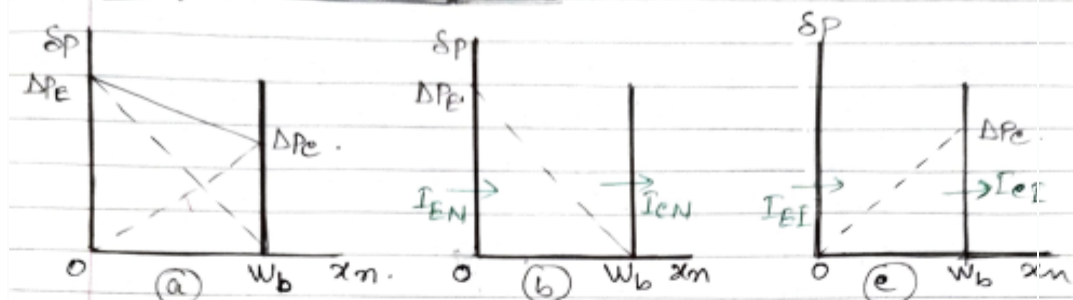


- $I_B = \text{small, as } I_E \text{ is emitted } h^+ \text{ I \& } I_C \text{ is collected } h^+ \text{ I} - I_C \sim I_E.$
 - 1. injected h^+ lost to recombination in B.
 - 2. h^+ reaching RB J_c .
 - 3. thermally generated e^- & h^+ making up reverse I_{sat} of J_c .
 - 4. e^- supplied by B contact for recombination with h^+ .
 - 5. e^- injected at FB J_e .
- I_b can be accounted by three dominant mechanisms:
 - (a) There must be some recombination of injected h^+ with e^- in the base, even with $W_b \ll L_p$. e^- lost to recombination must be resupplied through the base contact.
 - (b) Some e^- will be injected from n to p in the FB J_e , even if the emitter is heavily doped compared with the base. These e^- must also be supplied by I_b .
 - (c) Some e^- are swept into the base at the RB J_c due to thermal generation in the collector. This small current reduces I_b by supplying e^- to the base.
- Dominant sources of I_b are : (i) recombination in B (ii) injection into E region.

2. Derive the Ebers Moll equations of a PNP BJT and draw the Ebers Moll model (Diode coupled Model) diagram. Also explain the terms involved in it. [10]

CO4 L3

1.5 of Coupled-Diode Model



(a) Approximate h^+ distribution in B with J_E & $J_C = FB$.
 (b) Component due to injection & collection in normal mode.
 (c) Component due to inverted mode.

- ✓ $\rightarrow J_E$ & J_C are FB as shown $\therefore \Delta p_E$ & $\Delta p_C = +ve$ no.
- ✓ \rightarrow straight line h^+ distribution of (a) can be broken into 2 components of (b) & (c).
- ✓ (b) $\rightarrow h^+$ injected by E & collected by C. Resulting $I_s \rightarrow I_{EN}$ & I_{CN} as Normal mode components as they are due to injection from $E \rightarrow C$.
- ✓ (c) \rightarrow components of h^+ distribution results in I_{El} & I_{Cl} as Inverted mode injection from $C \rightarrow E$ these inverted components will be $-ve$ as they are opposite to J_E & J_C .

\rightarrow For symmetrical TS these components are described by:

$$I_{EP} = qA \frac{Dp}{Lp} \left[\frac{\Delta p_E (e^{Wb/Lp} + e^{-Wb/Lp}) - 2\Delta p_C}{e^{Wb/Lp} - e^{-Wb/Lp}} \right]$$

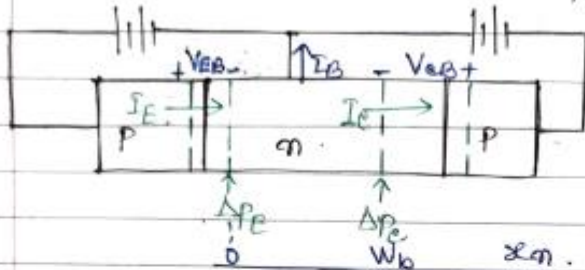
$$\Rightarrow I_{EP} = qA \frac{Dp}{Lp} \left(\Delta p_E \coth \frac{Wb}{Lp} - \Delta p_C \operatorname{csch} \frac{Wb}{Lp} \right)$$

$$I_C = qA \frac{Dp}{Lp} \left(\Delta p_E \operatorname{csch} \frac{Wb}{Lp} - \Delta p_C \coth \frac{Wb}{Lp} \right)$$

\rightarrow Excess h^+ conc. at the edge of E DR Δp_E & corresponding h^+ conc. on the C side Δp_C is

$$\rightarrow \Delta p_E = p_m (e^{qV_{EB}/kT} - 1) \quad \& \quad \Delta p_C = p_m (e^{qV_{CB}/kT} - 1) \quad - (1)$$

\rightarrow If JE is strongly FB $\rightarrow V_{EB} \gg kT$ & if JE is strongly RB then $V_{EB} \ll 0$. $\therefore \Delta p_E \propto p_m e^{qV_{EB}/kT}$ & $\Delta p_C \propto -p_m$



By Superposition Theorem.

$$\rightarrow I_{EN} = a \Delta p_E \quad \& \quad I_{CN} = b \Delta p_E \quad \text{with } \Delta p_C = 0$$

$$I_{E1} = -b \Delta p_C \quad \& \quad I_{C1} = -a \Delta p_C \quad \text{with } \Delta p_E = 0$$

\rightarrow 4 components are combined by linear superposition

$$I_E = I_{EN} + I_{E1} = a \Delta p_E - b \Delta p_C = A (e^{qV_{EB}/kT} - 1) - B (e^{qV_{CB}/kT} - 1)$$

$$I_C = I_{CN} + I_{C1} = b \Delta p_E - a \Delta p_C = B (e^{qV_{EB}/kT} - 1) - A (e^{qV_{CB}/kT} - 1) \quad - (2)$$

where $A = a p_m$ & $B = b p_m$.

\rightarrow I_E in normal mode is given by

$$I_{EN} = I_{ES} (e^{qV_{EB}/kT} - 1), \quad \Delta p_C = 0 \quad - (3)$$

I_{ES} = magnitude of I_{ESat} in normal mode, with $J_c = SC$ $\therefore V_{CB} = 0$.

$\therefore \Delta p_C = 0$ $\therefore V_{CB} = 0$.

\rightarrow I_C in inverted mode.

$$I_{C1} = -I_{CS} (e^{qV_{CB}/kT} - 1), \quad \Delta p_E = 0 \quad - (3)$$

I_{CS} = magnitude of I_{CSat} with $V_{EB} = 0$

\rightarrow -ve sign of I_{CS} \Rightarrow h^+ are injected opposite to the defined I_C for inverted mode.

→ Collected I for each mode of operation.

$$\left. \begin{aligned} I_{EN} &= \alpha_N I_{EN} = \alpha_N I_{ES} (e^{qV_{EB}/kT} - 1) \\ I_{EI} &= \alpha_I I_{EI} = -\alpha_I I_{ES} (e^{qV_{EB}/kT} - 1) \end{aligned} \right\} (4)$$

where α_N & α_I = ratios of collected I to injected I in each mode.

→ In inverted mode → injected I = I_{EI} & collected I = I_{EI} .

→ Total I obtained by superposition of components

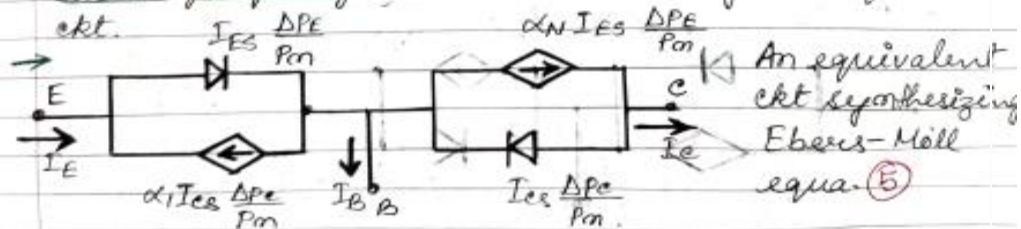
$$\left. \begin{aligned} I_E &= I_{EN} + I_{EI} = I_{ES} (e^{qV_{EB}/kT} - 1) - \alpha_I I_{ES} (e^{qV_{EB}/kT} - 1) \\ I_C &= I_{EN} + I_{CI} = \alpha_N I_{ES} (e^{qV_{EB}/kT} - 1) - I_{ES} (e^{qV_{EB}/kT} - 1) \end{aligned} \right\} (5)$$

→ These relations were derived by J.J. Ebers & J.L. Moll and are referred to as Ebers-Moll equations

→ Equa. (2) is for symmetrical T_s , these equa. varies in I_{ES} , I_{CS} , α_I & α_N due to asymmetry b/w the junc.

→ $\alpha_N I_{ES} = \alpha_I I_{CS}$ can be shown for non-symmetrical T_s .

→ Feature of Ebers-Moll equa. → I_E & I_C are described by terms resembling diode relations (I_{EN} & I_{CI}) + terms that give coupling b/w the properties of E & C (I_{EI} & I_{CN}). This coupled-diode property is illustrated by the equivalent



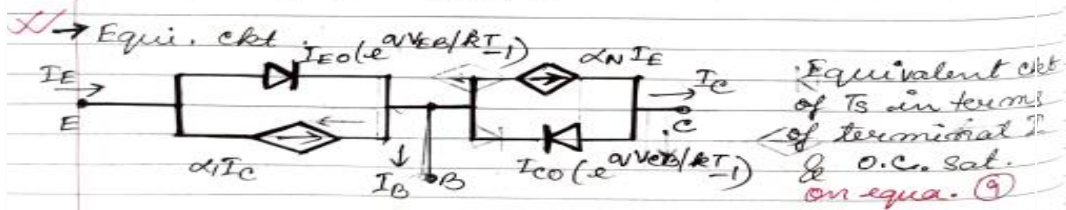
$$I_B = (1 - \alpha_N) I_{ES} \frac{DPE}{Pm} + (1 - \alpha_I) I_{CS} \frac{DPC}{Pm}$$

✓ → saturation I from the coupling term in each part of equa (5) can be removed by multiplying equa. (5a) by α_N & subtracting the result from (5b)

$$\left. \begin{aligned} I_C &= \alpha_N I_E - (1 - \alpha_N \alpha_I) I_{CS} (e^{qV_{EB}/kT} - 1) \\ I_E &= \alpha_I I_C + (1 - \alpha_N \alpha_I) I_{ES} (e^{qV_{EB}/kT} - 1) \end{aligned} \right\} (8)$$

✓ → The term $(1 - \alpha_N \alpha_I) I_{CS} = I_{CO}$ & $(1 - \alpha_N \alpha_I) I_{ES} = I_{EO}$
 I_{CO} mag. of I_{CSat} with I_E open ($I_E = 0$)
 $\& I_{EO} =$ " " I_{ESat} " I_C open.

✓ → Ebers-Moll equa. becomes (substituting in (8))

$$\left. \begin{aligned} I_E &= \alpha_I I_C + I_{EO} (e^{qV_{EB}/kT} - 1) \\ I_C &= \alpha_N I_E - I_{CO} (e^{qV_{EB}/kT} - 1) \end{aligned} \right\} (9)$$


3. a) Define and give expressions for
- Base transport factor
 - Emitter injection efficiency
 - Current transfer ratio
 - Base to collector current amplification factor.

[1+1+2+1]

[5]

CO2 L1

CO1 L2

Amplification with BJTs

→ BJT works as amplifier as $\Rightarrow I_C$ & I_E are controlled by the relatively small I_B .

→ Neglecting, ac analysis, I_{sat} at c, recombinations. I_C is made up of h^+ injected at E which are not lost to recombination in B.

∴ $I_C \propto h^+$ component of I_E .

$$I_C = \beta I_{E_p} \quad (1)$$

✓ β = proportionality factor = fraction of injected h^+ which make it across the B to C.

β = Base transport factor.

→ Total i_E is made up of h^+ component i_{E_p} & e^- component i_{E_m} due to e^- injected from B \rightarrow E.

$$I_{ET} = I_{E_p} + I_{E_m}$$

E \rightarrow B B \rightarrow E

II → Emitter Injection Efficiency γ .

$$\gamma = \frac{i_{EP}}{i_{EM} + i_{EP}} \quad (2)$$

→ For efficient T3 B & γ near to 1 $\Rightarrow i_e$ is max
 (vi) due to h^+ \therefore most of the injected h^+ should contribute to i_c ($\beta \approx 1$)

→ Relation b/w i_c & i_e .

$$\frac{i_c}{i_e} = \frac{\beta i_{EP}}{i_{EM} + i_{EP}} = \beta \gamma = \alpha \quad (3)$$

II → Current Transfer Ratio $\alpha = \beta \gamma$ \Rightarrow represent E to C I amplification.

→ as $\alpha < 1 \Rightarrow$ there is no real I. amplification
 i_e & i_b are more for amplification.

→ For $i_b \Rightarrow$ Rates at which e^- are lost from B by injection across JE (i_{EP} be rate of e^- i_{EM} i_{EP} recombination with h^+ in the B are included

→ For each case lost e^- must be resupplied through i_b .

→ If fraction of injected h^+ making it across the B without recombination = β , then $(1-\beta)$ is the fraction recombining in the B.

$$\therefore i_b = i_{EM} + (1-\beta)i_{EP} \quad \text{neglecting } I_{csat} \quad (4)$$

→ Relation b/w i_c & i_b is found from (1) & (4).

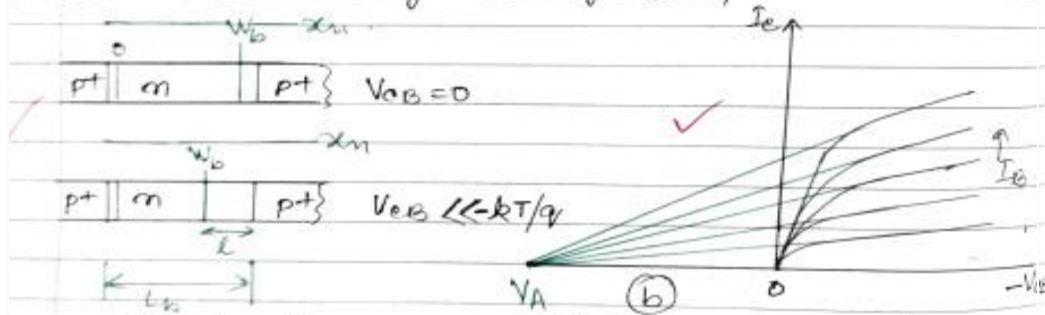
$$\frac{i_c}{i_b} = \frac{\beta i_{EP}}{i_{EM} + (1-\beta)i_{EP}} = \frac{\beta [i_{EP} / (i_{EM} + i_{EP})]}{1 - \beta [i_{EP} / (i_{EM} + i_{EP})]}$$

$$\Rightarrow \frac{i_c}{i_b} = \frac{\beta \gamma}{1 - \beta \gamma} = \frac{\alpha}{1 - \alpha} = \beta$$

⇒ Base-to-Collector Current Amplification Factor
 $= \beta$

b) Explain base narrowing effect (Early effect) in P⁺-N-P⁺ transistor.

7-3 Base Narrowing → Early Effect / Base Width Modulation



$W_b = L_b - l$
 $l \propto \sqrt{V_{CE}}$
 (a) β width decreases with increase in $I_C = RB$.
 (b) increase in I_C increases V_C . Extrapolation of curve to V_A .

- If β is lightly doped, DR at $I_C = RB$ can extend into n-type B.
- As V_C increases, space charge layer takes up more of the β width L_b & decrease W_b .

- This effect is called Base Narrowing / Early effect / Base-width Modulation. J.M. Early.
- Effect of this is shown in I_C characteristics of CE Conf.
- Decrease in W_b causes β to increase → I_C increases with V_C instead of remaining const.
- Slope of Early effect is almost linear with I_C & CE graph extrapolate to an intersection with V axis V_A — called Early V.

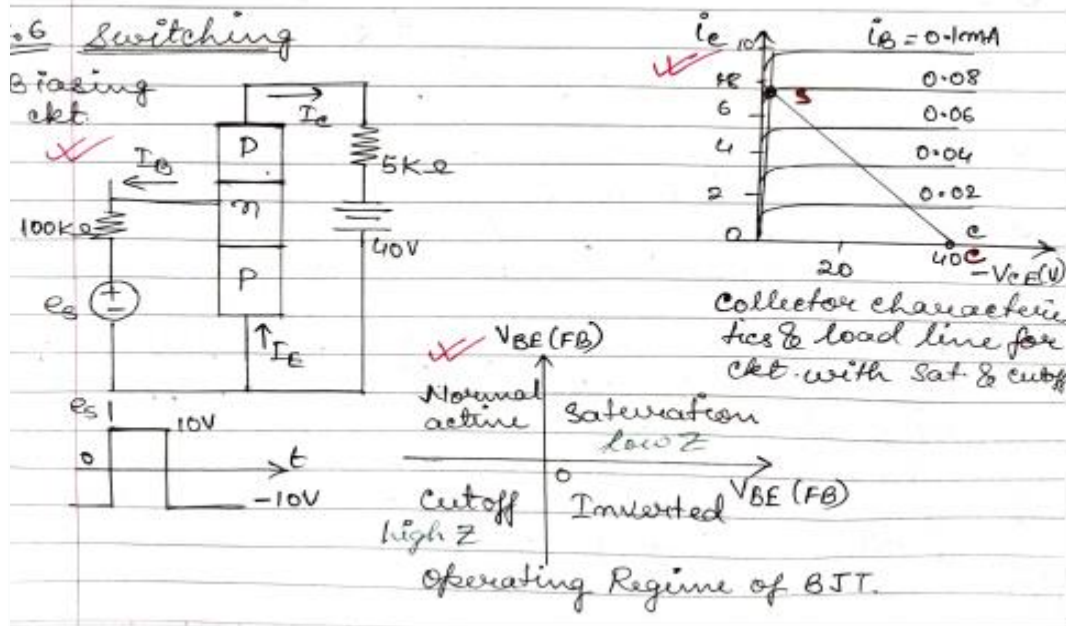
→ $l =$ length of DR into n B.

$$l = \left(\frac{2\epsilon V_{CE}}{qNd} \right)^{1/2}$$

- If RB on I_C increases far enough, W_b reduces such that I_C DR fills the entire β region giving Punch through condition → where W are swept directly from E → C & T_s action is lost. Punch through is a breakdown effect & should be avoided.
- Avalanche breakdown of I_C occurs before Punch-through is reached.

4. Explain switching operation of a transistor in cut off and saturation region. What are the various mechanisms of switching cycle?

[6+4] CO2 L2



- ✓ → If i_B is such that the operating pt lies somewhere b/w the 2 end pts of the load line T_S operates in the normal mode. i.e. $I_E = I_C$ & $I_C = \beta I_B$ with sufficient i_B flowing out of T_S .
- ✓ → $i_B = 0$ / $-ve$ pt C is reached at the bottom end of the load line & $i_C = \text{negligible}$ i.e. OFF state of T_S & device in cutoff regime.
- ✓ → If $i_B = +ve$ & quite large, device is driven to saturation region → **S**. This is a state of T_S where large i_C flows with only very small V_{CE} drop.
- ✓ → Beginning of sat. region corresponds to loss of R_B across the T_C .
- ✓ → In switching operation i_B swings from $+ve$ to $-ve$ driving the device from sat to cutoff.
- ✓ → From the regions of operation of BJT if $I_E = I_B$ & $I_C = R_B$ → ^{Normal} active mode, opposite to that is inverted mode.
Both I_E & $I_C = R_B$ → cutoff → very high impedance state of BJT & both I_E & $I_C = I_B$ → Saturation with low impedance state.

1.6.3 Switching Cycle →

→ Various mechanisms of a switching cycle are shown



(a) Switching effects in CE T_c

- t_0 = Cutoff.
- t_1 = Normal active region.
- t_2 = Beginning of saturation region.
- t_3 = Final saturated state.

(b) h^+ distribution in B during switching from Cutoff to saturation

(c) i_B , stored charge Q_B & i_C during turn on & turn off of T_c

- If the device is originally in cutoff, a step increase of i_B to I_B causes the h^+ distribution to increase as shown in (b)
- At t_2 - device enters saturation & the h^+ distribution reaches its final state at t_3
- As stored charge in B Q_B increases, i_C increases
- i_C does not increase beyond its value at the beginning of saturation: t_2 . $I_C \approx E_{cc} / R_C \gg I_{Csat}$.
- E_{cc} = value of ckt battery & R_C = load R.
- $I_C \approx 80mA$.
- i_C increases exponentially as Q_B rises to Q_B at t_3 → this rise time serves as one of the limitations of the T_c in switching application.

- when i_B is -ve to $-I_B$, stored charge Q_B must be withdrawn from the B before cutoff is reached.
- while Q_B is larger than Q_B i_C remains at I_C fixed by battery & R_C .
- there is storage delay time t_{sd} after the t_3 is switched & before i_C begins to fall towards 0.
- when $Q_B < Q_B$, i_C drops exponentially with characteristic fall time..
- Once stored charge is withdrawn, i_B cannot be maintained any longer at its large -ve value & must decay to small cutoff value.
- $i_B = -(1-\alpha_N) I_{ES} - (1-\alpha_I) I_{CE}$.

5a). Write the expression for current in photodiode resulting due to collection of optically generated carriers and open circuited voltage across the junction.

| | | |
|------|-----|----|
| [6] | CO2 | L2 |
| [4] | | |
| [5] | CO2 | L1 |
| [5] | | |
| [10] | CO2 | L2 |

8.1 Photodiodes \Rightarrow 2 terminal devices designed to respond to photon absorption \Rightarrow Photodiodes.
Some photodiodes have high sensitivity & response speed.

8.1.1 Current & Voltage in an Illuminated Junction

- \rightarrow Carriers generated within the DR are separated by the junc. field \vec{E} in a region $\approx h^+$ in p region.
- \rightarrow If junc is uniformly illuminated by photon with $h\nu > E_g$ the generation rate g_{op} (EHP/cm³s) gets added and participates in I .

The diagram shows a P-n junction with P on the left and n on the right. Light with energy $h\nu > E_g$ is incident on the P-region. Carriers are generated in a region of width W . The P-region has a diffusion length L_p and the n-region has a diffusion length L_n . The total width of the device is x_{no} . The built-in electric field E is shown pointing from n to p. A graph shows the current I versus voltage V , with $I_{op} = 0$ at $V = 0$ and a curve that rises exponentially for positive bias. The graph is labeled with g_1, g_2, g_3 and $g_3 > g_2 > g_1$.

Light absorption by device.

$\delta P_{op} = g_{op} \tau_p$
 $I_{op} = qAL_p g_{op}$
 I_{op} due to EHP generation within diffusion length W of the junction at n .

I - V characteristic of illuminated junc.

- $\rightarrow AL_p g_{op} =$ no. of h^+ created per second within diffusion length of the DR on n side (L_p of x_{no})
 - $\rightarrow AL_n g_{op} =$ no. of e^- generated per second within L_n of x_{po} .
 - $\rightarrow AW g_{op} =$ carriers generated within W .
 - $\rightarrow I$ due to all these optically generated carriers by the junc
- $$I_{op} = qA g_{op} (L_p + L_n + W)$$
- $L_p =$ length of h^+ diffusion $L_n =$ length of e^- diffusion

→ Thermally generated I due to carrier injection
 $I = I_{th} (e^{qV/kT} - 1)$ $I = qA \left(\frac{L_p p_{n0}}{\tau_p} + \frac{L_n n_{p0}}{\tau_n} \right) (e^{qV/kT} - 1)$ (1)
 $I_{th} = \text{thermally generated } I = \left(\frac{L_p p_{n0}}{\tau_p} + \frac{L_n n_{p0}}{\tau_n} \right) qA$

→ ∴ Total reverse I with illumination
 $I = I_{th} (e^{qV/kT} - 1) - I_{op}$
 $I = qA \left(\frac{L_p p_{n0}}{\tau_p} + \frac{L_n n_{p0}}{\tau_n} \right) (e^{qV/kT} - 1) - qA g_{op} (L_p + L_n + W)$ (2)

→ ∴ I-V graph is lowered by an amount \propto to the generation rate, g_{op} .

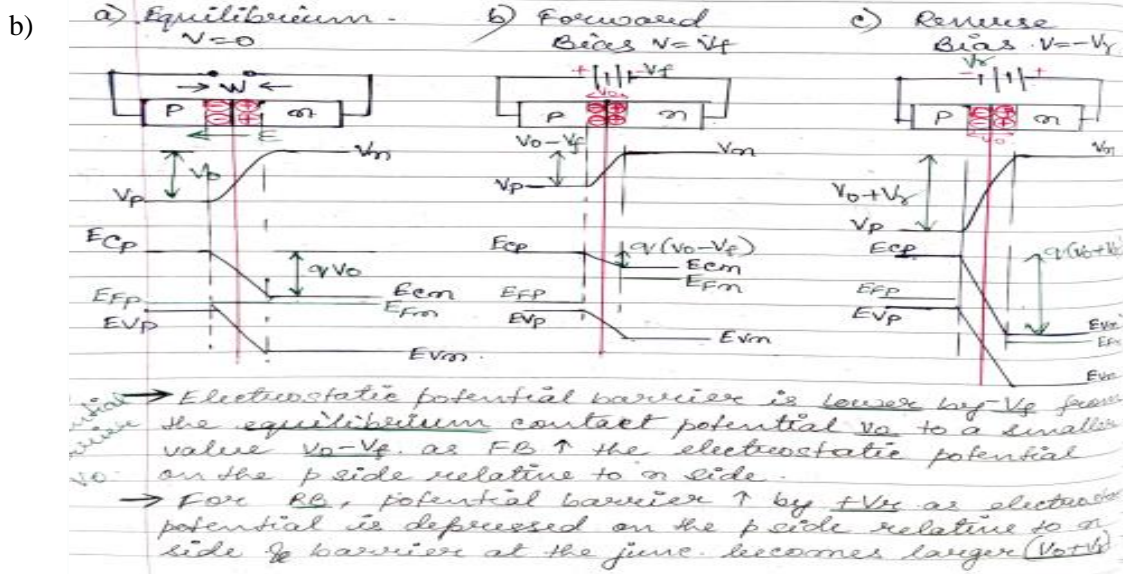
→ when device is SC $V=0 \Rightarrow$ current from the diode eqn. cancels & there is a I_{sc} from $n \rightarrow p$ equal to I_{sp} $I_{sc} = I_{op} \rightarrow$ I-V graph cross the I-axis at $-I_{sc}$ values $\propto g_{op}$.

→ when device is OC $I=0$ & $V = V_{oc}$.
 $V_{oc} = \frac{kT}{q} \ln \left(\frac{I_{op}}{I_{th}} + 1 \right)$

$\checkmark \Rightarrow V_{oc} = \frac{kT}{q} \ln \left[\frac{L_p + L_n + W}{(L_p/\tau_p) p_{n0} + (L_n/\tau_n) n_{p0}} g_{op} + 1 \right]$ (3)

$\checkmark \rightarrow$ For a symmetrical junction, $p_{n0} = n_{p0}$ & $\tau_n = \tau_p$
 (2) can be written as thermal generation rate $p_{n0}/\tau_n = g_{th}$ & the optical generation rate g_{op} . Neglecting generation within W .
 $\checkmark V_{oc} = \frac{kT}{q} \ln \frac{g_{op}}{g_{th}}$ for $g_{op} \gg g_{th}$.
 $g_{th} = p_{n0}/\tau_n = \text{equilibrium thermal generation-recombination rate.}$

Draw and explain the energy band diagram of a PN junction i) at equilibrium, ii) under forward bias and iii) under reverse bias. Explain qualitatively the effect of biasing on current using current equation and I-V characteristic graph.



→ For FB → EF at the transition region ↓ as the applied field V_f opposes the built-in field E
 For RB → EF at the transition region ↑ as applied field V_f & built in field are in same direction

→ Changes in EF E changes with W . W ↓ in FB & W ↑ in RB.

$$W = \left[\frac{2\epsilon V_0 (N_a + N_d)}{q N_a N_d} \right]^{1/2} = \left[\frac{2\epsilon V_0}{q} \left(\frac{1}{N_a} + \frac{1}{N_d} \right) \right]^{1/2} \text{ can be}$$

used to calculate width

→ Energy band separation in FB is given by $q(V_0 - V_f)$ and in RB by $q(V_0 + V_f)$

→ Under FB E_{cm} is above E_{cp} by the energy qV_f
 & for RB E_{cp} is above E_{cm} by qV_f .

→ In FB ⇒ barrier is lowered to $(V_0 - V_f)$ → more e^- get sufficient energy to diffuse from $n \rightarrow p$

∴ $I_{(e^- \text{ diff})}$ or $I_{(n \text{ diff})}$ increases. } FB
 & also more h^+ can diffuse from $p \rightarrow n$ } $I_{\text{diff}} \uparrow$
 ∴ $I_{(h^+ \text{ diff})}$ or $I_{(p \text{ diff})}$ increases.

→ For RB ⇒ barrier is large to $(V_0 + V_f)$ & e^- & h^+ can't cross over ∴ Diffusion I is negligible $I_{\text{diff}} \downarrow$

→ Drift I is not related to barrier potential height

* → Minority carriers contributing to drift component are generated by thermal excitation of EHP. I due to drift of generated carriers across the junction is called Generation Current as its magnitude depends on rate of generation of EHP. I_g can be increased by optical excitation at the junction for photodiode.

→ Total I crossing junc. = $I_{diff} + I_{drift}$.
 as the I s are in opposite direction net $I = 0$.
 → Under RB both $I_{diff} = 0$ & a small I_g current flows due to ~~the~~ minority carriers & large bias.

$I-V$ character of pn junc.

→ I_g is shown in the $I-V$ plot.
 → +ve direction of I is from $p \rightarrow n$ by applied $V = +ve$.
 → For $V = -ve$, small I_g flow due to carriers generated in the transition region.

→ At $V=0$, $I=0$ as I_g & I_{diff} cancel.

$$I = I_{(diff)} - |I_{(gen)}| = 0 \text{ for } V=0.$$

→ $V = V_f$ increases the probability that a carrier can diffuse across the junc. by the factor $\exp(qV_f/kT)$.
 I_{diff} under FB is given by its equilibrium value $\times \exp(qV/kT)$.

→ $V = -V_f$ I_{diff} is equilibrium value reduced by $\exp(-qV/kT)$.

At equilibrium $I_{diff} = |I_{(gen)}|$.

→ Total $I = \text{diffusion } I - \text{absolute value of } I_{gen} \Rightarrow I_0$.

$$I = I_0 (e^{qV/kT} - 1)$$

when $V = +V_f \Rightarrow \text{exponential term} \gg 1 \therefore I \uparrow$ exponentially with FB.

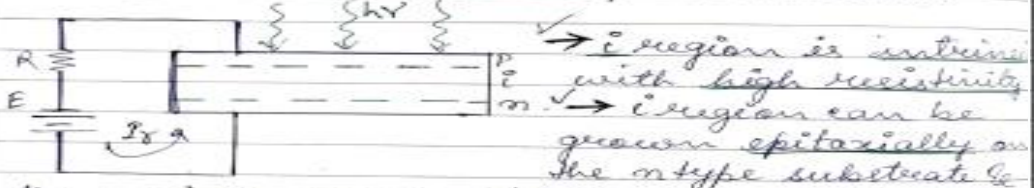
when $V = -V_f \Rightarrow \text{exp. term} \rightarrow 0 \therefore I = -I_0 \Rightarrow \text{Reverse Saturation Current}$.

6. (a) What are the advantages of using PIN photodiode in photo-detector? What is the figure of merit of a photo-detector?

1.3 Photodiodes

- When photodiode is operated in the 2nd quadrant of I-V characteristic, I is independent of V but I is \propto optical generation rate.
- This can be used to measure illumination or converting time-varying optical signal into electrical signal.
- In optical detection operation, detector's speed of response or BW is important.
- Eg → If a photodiode is to respond to a series of light pulses 10ns apart, the photogenerated minority carriers must diffuse to the junction & be swept across the other side in a time much $<$ 10ns.
- DR width w should be large so that most of the photons are absorbed within w rather than in the neutral n & p regions.
- When EHP is created in the DR, the EF sweeps the e^- to n side & h^+ to p side. As this carrier drift occurs very in very short time, \therefore response of photodiode is quite fast.
- When carriers are generated within the DR w , the detector is called depletion layer photodiode.

- One side of the junction is doped lightly so that w can be made large.
- Appropriate width for w is chosen b/w sensitivity & speed of response.
- If w is wide most of the incident photons will be absorbed in the DR giving high sensitivity.
- wide w → small junction capacitance $C_j = \frac{EA}{w}$ which reduces RC time const of the detector ckt.
- But w must not be so wide that the time needed for drift of photogenerated carriers out of the DR is more → low BW.
- To control w → p-i-n photodiode.



- The p region can be obtained by implantation.
- When pin diode is in RR the applied voltage appears almost entirely across the i region.
- If carrier lifetime τ within the i region is long compared with the drift time, most of the photogenerated carriers will be collected by the n & p region.
- Important figure of merit of the photodiode is the external quantum efficiency η_{ext} defined as the no. of carriers that are collected for every photon impinging on the detector.

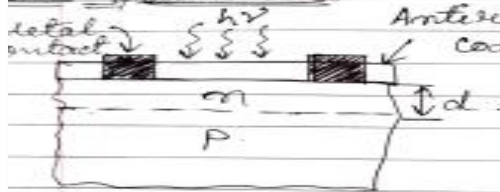
- ✓ → For a photo I density J_{op} we collect J_{op}/q carriers per unit area per sec.
- ✓ → For an incident optical P. density P_{op} , the no. of photons shining on the detector per unit area per sec. is $P_{op}/h\nu$.
 $\therefore \eta_e = (J_{op}/q) / (P_{op}/h\nu)$
- For a photodiode with no gain $\eta_{max} = 1$.
- ✓ → If a low level optical signal are to be detected it is expected that the device operate in avalanche regions, where for every photo-generated carriers there is a large M due to avalanche multiplication giving gain \approx external quantum eff. $> 100\%$.
- Avalanche photodiodes APD are useful as detectors in fiber optic system.
- ✓ → these photodiodes are sensitive to photons with energies near E_g . If $h\nu < E_g \rightarrow$ photon will not be absorbed & if $h\nu > E_g \rightarrow$ photons will be absorbed very near to surface where recombination rate is high. \therefore a photodiode material is chosen with band gap corresponding to a particular region of spectrum.
- Detectors sensitive to longer wavelength can be designed such that photons can excite e^- from the material-intrinsic detectors.
- wider E_g material can be used as a window through which light is transmitted to the absorbing region - using multilayer junc.
- ✓ → InGaAs has $E_g = 0.75\text{eV}$ sensitive to a $\lambda = 1.55\ \mu\text{m}$.

(b) Write a note on LED mentioning the applications of LED. What is external quantum efficiency of an LED?

- Light-Emitting Diodes → when carriers are injected across the FB junc. the I is usually accounted for by recombination in the injection region & in the neutral regions near the junc. In sc with indirect BG, Si, Ge the recombination releases heat to the lattice.
- For direct sc light is emitted due to recombination across the FB junc. → this effect is Injection Electroluminescence → gives an imp. application of diodes as generator of light.
- LED application
 ↳ Digital display, traffic & automotive signals and other illumination.
- Sc laser is also another device making use of radiative recombination in a FB junc.
- Lasers emit coherent light in much narrower λ bands than LED with more collimation (directionality) & useful for OPC.
- LASER - Light amplification by stimulated Emission of Radiation.
- For LED, freq of the photon (color) depends on E_g of sc given by Planck Relation $h\nu = E_g$
 $E_g (\text{eV}) = 1.24 / \lambda (\text{nm})$
- Imp. metric of LED = external quantum eff.
 η_{ext} = defined as light o/p / electrical P_{in}
 = (Internal radiative efficiency) \times (Extraction efficiency).

Explain the operation of a typical solar cell with a neat diagram. Define fill factor with diagram and describe efficiency of power generation of a solar cell.

1.2 Solar Cells



power can be delivered to an external ckt by an illuminated junction. it is possible to convert solar energy into electrical energy using 1st quarter junction.

- But less power is delivered by a single junction. $V_{is} < V_o$ for Si $V_{oc} < 1V$ & $I_{op} = 10-100mA$ for a junction of area = $1cm^2$. Single junction delivers less.
- If many devices are used then power is significant.
- Arrays of p-n junctions are used to supply electric power for satellites, over a long period of time.
- Array of p-n junctions can be placed on the satellite body / placed in solar cell paddles attached to the main body.
- Eg → a solar wing measuring 74m tip to tip can each have 32,800 solar cells & can generate 62KW of power.
- To utilize maximum amount of optical energy that is available, solar cells are designed with large area junction located near device surface.

Planar junction is formed by diffusion or ion implantation to the surface is coated with appropriate materials to reduce reflection & to decrease surface recombination.

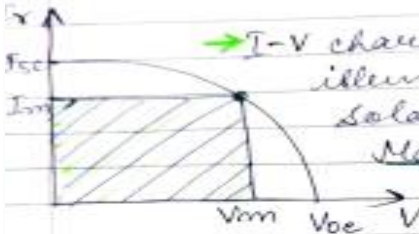
- Many compromises are made in solar cell design.
- Junction depth $d < L_p$ in the n material to allow hole generation near the surface to diffuse to the junction before they recombine. If p region thickness must be such that \bar{E} generated in this region can diffuse to the junction before they recombine.
- Proper matching is needed b/w L_n, p region thickness & α absorption coeff.

$$I(x) = I_0 e^{-\alpha x}$$
- If contact potential V_o is large, photovoltage will be large. ∴ heavy doping is needed.
- Series resistance of the device should be very small so that power is not lost to heat due to ohmic losses in the device. Series resistance of only few ohms can seriously reduce the output of a solar cell.
- Since area is large, resistance of p-type body of the device can be made small.
- Contacts to the thin n region requires special design. If this region is contacted at the edge, I must flow along the thin n region to the contact giving a large series resistance.
- To prevent this the contact must be distributed over the n surface by providing by providing small contact fingers.

Metal contact fingers



→ Narrow contacts reduce the series resistance without interfering with the incoming light



→ I-V characteristics of 4th quadrant illuminated solar cells. Max. power I_x plotted upward. Rectangle.

→ V_{oc} & I_{sc} are determined for a given light level by the cell properties.

→ Max Power is delivered to the load by the solar cell when product $V \times I = \text{max}$.

→ Max Power delivered is given by shaded rectangle $< I_{sc} V_{oc}$ $V_m I_m < V_{oc} I_{sc}$

→ Ratio $\frac{I_m V_m}{I_{sc} V_{oc}}$ = fill factor = figure of merit for solar cell design. FOM.

→ Solar cell application

↳ Outer space

↳ Terrestrial application even though solar intensity is reduced by atmosphere.

→ Currently total power generation $\sim 15 \text{ TW}$ to annual energy usage of $\sim 500 \text{ quads}$ ($1 \text{ quad} = 10^{15}$ or a quadrillion BTU) with an annual increase of 1-2%. From this 80% comes from fossil fuels.

→ Photovoltaic gives a capacity of 100 GW .