



### VTU BE Degree Examination – Feb/Mar 2022



## 17EC71

## Module-4

7\_a<sup>-</sup> Plot the field pattern for an array of 2 isotropic sources with equal amplitude and same phase. Take  $d = \lambda/2$ . Find Directivity of a source with a sine squared pattern (doughnut) (power pattern). (07 Marks)  $(66 Marks)$ 

State and explain power theorem.

Obtain the field pattern for a linear uniform array of isotropic antennas for  $n = 6$ ,  $d =$ a

 $(08$  Marks)

- $\partial = -d$ , Obtain an expression for radiation resistance of a short dielectric dipole.  $(06$  Marks) b.  $(06 Marks)$
- Define and explain the principle of pattern multiplication.  $\overline{c}$

### Module-5/





 $\mathcal{A} \times \mathcal{A} \times \mathcal{A}$ 



## **Scheme & Solution-VTU-QP**



# VTU Question Paper– Jan/Feb. 2022

As with the multicavity klystron, the operating mechanism is best understood by considering the behavior of individual electrons. This time, however, the reference electron is taken as one that passes the gap on its way to the repeller at the time when the gap voltage is zero and going negative. This electron is of course unaffected, overshoots the gap, and is ultimately returned to it, having penetrated some distance into the repeller space. An electron passing the gap slightly earlier would have encountered a slightly positive voltage at the gap. The resulting acceleration would have propelled this electron slightly farther into the repeller space, and the electron would thus have taken a slightly longer time than the reference electron to return to the gap. Similarly, an electron passing the gap a little after the reference electron will encounter a slightly negative voltage. The resulting retardation will,shorten its stay in the repeller space. It is seen that, around the reference electron, earlier electrons take longer to return to the gap than later electrons, and so the conditions are right for bunching to take place. The situation can be verified experimentally by throwing a series of stones upward. If the earlier stones are thrown harder, i.e., accelerated more than the later ones, it is possible for all of them to come back to earth simultaneously, i.e., in a bunch.

It is thus seen that, as in the multicavity klystron, velocity modulation is converted to current modulation in the repeller space, and one bunch is formed per cycle of oscillations. It should be mentioned that bunching is not nearly as complete in this case, and so the Reflex Klystron Oscillator is much less efficient than the multicavity klystron.



**2 (c)** A Smith chart is **developed by examining the load where the impedance must be matched**. Instead of considering its impedance directly, you express its reflection coefficient L, which is used to characterize a load (such as admittance, gain, and trans conductance). The L is more useful when dealing with RF frequencies. [06] CO2 L2





 $L1,$ 

 $L2$ 

In a microwave integrated circuit a strip line can be easily fabricated on a dielectric substrate by using printed-circuit techniques. A parallel stripline is similar to a two-conductor transmission line, so it can support a quasi-TEM mode. Consider a TEM-mode wave propagating in the positive z direction in a lossless strip line  $(R = G = 0)$ . The electric field is in the y direction, and the magnetic field is in the x direction. If the width  $w$  is much larger than the separation distance  $d$ , the fringing capacitance is negligible. Thus the equation for the inductance along the two conducting strips can be written as

$$
L = \frac{\mu_c d}{w} \qquad H/m \tag{11-2-1}
$$

where  $\mu_c$  is the permeability of the conductor. The capacitance between the two conducting strips can be expressed as

$$
C = \frac{\epsilon_d w}{d} \qquad F/m \qquad (11-2-2)
$$

where  $\epsilon_d$  is the permittivity of the dielectric slab.

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If the two parallel strips have some surface resistance and the dielectric substrate has some shunt conductance, however, the parallel stripline would have some losses. The series resistance for both strips is given by

$$
R = \frac{2R_s}{w} = \frac{2}{w} \sqrt{\frac{\pi f \mu_c}{\sigma_c}} \qquad \Omega/m \qquad (11-2-3)
$$

where  $R_s = \sqrt{\frac{\pi f \mu_c}{\sigma_c}}$  is the conductor surface resistance in  $\Omega$ /square and  $\sigma_c$  is the conductor conductivity in  $\mathbb{U}/m$ . The shunt conductance of the strip line is

$$
G = \frac{\sigma_d w}{d} \qquad \qquad \text{U/m} \tag{11-2-4}
$$

where  $\sigma_d$  is the conductivity of the dielectric substrate.



where  $e_t$  is the radiation efficiency of the transmitting antenna. For a nonisotropic transmitting antenna, the power density of (2-113) in the direction  $\theta_t$ ,  $\phi_t$  can be written as

$$
W_t = \frac{P_t G_t(\theta_t, \phi_t)}{4\pi R^2} = e_t \frac{P_t D_t(\theta_t, \phi_t)}{4\pi R^2}
$$
(2-114)

where  $G_t(\theta_t, \phi_t)$  is the gain and  $D_t(\theta_t, \phi_t)$  is the directivity of the transmitting antenna in the direction  $\theta_t$ ,  $\phi_t$ . Since the effective area  $A_t$  of the receiving antenna is related to its efficiency  $e_r$  and directivity  $D_r$  by

$$
A_r = e_r D_r(\theta_r, \phi_r) \left(\frac{\lambda^2}{4\pi}\right) \tag{2-115}
$$

the amount of power  $P_r$  collected by the receiving antenna can be written, using  $(2-114)$ and  $(2-115)$ , as

$$
P_r = e_r D_r(\theta_r, \phi_r) \frac{\lambda^2}{4\pi} W_t = e_t e_r \frac{\lambda^2 D_t(\theta_t, \phi_t) D_r(\theta_r, \phi_r) P_t}{(4\pi R)^2} |\hat{\rho}_t \cdot \hat{\rho}_r|^2 \tag{2-116}
$$

or the ratio of the received to the input power as

$$
\frac{P_r}{P_t} = e_t e_r \frac{\lambda^2 D_t(\theta_t, \phi_t) D_r(\theta_r, \phi_r)}{(4\pi R)^2}
$$
\n(2-117)

The power received based on (2-117) assumes that the transmitting and receiving antennas are matched to their respective lines or loads (reflection efficiencies are unity) and the polarization of the receiving antenna is polarization-matched to the impinging wave (polarization loss factor and polarization efficiency are unity). If these two factors are also included, then the ratio of the received to the input power of (2-117) is represented by

$$
\frac{P_r}{P_t} = e_{cdt}e_{cdr}(1-|\Gamma_t|^2)(1-|\Gamma_r|^2)\left(\frac{\lambda}{4\pi R}\right)^2 D_t(\theta_t, \phi_t)D_r(\theta_r, \phi_r)|\hat{\mathbf{p}}_t \cdot \hat{\mathbf{p}}_r|^2
$$
\n(2-118)

For reflection and polarization-matched antennas aligned for maximum directional radiation and reception, (2-118) reduces to

$$
\frac{P_r}{P_t} = \left(\frac{\lambda}{4\pi R}\right)^2 G_{0t} G_{0r}
$$
\n(2-119)

