

$$
\Delta P = J_1 \Delta \delta = \left[\frac{\partial P}{\partial \delta}\right] \Delta \delta \tag{6.69}
$$

$$
\Delta Q = J_4 \Delta |V| = \left[\frac{\partial Q}{\partial |V|}\right] \Delta |V| \tag{6.70}
$$

 (6.69) and (6.70) show that the matrix equation is separated into two decoupled equations requiring considerably less time to solve compared to the time required for the solution of (6.54). Furthermore, considerable simplification can be made to eliminate the need for recomputing J_1 and J_4 during each iteration. This procedure results in the decoupled power flow equations developed by Stott and Alsac[75-76]. The diagonal elements of J_1 described by (6.55) may be written as

$$
\frac{\partial P_i}{\partial \delta_i} = \sum_{j=1}^n |V_i||V_j||Y_{ij}|\sin(\theta_{ij} - \delta_i + \delta_j) - |V_i|^2|Y_{ii}|\sin\theta_{ii}
$$

Replacing the first term of the above equation with $-Q_i$, as given by (6.53), results in

$$
\frac{\partial P_i}{\partial \delta_i} = -Q_i - |V_i|^2 |Y_{ii}| \sin \theta_{ii}
$$

$$
= -Q_i - |V_i|^2 B_{ii}
$$

Where $B_{ii} = |Y_{ii}| \sin \theta_{ii}$ is the imaginary part of the diagonal elements of the bus admittance matrix. B_{ii} is the sum of susceptances of all the elements incident to bus *i*. In a typical power system, the self-susceptance $B_{ii} \gg Q_i$, and we may neglect Q_i . Further simplification is obtained by assuming $|V_i|^2 \approx |V_i|$, which yields

$$
\frac{\partial P_i}{\partial \delta_i} = -|V_i| B_{ii} \tag{6.71}
$$

Under normal operating conditions, $\delta_j - \delta_i$ is quite small. Thus, in (6.56) assuming $\theta_{ii} - \delta_i + \delta_j \approx \theta_{ii}$, the off-diagonal elements of J_1 becomes

$$
\frac{\partial P_i}{\partial \delta_j} = -|V_i||V_j|B_{ij}
$$

Further simplification is obtained by assuming $|V_j| \approx 1$

$$
\frac{\partial P_i}{\partial \delta_j} = -|V_i| B_{ij} \tag{6.72}
$$

Similarly, the diagonal elements of J_4 described by (6.61) may be written as

$$
\frac{\partial Q_i}{\partial |V_i|} = -|V_i||Y_{ii}|\sin\theta_{ii} - \sum_{j=1}^n |V_i||V_j||Y_{ij}|\sin(\theta_{ij} - \delta_i + \delta_j)
$$

replacing the second term of the above equation with $-Q_i$, as given by (6.53), results in

$$
\frac{\partial Q_i}{\partial |V_i|} = -|V_i||Y_{ii}|\sin\theta_{ii} + Q_i
$$

Again, since $B_{ii} = Y_{ii} \sin \theta_{ii} \gg Q_i$, Q_i may be neglected and (6.61) reduces to

$$
\frac{\partial Q_i}{\partial |V_i|} = -|V_i| B_{ii} \tag{6.73}
$$

Likewise in (6.62), assuming $\theta_{ij} - \delta_i + \delta_j \approx \theta_{ij}$ yields

$$
\frac{\partial Q_i}{\partial |V_j|} = -|V_i| B_{ij} \tag{6.74}
$$

With these assumptions, equations (6.69) and (6.70) take the following form

$$
\frac{\Delta P}{|V_i|} = -B' \,\Delta \delta \tag{6.75}
$$

$$
\frac{\Delta Q}{|V_i|} = -B'' \Delta |V| \tag{6.76}
$$

Here, B' and B'' are the imaginary part of the bus admittance matrix Y_{bus} . Since the elements of this matrix are constant, they need to be triangularized and inverted only once at the beginning of the iteration. B' is of order of $(n - 1)$. For voltage-controlled buses where $|V_i|$ and P_i are specified and Q_i is not specified, the corresponding row and column of Y_{bus} are eliminated. Thus, B'' is of order of $(n-1-m)$, where m is the number of voltage-regulated buses. Therefore, in the fast decoupled power flow algorithm, the successive voltage magnitude and phase angle changes are

$$
\Delta \delta = -[B']^{-1} \frac{\Delta P}{|V|} \tag{6.77}
$$

$$
\Delta|V| = -[B'']^{-1} \frac{\Delta Q}{|V|} \tag{6.78}
$$

The fast decoupled power flow solution requires more iterations than the Newton-Raphson method, but requires considerably less time per iteration, and a power flow solution is obtained very rapidly. This technique is very useful in contingency analysis where numerous outages are to be simulated or a power flow solution is required for on-line control.

The value of $Q_3^{(1)}$ is used as Q_3^{sch} for the computation of voltage at bus 3. The complex voltage at bus 3, denoted by $V_{c3}^{(1)}$, is calculated

$$
V_{c3}^{(1)} = \frac{\frac{P_3^{2\alpha} - j Q_3^{2\alpha}}{V_3^{(6)}} + y_{13} V_1 + y_{23} V_2^{(1)}}{y_{13} + y_{23}}
$$

=
$$
\frac{\frac{2.0 - j1.16}{1.04 - j0} + (10 - j30)(1.05 + j0) + (16 - j32)(0.97462 - j0.042307)}{(26 - j62)}
$$

= 1.03783 - j0.005170

Since |V₃| is held constant at 1.04 pu, only the imaginary part of $V_{c3}^{(1)}$ is retained, i.e, $f_3^{(1)} = -0.005170$, and its real part is obtained from

$$
e_3^{(1)} = \sqrt{(1.04)^2 - (0.005170)^2} = 1.039987
$$

Thus

$$
V_3^{(1)} = 1.039987 - j0.005170
$$

For the second iteration, we have

$$
V_2^{(2)} = \frac{\frac{P_2^{sch}-jQ_2^{sch}}{V_2^{(1)}} + y_{12}V_1 + y_{23}V_3^{(1)}}{y_{12} + y_{23}}
$$

\n
$$
= \frac{\frac{-4.0+j2.5}{37462+j042307} + (10 - j20)(1.05) + (16 - j32)(1.039987 + j0.005170)}{(26 - j52)}
$$

\n= 0.971057 - j0.043432
\n
$$
Q_3^{(2)} = -\Im\{V_3^{*(1)}[V_3^{(1)}(y_{13} + y_{23}) - y_{13}V_1 - y_{23}V_2^{(2)}]\}
$$

\n
$$
= -\Im\{(1.039987 + j0.005170)[(1.039987 - j0.005170)(26 - j62) - (10 - j30)(1.05 + j0) - (16 - j32)(0.971057 - j0.043432)]\}
$$

\n= 1.38796
\n
$$
V_{c3}^{(2)} = \frac{\frac{P_3^{sch}-jQ_3^{sch}}{V_3^{(1)}} + y_{13}V_1 + y_{23}V_2^{(2)}}{y_{13} + y_{23}}
$$

\n
$$
= \frac{\frac{2.9 - j1.38796}{1.039987 + j0.00517} + (10 - j30)(1.05) + (16 - j32)(.971057 - j.043432)}{(26 - j62)}
$$

\n= 1.03908 - j0.00730
\n(26 - j62)

Since $|V_3|$ is held constant at 1.04 pu, only the imaginary part of $V_{c3}^{(2)}$ is retained, i.e, $f_3^{(2)} = -0.00730$, and its real part is obtained from

$$
e_3^{(2)} = \sqrt{(1.04)^2 - (0.00730)^2} = 1.039974
$$

$$
V_3^{(2)} = 1.039974 - j0.00730
$$

The process is continued and a solution is converged with an accuracy of 5×10^{-5} pu in seven iterations as given below.

The final solution is

 α r

 $V_2=0.97168\angle{-2.6948}^{\circ}$ pu

 $S_3 = 2.0 + j1.4617$ pu $V_3 = 1.04\angle -0.498$ ° pu $S_1 = 2.1842 + i1.4085$ pu

Line flows and line losses are computed as in Example 6.7, and the results expressed in MW and Mvar are

 $\begin{aligned} S_{12} &= 179.36 + j118.734 \quad S_{21} = -170.97 - j101.947 \quad S_{L\,12} = 8.39 + j16.79 \\ S_{13} &= 39.06 + j22.118 \quad \quad S_{31} = \, -38.88 - j \,\, 21.569 \quad S_{L\,13} = 0.18 + j0.548 \end{aligned}$ $S_{23} = -229.03 - j148.05$ $S_{32} = 238.88 + j167.746$ $S_{L23} = 9.85 + j19.69$

The power flow diagram is shown in Figure 6.13, where real power direction is indicated by \rightarrow and the reactive power direction is indicated by \mapsto . The values within parentheses are the real and reactive losses in the line.

 $\frac{\partial Q_2}{\partial \mathfrak{c}} = |V_2||V_1||Y_{21}|\cos(\theta_{21}-\delta_2+\delta_1)+|V_2||V_3||Y_{23}|$ $\overline{\partial \delta_2}$ $\cos(\theta_{23} - \delta_2 + \delta_3)$ ∂Q_2 $= -|V_2||V_3||Y_{23}|\cos(\theta_{23}-\delta_2+\delta_3)$ $\overline{\partial \delta_3}$ ∂Q_2 = $-|V_1||Y_{21}|\sin(\theta_{21}-\delta_2+\delta_1)-2|V_2||Y_{22}|\sin\theta_{22} \overline{\partial|V_2|}$ $|V_3||Y_{23}|\sin(\theta_{23}-\delta_2+\delta_3)$

The load and generation expressed in per units are

$$
S_2^{sch} = -\frac{(400 + j250)}{100} = -4.0 - j2.5 \quad \text{pu}
$$

$$
P_3^{sch} = \frac{200}{100} = 2.0 \quad \text{pu}
$$

The slack bus voltage is $V_1 = 1.05\angle 0$ pu, and the bus 3 voltage magnitude is $|V_3| = 1.04$ pu. Starting with an initial estimate of $|V_2^{(0)}| = 1.0$, $\delta_2^{(0)} = 0.0$, and $\delta_3^{(0)} = 0.0$, the power residuals are computed from (6.63) and (6.64)

$$
\Delta P_2^{(0)} = P_2^{sch} - P_2^{(0)} = -4.0 - (-1.14) = -2.8600
$$

\n
$$
\Delta P_3^{(0)} = P_3^{sch} - P_3^{(0)} = 2.0 - (0.5616) = 1.4384
$$

\n
$$
\Delta Q_2^{(0)} = Q_2^{sch} - Q_2^{(0)} = -2.5 - (-2.28) = -0.2200
$$

Evaluating the elements of the Jacobian matrix with the initial estimate, the set of linear equations in the first iteration becomes

Obtaining the solution of the above matrix equation, the new bus voltages in the first iteration are

Voltage phase angles are in radians. For the second iteration, we have

and

and

The solution converges in 3 iterations with a maximum power mismatch of 2.5×10^{-4} with $V_2 = 0.97168\angle -2.696^\circ$ and $V_3 = 1.04\angle -0.4988^\circ$. From (6.52) and (6.53), the expressions for reactive power at bus 3 and the s tive powers are

$$
Q_3 = -|V_3||V_1||Y_{31}|\sin(\theta_{31} - \delta_3 + \delta_1) - |V_3||V_2||Y_{32}|
$$

$$
\sin(\theta_{32} - \delta_3 + \delta_2) - |V_3|^2|Y_{33}|\sin \theta_{33}
$$

 $P_1 = |V_1|^2 |Y_{11}| \cos \theta_{11} + |V_1||V_2||Y_{12}| \cos(\theta_{12} - \delta_1 + \delta_2) + |V_1||V_3|$ $|Y_{13}|\cos(\theta_{13}-\delta_1+\delta_3)$

$$
Q_1 = -|V_1|^2|Y_{11}|\sin\theta_{11} - |V_1||V_2||Y_{12}|\sin(\theta_{12} - \delta_1 + \delta_2) - |V_1||V_3|
$$

\n
$$
|Y_{13}|\sin(\theta_{13} - \delta_1 + \delta_3)
$$

Upon substitution, we have

 $Q_1 = 1.4085$ pu

Finally, the line flows are calculated in the same manner as the line flow calculations in the Gauss-Seidel method described in Example 6.7, and the power flow

diagram is as shown in Figure 6.13.
A program named **linewton** is developed for power flow solution by the
Newton-Raphson method for practical power systems. This program must be precoded by the lippus program. busout and lineflow programs can be used to print
the load flow solution and the line flow results. The format is the same as the Gauss-Seidel. The following is a brief description of the Ifnewton program.

****** ALL THE BEST ******