

Scanned by CamScanner

ol. $disc$ α the \mathcal{C} r. **AHO** Dno Athere cha rhod par col wes them is usually standard orwer source α Curre natin Ċ4 ow-voltage α the U voltare rom the ہ
عو standard 480 V V σv $7 - 6$ $\frac{1}{2}$ $20,000$ 6 $220 - 6$ $#$ t allornating set Conver Reeli w direc $d \cdot c$ to $\overline{(\begin{smallmatrix} a,c \end{smallmatrix})}$ current sectionalization? What is field. electrodes are avranged \sqrt{m} discharge Ike own T-R se ω its lead provered partitioning Known a electrical This \int_{0}^{∞} because use brumar ization. This sectiona requiremen α van bower within $\sqrt{2}$ $mu \alpha$ 忧 nee 20000 more $Im(\sigma m)$ can tration $\bigcap_{\mathbf{F}}$ \mathbf{v} cm \mathcal{C} ρσωσ cust g_u ren Con $\frac{1}{10}$ at the Se als σ neo Jur $i\delta i$ that ϵ Smal $_{\text{max}}$ \mathscr{L} α particl $+0$ renstance Ò \overline{O}

Scanned by CamScanner

Date: it is not easy installed once $i.e.$ O. another places mó particles with bh 20150 curve q_{\bullet}^{σ} 8O $\overline{90}$ 60 50° э $200, 300, 400$ 600 500 100 ô c_{f^m} 1000 ratio W power corona ī the $anE6P$ nchim σ α Ίų pica 900 coron Watt $\frac{1}{1}$ Power Consume mi m Ξ per der 不错

1b Electrostatis separator

Separation Electrostatic u se $\overline{2}$ m λb rr mha Ċτ m ω ky the are generated Electrics 'c tonces $\frac{1}{2}$ parti'd rric hara ٥٦ α n ES Source. nee α proce one π $potential$ to genera q, which H_{α} ono ces electrical vidual hYr arti i md are ċ tactors ha \mathcal{H}_{α} can \mathbf{c}_n proug Intern 1. decki i'ild. σ $2.$ Particle size. $\frac{3}{2}$. Relative humidit $4.$ Perpocative H_{c} Particle bed. 5° Inter Brode ۱. tance. Advantages. Electrostatio t. the particles onces wor $6n$ **ATTACHED** Scanned by CamScanner

Date: $1¹$ they do not to be separated only athech $0 + Re$ medium. in particle the **DAL** $209. h$ 0 llow electri m fke particu p lication Sha E.S. forces rection σ may h0 $\overline{\mathbf{3}}$. the $holaxi$ changing J_{0} eithe σ α_{14} the 0 dencetion $4k$ ا ہے σ arranged $C.S$ $4.$ $oneo$ mai work \overline{m} Compinent with other Lonced λu \mathfrak{g} gravitational 0^h C_{44+} oher E.S separation 5° the. α inde the $\mathcal{S}u$ بہ oи **SLU** Lace generate do bo in n electric $0.2.4$ 0 α charge Ma anta imitation H_a marks mass max. can \sim tively work $\n *upon*$. s_1 $2e$ the aualcol makerial ô., \pm $a₀$ ۵ the which sma m m_l Lonution Como \mathfrak{a} ľ σ Comme nution He ducing action 900 σ material dur i especially a mineral ore to minute. Scanned by CamScanner

 $cons_s$ σ a drum mad This 29 othe S ome ϵ_{cm} ٥H m_{a} throu λ Comprusion 280 ä Ņ ohr. sub $_{\text{both}}$ S C_{LITI}te bans km an μ o D 2_u 5 kV_A О -11 0 忧 Volta 'n t ai_s $10₀$ k ۱Q λч λ risible املی Corona cing ~ 0 b الہ \mathcal{H}_\bullet Ω¥ voide des Thous $10⁷$ n_{ω} α sation α **CCCLIBLE** 2000 ve ıЗ $10m$ the citys lee ives Ο \overline{a} Con a **CD** 忧 atts b ŀ٥ Cauring ٦o **as** H_{α} λ we QCO roton $2a$ amon m Come $\overline{11}$ VΤ $91G$ Sim \mathfrak{c} Ωr $#$ わ rab'll oas $s31$ pat σ nese Q Qa σ tor Λ Con œ hνi her ¢ α bination pinning $_{\text{C}\sigma m}$ h unn created b o kа \mathcal{O} enot $\sqrt{2}$ the trode Corona ä, As est A considerable for Scanned by CamScanne

which hing ru ticles are The Conduc hom Н. 200 ö. λ LK.O Lom am $+t$ ound $\cot \alpha$ α i ρ_{cm} duchno α \mathcal{S} e m_{6} لمزلها Plate separtion ele 之. h'c unit " trosta 0 s_{e} parator 01 09 Com α \cdot 9100 ô $#_o$ hro Con pay ý. Ab $\frac{1}{2}$ \mathbf{a} oara h 'on wo $\overline{10}$ 100 wide ò rang min 0 \mathcal{L} Seac San kere 2.5 Bene , leah's ø oal Bene \Im ning ind Ω ΨĄ ex tall the rou m \mathcal{P} proves ene Ore 枨 t moving material. 9 que Scanned by Cam!

Poursend's uvoien.
Ref to fig (), let us assume that no les are Ref to fig (), let us assume .
emitted from Cathode: When one e collide emitted from Cathodi. Who i'm and an e are! This is called lone'ring collision tr het a be aug no. of ionising collinion made by an e/con travel in the direction of the fld (x depends on gas pre p and E/P. and called T's first ionication coeff). Ata distance a from cathode, let no. of clecho be no when there me es travel a further distance of da, they give reserts annular & At $x = 0$, $n_x = n_0$ Also $\frac{dn_2}{dx}$ = λn_2 $ln n_2$ = n_0 exp (x) $3heonoofe;teachingauroole (z=d)
3heonoerb(a))$ Md= no exp (x) - (3) Namber of new es créated, on aug by each e in $\frac{m_{d-m_0}}{n_0}$ = exp(xd) - 1

2a

.' Avg current in gap, which is equal to no. of es travelling per second will be. where I_0 is initial current at the cathode. Current Growth In The Presence of 2ndary Processes Single avalanche process becomes complete when initial sets les reaches anode. Since amphification of es [exp (od) is occuring in flo, probability of additional new es being liberated in gapty further avalanches. Other mechanisme are) the low liberated may have sufficient energy to cause liberations of es from the cathode when they impinge on it if Ecuted mots oratom in avalanches may units photo-envision. Eu) Metastable particle may diffuse back cousing e envission.

 $I = I_0 exp(\alpha d)$ $(1 - 3) \int exp[\alpha d) - 1.$ Pownsend s Criterian for Breakdown Above egn. gives total aug current in a gap before occuraire of breakdown. As d?, denominates of agn tends to zero and at gome vitical distance $dzds = 1 - \gamma |exp(\alpha d)^{-1}| = 0$. fon values d (ds $I \cong I_0$ and external source for reply of Io is removed I becomes zero. It de des IX and current will be liméted only by curso of fower supply and ext ckt. This condition les called M's breakdours oriterion and can be written as γ [exp (xd)-1] = 1 Normally esp (x d) in very large andabore equi reduces to $\sqrt{8\exp(\alpha d)} = 1$ f

e sproduced by these enussion. A 2ndary ionization coeff & is defined as net number of 2ndary es produced by per incident tre ions photon, éxcited particle on metastable particle and total value of 8 is num of individual coeffs due to 3 diff processes. $y' = y_1 + y_2 + y_3$ V - T's 2ndary ionization coeff and function of P and E/P Following T's procedure of wovent growth let us assume, no! = no. of 2ndary es produced due to 2ndary processes no" - total no of es leaving cathode $m_0^{\prime\prime} = m_0 + m_0^{\prime}$ To tal no. of es n reaching ande become $n = n_0$ "exploid = $(m_0 + n_0)$ exploid and $n_0' = 3\sqrt{n} + (n_0 + n_0)$ Eliminating no n= noexp (ad) $exb(d)$

2b **Limitation In Townsend's Mechanism**

• Townsend mechanism when applied to breakdown at atmospheric pressure was found to have certain drawbacks.

• A) according to the Townsend theory, current growth

occurs as a result of ionization processes only. But in practice, breakdown voltages were found to depend on the gas pressure and the geometry of the gap.

B) Mechanism predicts time lags of the order of 10-5S, while in actual practice breakdown was observed to occur at very short times of the order of 10-8S.

C) Townsend mechanism predicts a very diffused form of discharge, in actual practice, discharges were found to be filamentary and irregular.

The Townsend mechanism failed to explain all these observed phenomena and as a result, around 1940, Raether and, Meek and Loeb independently proposed the **Streamer theory**.

The County of the County Townserd's Criteria α ex \overline{O} -1 enables H_a evaluation. breakdown voltage the σ 296 va Ε p when $+\rho$ copous Curvice 。 cath. $\frac{1}{100}$ low ano α isfor $\frac{1}{10}$ Shace λ d are min! 1 $+$ وت μ ard α σ \overline{u} percine O Vales matches are $s l$ θ \overline{a} n \mathbf{a} ong pressure ativel 910 oue. 0 θ Breakdron $V_{\mathcal{O}}$ IJ $m+n$ tion a_o Q q oren d \mathcal{O} threshold ね hopid $a \sqrt{m}$ \bullet ä 日期 θ 4 0.15874 u. -14.8 1.70 \bar{q}^i 46 \mathbb{E} 0 e ϵ Pd Ξ n'ales : b oth Kim Scanned by CamScanner 3

 \times $\overline{1}$ $V_{\mathbf{b}}$ E_{2} $10₁$ W \overline{d} $+$ o ϵ b d N \mathbf{b} $P_{\mathbf{y}}$ pd p A his Show breakdown Le VO $+$ cn б. a V \overline{M} hm qα unige o đ ine h, $\overline{\sigma}$ pood ϵ eletro de ress wy α_1 d und. V \overline{a} w aschen's au . $h_{\mathfrak{m}}$ 9τ, dos mot mean Tha $\overline{10}$ α JH.R Pxo $Domh$ n e Some 0 אפ ۹ **107** rela hon Prod å Ѥ near. \sim

Scanned by CamScanner

Scanned by CamScanner

This is equation to a straight with ine the value uron obe l_{12} \vert Oma Æ v_0 to curve $\frac{2c}{2c}$ α cui this \mathcal{B} A and $bosh 1$. × 100 . preakd voltages a $+w$ BLUY he mus there Correce Hо mg resure one open nQ O b om $#$ smal leng $1 - 2$ ler **ga** longer other \mathbf{b} o'm I Ł 0.91 the Ŝ Ō. σ N olta $A +$ $hr1.0$ straight and curve! ine intersection between 900 22 with break down occurs This means MO Pasches's minimum Small $\hbar \omega$ vottages va $curve¹$ the indicator $\vert \hspace{-0.2em} \vert_{\mathsf{nm}} \hspace{-0.2em} \vert$ the CLUMY Δn 3 parken mum price ϵ wn Voltak on α Potenhal boon concespondin Fo S par Nolting oven previou $|m|$ shown Λ ari

Scanned by CamScanne

uabergrid Date: plain He \circ ex there \mathcal{I} α c exist menigner spar potemn $Im \tfrac{1}{100}$ Ween \mathfrak{S} \sim \mathfrak{m} an assumeng com ί'n θ Can đ ϵ α P. Completerio one ration α õ thavern conff c energics. Ssuming ownsen Second onization Coe his i $\sqrt{2}$ ma Va ひん trom Æ Crossin α islor Δ ay n. $a \wedge$ molecu p d mm's en ucconsive ø that m α P_o min σ_{ba} ence 0 meg $0₂$ 500 $41n$ tto case J. pd σ P_q min Cross 0.4 their Λ and 积 sparke ł m 0σ Point Pd Correspondin mem 10 1 om pation Einey ϱ minimus an hence α potembo **ABS - BOOTS - 277-28**

Scanned by CamScanner

expression minimum A_0 sea ling obtained general unng can d p renio U. Ε Ae X Έ b PA ₹ E B pd Ξ $\overline{\rho}$ $e^{\beta \rho d}$ Þ $\dddot{}$ \overline{P} A </u> Bpd λ . α d PA . Xd ∞ Bpd $\overline{\sqrt{b}}$ ϱ ĭ, $\overline{\sigma}$ \rightarrow $\ddot{}$ Ė h ₽ PA Compan a° \circ $\mathcal{V}(\mathcal{L})$ ln $\overline{+}$ n Bp $\sqrt{2}$ $\frac{1}{d}$ VЬ d \overline{z} $\overline{\rho}$ Χ $P A$

Scanned by CamScanner

Scanned by CamScanner

4.a **THERMAL BREAKDOWN**

- The breakdown voltage of a solid dielectric should increase with its thickness.
- This is true only up to a certain thickness above which the heat generated in the dielectric due to the flow of current determines the conduction.
- When an electric field is applied to a dielectric, conduction current, however small it may be, flows through the material.
- The current heats up the specimen and the temperature rises.
- The heat generated is transferred to the surrounding medium by conduction through the solid dielectric and by radiation from its outer surfaces.
- Equilibrium is reached when the heat used to raise the temperature of the dielectric, plus the heat radiated out, equals the heat generated. The heat generated under d.c. stress *E is given as*

$$
W_{\text{d.c.}} = E^2 \sigma \qquad W/\text{cm}^3
$$

- where, б is the d.c. conductivity of the specimen.
- Under a.c. fields, the heat generated

$$
W_{\text{a.c.}} = \frac{E^2 f \epsilon_r \tan \delta}{1.8 \times 10^{12}} \qquad W/cm^3
$$

• where, *f*= frequency in Hz,

•

• δ= loss angle of the dielectric material, and

• *E- rms value.*

•

• The heat dissipated *(W^T)is given by*

$$
W_T = C_V \frac{dT}{dt} + \text{div} (K \text{ grad } T)
$$

- where, C_V = specific heat of the specimen,
- T= temperature of the specimen,
- *K = thermal conductivity of the specimen, and*
- \cdot $t =$ time over which the heat is dissipated.
- Equilibrium is reached when the heat generated (Wa.c.or Wd.c.) becomes equal to Equilibrium is reached when the heat generated Wa.c.or Wd.c becomes equal to the heat dissipated *(WT).*
- *In actual practice there is always some heat that is radiated* out*).*
- *Breakdown occurs when Wd.c. or Wa.c. exceeds WT.*
- *The thermal instability condition is shown in Fig.*
- *Here, the heat lost is shown by a straight line*
- *the heat generated at fields E1 and E2 are shown by separate curves.*
- *At field EI breakdown occurs both at temperatures TA and TB*
- *In the temperature region of TA and TB heat generated is less than the heat lost for the field E2 the breakdown will not occur.*
- *This is of great importance to practising engineers, as most of the insulation failures in high voltage power apparatus occur due to thermal breakdown.*
- *Thermal breakdown sets up an upper limit for increasing the breakdown voltage when the thickness of the insulation is increased.*
- *For a given loss angle and applied stress, the heat generated is proportional to the frequency and hence thermal breakdown is more serious at high frequencies.*

It can be seen

•

• since the power loss under a.c. fields is higher, the heat generation is also high, and hence the thermal breakdown stresses are lower under a.c. conditions than under d.c. conditions.

4b **ELECTROMECHANICAL BREAKDOWN**

When solid dielectrics are subjected to high electric fields, failure occurs due

to electrostatic compressive forces which can exceed the mechanical compressive strength.

If the thickness of the specimen is *4\$ and is compressed to* a thickness *d under an applied voltage V, then the electrically developed compressive stress is in equilibrium*

$$
\varepsilon_0 \, \varepsilon_r \, \frac{V^2}{2d^2} = Y \, \ln \left[\frac{d_0}{d} \right]
$$

where *Y is the Young's modulus.*

Usually, mechanical instability occurs when

$$
V^2 = d^2 \left[\frac{2Y}{\epsilon_0 \epsilon_r} \right] \ln \left[\frac{d_0}{d} \right]
$$

Substituting this in Equation the highest apparent electric stress before breakdown,

d/d_0 = 0.6 or d_0/d = 1.67

- The above equation is only approximate as *Y depends on the mechanical stress.*
- When the material is subjected to high stresses the theory of elasticity does not hold good, and plastic deformation has to be considered

$$
E_{\text{max}} = \frac{V}{d_0} = 0.6 \left[\frac{Y}{\epsilon_0 \epsilon_r} \right]^{\frac{1}{2}}
$$

5. **Suspended Particle Theory**

- In commercial liquids, the presence of solid impurities cannot be avoided.
- These impurities will be present as fibres or as dispersed solid particles.
- The permittivity of these particles (**ε1**) will be different from the permittivity of the liquid **(ε2**).
- If we consider these impurities to be spherical particles of radius r, and if the applied field is E, then the particles experience a force F, where

$$
F = r^3 \frac{\varepsilon_1 - \varepsilon_2}{\varepsilon_1 + 2\varepsilon_2} E \cdot \frac{dE}{dx}
$$

- and this force is directed towards a place of higher stress if ϵ 1 > ϵ 2 and towards a place of lower stress if $ε1 < ε2$ when $ε1$ is the permittivity of gas bubbles.
- The force given above increases as the permittivity of the suspended particles (ϵ 1) increases. If ε1 → ∞

$$
F = r^3 \frac{1 - \varepsilon_2/\varepsilon_1}{1 + 2\varepsilon_2/\varepsilon_1} E \frac{dE}{dx}
$$

Let $\varepsilon_1 \to \infty$

$$
F = r^3 E \cdot \frac{dE}{dx}
$$

- Force will tend the particle to move towards the strongest region of the field.
- In a uniform electric field which usually can be developed by a small sphere gap, the field is the strongest in the uniform field region**. Here** *dE/dx → 0 so that the force on the particle is zero and the particle* **remains in equilibrium**.
- Particles will be dragged into the uniform field region.
- Permittivity of the particles is higher than that of the liquid, the presence of particle in the uniform field region will cause flux concentration at its surface.
- Other particles if present will be attracted towards the higher flux concentration.
- The movement of the particle under the influence of electric field is opposed by the viscous force posed by the liquid and since the particles are moving into the region of high stress, diffusion must also be taken into account.
- We know that the viscous force is given by (Stoke's relation) *FV = 6πnrν where* η is the viscosity of liquid, *r the raidus of the particle and v the velocity of the particle.*
- Equating the electrical force with the viscous force we have

However, if the diffusion process is included, the drift velocity due to diffusion will be given by

$$
v_d = -\frac{D}{N}\frac{dN}{dx} = -\frac{KT}{6\pi\eta r}\frac{dN}{Ndx}
$$

where *D = KT/6πηr a relation known as Stokes-Einstein relation. Here K is Boltzmann's constant and T the absolute temperature.*

At any instant of time, the particle should have one velocity and, therefore, equation

$$
v = vd
$$

We have

$$
-\frac{KT}{6\pi\eta r} \cdot \frac{dN}{Ndx} = \frac{r^2 E}{6\pi\eta} \cdot \frac{dE}{dx}
$$

$$
\frac{KT}{r} \frac{dN}{N} = -r^2 E dE
$$

$$
\frac{KT}{r} \ln N = -\frac{r^2 E^2}{2}
$$

• It is clear that the breakdown strength *E*

depends upon the concentration of particles N, radius r of particle, viscosity η of liquid and temperature *T of the liquid.*

- It has been found that liquid with solid impurities has lower dielectric strength as compared to its pure form.
- larger the size of the particles impurity the lower the overall dielectric strength of the liquid containing the impurity.
- If there is only a single conducting particle between the electrodes, it will give rise to local field enhancement depending on its shape.
- If this field exceeds the breakdown strength of the liquid, local breakdown will occur near the particle, and this will result

in the formation of gas bubbles which may lead to the breakdown of the liquid.

- The values of the breakdown strength of liquids containing solid impurities was found to be much less than the values for pure liquids.
- The impurity particles reduce the breakdown strength, and it was also observed that the larger the size of the particles the lower were the breakdown strengths
- **Cavitation and the Bubble Theory**
- experimentally observed that in many liquids, the breakdown strength depends

strongly on the applied **hydrostatic pressure.**

suggesting that a change of phase of the medium is involved in the breakdown process, which means kind of **vapour bubble** formed is responsible for breakdown.

The following processes are responsible for **formation of the vapour bubbles**:

- *(a)* **Gas pockets at the surfaces of the electrodes;**
- *(b)* **electrostatic repulsive forces between space charges which may be sufficient to overcome the surface tension**;
- (c) **gaseous products due to the dissociation of liquid molecules by electron collisions**
- *(d)* **vaporizations of the liquid by corona type discharge from sharp points and irregularities on the electrode surfaces**.
- Once a bubble is formed it will elongate in the direction of the electric field under the influence of electrostatic forces.
- The volume of the bubble remains constant during elongation.
- Breakdown occurs when the voltage drop along the length of the bubble becomes equal to the minimum value on the Paschen's curve
- The breakdown field is given as

- of the liquid,
- *ε 1 is the permittivity of the liquid,*
- *ε2 is* the permittivity of the gas bubble,
- *r is the initial radius of the bubble assumed as a* sphere
- *Vb is the voltage drop in the bubble (corresponding to minimum on the*
- Paschen's curve).
- From this equation, it can be seen that the breakdown strength depends on the initial size of the bubble which in turn is influenced by the hydrostatic pressure and temperature of the liquid.
- This theory does not take into account the production of the initial bubble and hence the results given by this theory do not agree well with the experimental results

6a

When a solid dielectric subjected to electrical stresses for a long time fails, normally two kinds of visible markings are observed on the dielectric materials. They are:

- (a) *the presence of a conducting path across the surface of the insulation;*
- *(b) a mechanism whereby leakage current passes through the conducting path*

finally leading to the formation of a spark.

Insulation deterioration occurs as a result of these sparks.

The spreading of spark channels during *tracking, in the form of the branches of a* tree is called **treeing***.*

- Consider a system of a solid dielectric having a conducting film and two electrodes on its surface.
- In practice, the conducting film very often is formed due to moisture.
- On application of voltage, the film starts conducting, resulting in generation of heat, and the surface starts becoming dry.
- The conducting film becomes separate due to drying, and so sparks are drawn damaging the dielectric surface.

With organic insulating materials such as paper and bakelite, the dielectric carbonizes at the region of sparking, and the carbonized regions act as permanent conducting channels resulting in increased stress over the rest of the region.

This is a cumulative process, and insulation failure occurs when carbonized tracks bridge the distance between the electrodes.

This phenomena, called **tracking** is common between layers of bakelite, paper and similar dielectrics built of laminates

- **Treeing** occurs due to the erosion of material at the tips of the spark.
- Erosion results in the roughening of the surfaces, and hence becomes a source of dirt and contamination.
- This causes increased conductivity resulting either in the formation of a conducting path bridging the electrodes or in a mechanical failure of the dielectric.

- **Arrangement for study of treeing phenomena. 1 and 2 are electrodes**
- When a dielectric material lies between two electrodes as shown in Fig. there is a possibility for two different dielectric media, the air and the dielectric, to come in series.
- The voltages across the two media are as shown *(V1 across the air gap, and V2* across the dielectric).
- The voltage *V1 across the air gap is given as,*

$$
V_1 = \frac{V d_1}{d_1 + \left(\frac{\varepsilon_1}{\varepsilon_2}\right) d_2}
$$

where *V is the applied voltage.*

- Since ε2>ε1, most of the voltage appears across **d1 the air gap.**
- *Sparking will* occur in the air gap and, charge accumulation takes place on the surface of the insulation.
- Sometimes the spark erodes the surface of the insulation.
- As time passes, breakdown channels spread through the insulation in an irregular "tree" like fashion leading to the formation of conducting channels.
- This kind of channelling is called **treeing**
- Usually, **tracking** occurs even at **very low voltages** of the order of about 100 V, whereas **treeing** requires **high voltage**.
- For testing of tracking, low and medium voltage tracking tests are specified.
- These tests are done at low voltages but for times of about 100 hr or more.
- The insulation should not fail.
- Sometimes the tests are done using 5 to 10 kV with shorter durations of 4 to 6 hr.
- The numerical value of voltage that initiates or causes the formation of a track is called the **"tracking index"** and this is used to qualify the surface properties of dielectric materials.
- Treeing can be prevented by having clean, dry, and undamaged surfaces and a clean environment.
- The materials chosen should be resistant to tracking.
- Sometimes moisture repellant greases are used.
- But this needs frequent cleaning and regreasing.
- Increasing creepage distances should prevent tracking, but in practice the presence of moisture films defeat the purpose.
- Usually, treeing phenomena is observed in capacitors and cables.

6b

 \bullet

- It has been recognized that one process that gives high breakdown strength to a gas is the electron attachment in which free electrons get attached to neutral atoms or molecules to form negative ions.
- Since negative ions like positive ions are too massive to produce ionization due to collisions, attachment represents an effective way of removing electrons which otherwise would have led to current growth and breakdown at low voltages.
- The gases in which attachment plays an active role are called **electronegative gases.**
- The most common attachment processes encountered in gases are
- *(a) the direct* attachment in which an electron directly attaches to form a negative ion,
- (b) the dissociative attachment in which the gas molecules split into their constituent atoms and the electronegative atom forms a negative ion.
- These processes may be symbolically represented as:
- *(a) Direct attachment*

• *(b) Dissociative attachment*

$AB + e \longrightarrow A + B^-$

•

•

•

•

- A simple gas of this type is oxygen. Other gases are sulphur hexafluoride, freon, carbon dioxide, and fluorocarbons.
- In these gases, 'A" is usually sulphur or carbon atom, and 'B* is oxygen atom or one of the halogen atoms or molecules.
- With such gases, the Townsend current growth equation is modified to include ionization and attachment
- . **An attachment coefficient (η) is defined, as the number of attaching collisions made by one electron drifting one centimetre in the direction of the field.**
- Under these conditions the current reaching the anode, can be written as:

$$
I = I_0 \frac{\left[\left(\alpha/(\alpha - \eta) \right) \exp \left(\alpha - \eta \right) d \right] - \left[\eta/(\alpha - \eta) \right]}{1 - \left\{ \gamma \frac{\alpha}{(\alpha - \eta)} \left[\left\{ \exp \left(\alpha - \eta \right) d \right\} - 1 \right] \right\}}
$$

• The Townsend breakdown criterion for attaching gases can also be deduced by equating the denominator to zero:

$$
\gamma \frac{\alpha}{(\alpha - \eta)} [\exp(\alpha - \eta)d - 1] = 1
$$

- This shows that for $\eta < \alpha$ breakdown is always possible irrespective of the values of α , η , γ
- If on the other hand, η > α *Eq approaches an asymptotic form with* increasing value of d,

$$
\gamma \frac{\alpha}{(\alpha - \eta)} = 1 \; ; \; \text{or} \; \alpha = \frac{\eta}{(1 - \gamma)}
$$

• This condition puts a limit *for E/p below which no breakdown is possible irrespective* of the value of *d, and the limit value is called the critical E/p. Critical E/p for SF6 is 117 V cm⁻¹ torr⁻¹,* and for CCl2F2 it is 121 V cm⁻¹ torr⁻¹ (both at 20⁰C).

• The mechanism of spark breakdown is considered as a function of ionization processes under uniform field conditions.

But in practical engineering designs, the breakdown due to rapidly changing voltages or impulse voltages is of great importance.

- Actually, there is a time difference between the application of a voltage sufficient to cause breakdown and the occurrence of breakdown itself.
- This time difference is called the **time lag.**
- The Townsend criterion for breakdown is satisfied, only if at least one election is present in the gap between the electrodes.
- In the case of applied d.c. or slowly varying (50 Hz a.c) voltages, there is no difficulty in satisfying this condition.
- With rapidly varying voltages of short duration (=10-6 *s), the initiatory electron may not* be present in the gap.
- In the absence of such an electron breakdown cannot occur.
- The time *t which lapses between the application of the voltage sufficient to cause* breakdown and the appearance of the initiating electron is called a **statistical time lag** (ts)of the gap.
- The appearance of electrons is usually statistically distributed.
- After the appearance of the electron, a time *tt is required for the ionization processes to* develop fully to cause the breakdown of the gap
- this time is called the **formative time lag** *(tt).*
- *The total time ts + tt=t is called the total time lag.*
- Time lags are of considerable practical importance.
- For breakdown to occur the applied voltage *V should be greater than the static breakdown voltage Vs as shown* in Fig.
- The difference in voltage ΔV = V- *Vs is called the overvoltage*
- ratio *V/Vs is called the impulse ratio.*

Breakdown on the front of the applied impulse voltage wave

• *The variation of tt with overvoltage (ΔV)* is shown in Fig.

7a *Withstand Voltage*

•

• The voltage which has to be applied to a test object under specified conditions in a **withstand test** is called the **withstand voltage**

7b *Hundred Per Cent Flashover Voltage*

The voltage that causes a flashover at each of its applications under specified conditions when applied to test objects is specified as **hundred per cent flashover voltage**.

7c *Fifty Per Cent Flashover Voltage*

- This is the voltage which has a probability of **50% flashover,** when applied to a test object.
- This is normally applied in impulse tests in which the loss of insulation strength is temporary

7d *Disruptive Discharge Voltage*

- This is defined as the voltage which produces the loss of dielectric strength of an insulation.
- It is that voltage at which the electrical stress in the insulation causes a failure which includes the collapse of voltage and passage of current.
- In solids, this causes a permanent loss of strength.
- In liquids or gases only temporary loss may be caused.
- When a discharge takes place between two electrodes in a gas or a liquid or over a solid surface in air, it is called **fiashover**.
- If the discharge occurs through a solid insulation it is called **puncture**.

7e *Creepage Distance*

It is the shortest distance on the contour of the external surface of the insulator unit or between two metal fittings on the insulator.

Solution: 1st Set. Since there is gradual increase in current upto gap distance of 3 mm, slope between any two points

$$
\left(\frac{\ln I/I_0}{x}\right)
$$

will give us the value of α .

Let us take gap distances of 2 and 2.5 mm.

The respective $\ln I/I_0$ are

$$
\ln\left(\frac{1.5 \times 10^{-12}}{6 \times 10^{-14}}\right) = 3.2188
$$

$$
\ln\left(\frac{1.5 \times 10^{-12}}{6 \times 10^{-14}}\right) = 4.5362
$$

$$
\therefore \text{ The slope } = \frac{4.5362 - 3.2188}{0.05} = 26.34
$$

and

Since there is sudden rise in current at the last observation, this is used to evaluate γ . We know that

0f

0ſ

 $I = \frac{I_0 e^{\alpha x}}{1 - \gamma (e^{\alpha x} - 1)}$ $\frac{I}{I_0} = \frac{7}{6} \times 10^7 = \frac{e^{26.34 \times 0.5}}{1 - \gamma(e^{13.17} - 1)}$

$$
= \frac{5.24 \times 10^5}{1 - 5.24 \times 10^5 \gamma}
$$

or

$$
\frac{7}{6} \times 10^7 \frac{1}{5.24 \times 10^5} = \frac{1}{1 - 5.24 \times 10^5 \gamma}
$$

or
0.0449 = 1 - 5.24 × 10⁵ γ
or
0.9551 = 5.24 × 10⁵ γ

or
$$
\gamma = 0.182 \times 10^{-5} / \text{cm}.
$$

Set-II. For the same gap distance the slope will be $\alpha = \ln(12/8)/0.05 = 8.1$ collisions/cm and therefore

$$
\frac{I}{I_0} = 2 \times 10^5 = \frac{e^{8.1 \times 0.5}}{1 - \gamma(e^{4.05} - 1)}
$$

$$
2 \times 10^5 = \frac{57.39}{1 - \gamma(56.39)}
$$

$$
\frac{200 \times 10^3}{57.39} = 3.4849 \times 10^3 = \frac{1}{1 - 56.39\gamma}
$$

$$
2.87 \times 10^{-4} = 1 - 56.39\gamma
$$

$$
56.39\gamma = 1.0
$$

$$
\gamma = 1.7 \times 10^{-2} \text{ collisions/cm}
$$

 $\overline{\text{O}}$

0f